

# Parameterizations of black-hole spacetimes beyond circularity

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(in collaboration with Astrid Eichhorn and Aaron Held)

## Why black holes beyond GR?

- Test of validity of GR from near-horizon observations  
[Broderick '21; E.H.T. Collaboration '19, '22]
- Learn about true nature of black-holes  
→ Glimpse on quantum gravitational effects?  
→ How black-holes beyond GR look like?  
[Held '19]
- Usual parameterized deviations from Kerr metric assume symmetries which might not hold beyond GR:
  - Circularity  
[Lewis '32; Papapetrou '66; Konoplya, Rezzolla, Zhidenko '16]
  - Hidden constant of motion  
[Benenti, Francaviglia '79; Vigeland, Yunes, Stein '11; Johannsen '13]

## Parameterizations and symmetries

**1<sup>st</sup> key result:** Usual parameterizations assume some symmetries

	Symmetries	Metric components	Free functions
BHs:	axisymmetry + stationarity	$10 \rightarrow 6$	at least 6
add:	circularity	5	4 (LP form)
add:	hidden constant of motion	5	10 of 1 coord.

- Most general circular metric: [Johannsen '13]

$$ds_{RZ}^2 = -g_{tt}dt^2 - 2g_{t\phi_{BL}}dt d\phi_{BL} + g_{\phi_{BL}\phi_{BL}}d\phi_{BL}^2 + g_{rr}dr^2 + g_{\theta\theta}d\theta^2$$

→ 5 non-zero metric components & 5 free functions ✓

→ 1 free function can be removed by coordinate transformation  
⇒ 5 metric components & 4 free functions in some coordinates

- Most general circular metric with new Carter-like hidden constant of motion:  
[Benenti, Francaviglia '79]

$$g^{\mu\nu}\partial_\mu\partial_\nu = \frac{1}{S_{x_1} + S_{x_2}} \left[ (G_{x_1}^{ij} + G_{x_2}^{ij}) \partial x_i \partial x_j + \Delta_{x_1} \partial x_1^2 + \Delta_{x_2} \partial x_2^2 \right]$$

with  $i, j$  indices related to Killing coordinates, and 1, 2 to explicit coordinates

→ Forms a subclass of circular metrics

**2<sup>nd</sup> key result:** Non-circular parameterization of deviations from Kerr (HP coordinates)

$$ds_{NC}^2 = - \left( \frac{r^2 - 2Mr + a^2\chi^2}{r^2 + a^2\chi^2} \right) (1 + \Delta_1(r, \chi)) du^2 + 2(1 + \Delta_2(r, \chi)) du dr - 4 \frac{Mar}{r^2 + a^2\chi^2} (1 - \chi^2)(1 + \Delta_3(r, \chi)) du d\phi - 2a(1 - \chi^2)(1 + \Delta_4(r, \chi)) dr d\phi + \frac{r^2 + a^2\chi^2}{1 - \chi^2} (1 + \Delta_5(r, \chi)) d\chi^2 + \frac{1 - \chi^2}{r^2 + a^2\chi^2} \left( (a^2 + r^2)^2 - a^2 (r^2 - 2Mr + a^2) (1 - \chi^2) \right) (1 + \Delta_6(r, \chi)) d\phi^2$$

- Asymptotic flatness:  $\Delta_{HP,i}(r, \chi) \xrightarrow{r \rightarrow \infty} 0$  ✓
- Correct Newtonian limit:  $\Delta_{HP,i}(r, \chi) \sim \mathcal{O}(\frac{1}{r^2})$  ✓
- Flat Minkowski limit: only reached for  $M \rightarrow 0$  if  $\Delta_{HP,i}(r, \chi) \sim M$  & no extra-hair
- Limit  $a \rightarrow 0$  breaks spherical symmetry due to  $\chi$ -dependence of  $\Delta_{HP,1/2}$  and  $\Delta_{HP,1/2/5/6}$  encoding additional (quantum) hair
- Can reduce to circular spacetimes if e.g.  $\Delta_{HP,i} = \varepsilon, \forall i$  or for  $\Delta_{HP,5} \neq 0, \Delta_{HP,i \neq 5} = 0$  (non-trivial) ✓

Challenge: directly relating non-circular HP parameterization to general circular BL parameterization

→ Possible via a counting argument: 14 free functions – 9 constraints = 5 free functions ✓

## What are circularity & hidden constant of motion?

- Consider axisymmetric, stationary & asymptotically flat spacetimes  
→ Two commuting Killing vectors  $\xi_1, \xi_2$
- Circularity holds if [Papapetrou '66]
 
$$\xi_1^\mu R_\mu^{[\nu} \xi_2^{\kappa]} = 0 \text{ everywhere,}$$

$$\xi_2^\mu R_\mu^{[\nu} \xi_1^{\kappa]} = 0 \text{ everywhere.}$$
 → Implies an isometry. In BL coordinates,  $t \rightarrow -t$  &  $\phi_{BL} \rightarrow -\phi_{BL}$
- In GR circularity holds because  $R_{\mu\nu} = 0$ . But some non-vacuum GR & beyond GR spacetimes break circularity  
[Ioka, Sasaki '03, '04; Birk, Stergioulas, Muller '11]  
[Anson, Babichev, Charmousis, Hassaine '20; Minamitsuji '20]
- Killing vectors  $K_\mu \rightarrow$  isometries  
Rank-2 Killing tensors  $K_{\mu\nu} \rightarrow$  hidden constants of motion  $C = K_{\mu\nu} u^\mu u^\nu$   
→ Leads to separability of geodesic equations

