Quantum gravity and cosmological perturbations:



Effects of a modified Friedmann equation on gauge-invariant perturbations

arXiv:2211.04500

Lisa Mickel with Steffen Gielen

Cosmology: a testing ground for quantum gravity

- Explicit calculations in quantum gravity models are difficult
- Simpler in cosmological scenario due to isotropy and homogeneity
- First test for quantum gravitational models: Recover FLRW dynamics at late times
 - \rightarrow Can we also make statements about perturbations?

Cosmological perturbation theory

- Include linear inhomogeneous perturbations in the
 - o **metric**, e.g.

$$a(t) \rightarrow a(t)(1 - \psi(t, x^i))$$

 \circ and **matter** content $\rho(t) \to \rho(t) + \delta \rho(t, x^i)$

Modified Friedmann equation

• **General form** of a modified Friedmann equation that encodes the universe dynamics

Hubble parameter
$$H=\frac{a'}{a}$$
 $H^2=\frac{8\pi G}{3}$ $N^2 \rho \mathcal{F}$ Generic function that encodes modification GR: $\mathcal{F}=1$

Unchanged matter dynamics

Example: A bouncing universe in Group Field Theory cosmology

- Field theory on a group manifold: quanta of space
- Evolution of the universe from the expectation value of the volume operator over coherent state $\ \langle \hat{V} \rangle = a^3$
- Spectrum of \hat{V} bounded from below \longrightarrow **bounce**
- Evolution with respect to matter clock

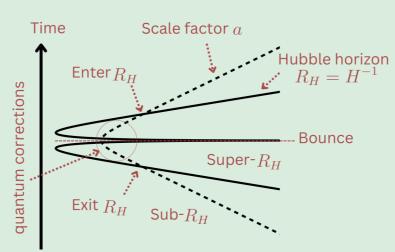
Separate Universe Picture



- Model the universe as an **ensemble of independent homogeneous patches**
- Each patch behaves like an FLRW universe
- Perturbations = deviations from the mean

$$\psi_p = \frac{1}{3} \left(1 - \frac{V_p}{V_{bg}} \right) \quad V_{bg} := \frac{1}{N_{\text{patches}}} \sum_p V_p$$

Perturbation Equations



- Before and after the bounce modes are super-horizon and approximately homogeneous
- Work in the SU picture and **perturb the Friedmann equation** at linear order to obtain perturbation
 equations

$$H\psi' = -H^2 \left(\tilde{\Phi} + \frac{\delta \rho}{2\rho} + \frac{\delta \mathcal{F}}{2\mathcal{F}} \right)$$

Relational Clock

• Massless scalar field ϕ follows Klein-Gordon equation

$$\phi'' - \frac{N'}{N}\phi' + 3H\phi' = 0$$

 Assign value of scalar field to each quanta of space



ightarrow time coordinate

Gauge invariant perturbation variables

- Invariance under **infinitesimal diffeomorphisms** acting on perturbed coordinate system \rightarrow relate to **physical quantities** (e.g. CMB)
- Comoving curvature perturbation

$$\mathcal{R} = \psi + \frac{H}{\phi'} \delta \phi$$

- o Conserved on super-horizon scales in GR
- In the comoving gauge $\delta \phi = 0 o \mathcal{R} = \psi$

Special Case $\mathcal{F} = \mathcal{F}(\rho)$

- Modification ${\cal F}$ is a function of the energy density only, such that $\,\delta{\cal F}\propto\delta\rho\,$
- \rightarrow Analogue to the diffeomorphism constraint of GR:

$$\psi' + H\tilde{\Phi} = 0$$

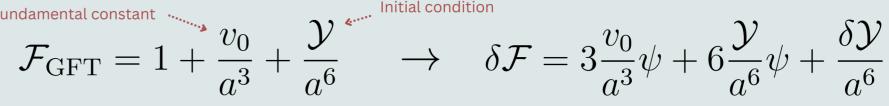
- \Longrightarrow Recover GR conservation law for $\mathcal R$
- Generalises the LQC case $\,\mathcal{F}_{\mathrm{LQC}} = 1 rac{
 ho}{
 ho_{\mathrm{c}}}$

General case $\mathcal{F} \neq \mathcal{F}(\rho)$

For a massless scalar field

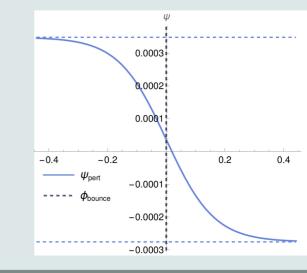
$$\mathcal{R}' = \psi' = -H \frac{\delta \mathcal{F}}{2\mathcal{F}}$$

• Consider the **GFT example**



 \Longrightarrow Dynamics in ${\mathcal R}$ around bounce

 Evolution from perturbation equations agree with quantum evolution up to second order effects



Conclusion

- Statements about large scale perturbations from a modified Friedmann equation
 - without specifying inhomogeneities in the quantum framework

Outlook

- Study **inhomogeneities in the quantum framework** (model dependent)
 - GFT: Add additional scalar fields to obtain a relational framework