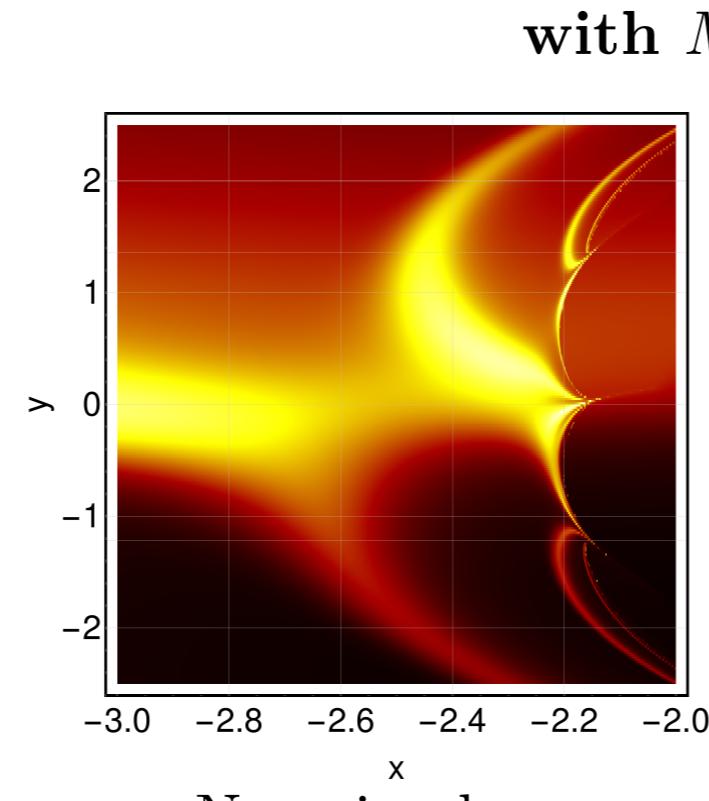
## Which choice of coordinates?

Boyer-Lindquist ( $t, r, \theta, \phi_{BL}$ )	Horizon-penetrating ( $u = t - r^*, r, \chi, \phi$ )
Requires fine tuning at horizon [Johannsen '13]	Singularity-free at horizon
<b>Circular spacetime metrics</b>	<b>Non-circular spacetime metrics</b>

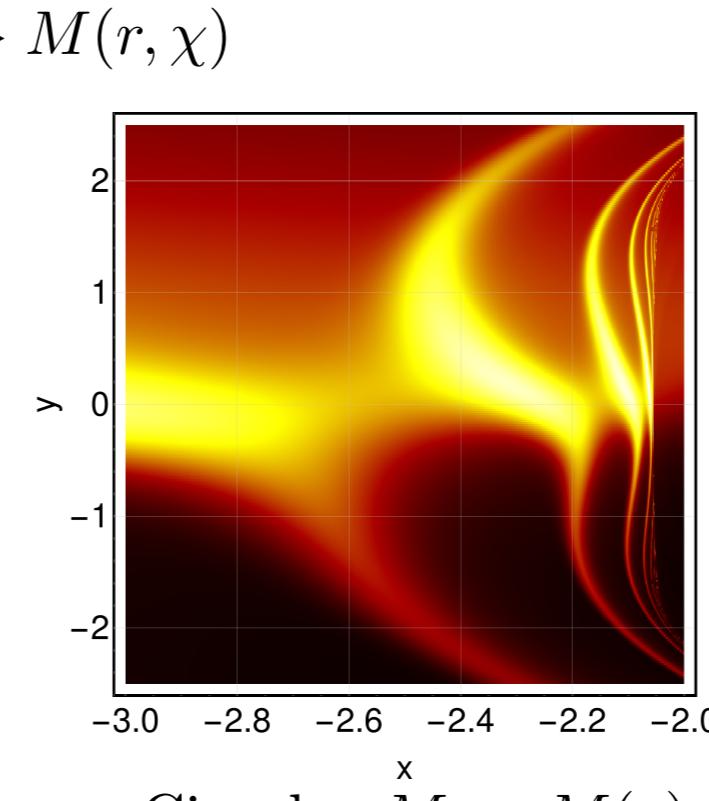
## Open question

- Minimal, general parameterization of axisymmetric, stationary & asymptotically flat BH spacetimes?

**Locality-based regular, non-circular black-hole with  $M \rightarrow M(r, \chi)$**



Non-circular



Circular  $M \rightarrow M(r)$

**3<sup>rd</sup> key result:**  
Specific image features:

- Cusps (in shadow boundary and photon rings)
- A dent (in  $y = 0$  image axis)
- A broken reflection symmetry (about  $y = 0$  image axis)

**4<sup>th</sup> key result:**  
No equivalent for  $a \rightarrow a(r, \chi)$

- Can we establish a 1-to-1 map between non-circularity and specific image features?