# Calibrating the continuum approximation of discrete quantum gravity









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## Outline

- The Causal Set Way : Quantising the causal structure

  - Does the Continuum emerge from Dynamics?
  - Uniqueness?
- Calibrating the continuum approximation: new techniques

The Continuum Approximation of quantum gravity: With or Without Discreteness

• The Continuum Approximation: random discretisation

#### The Continuum Approximation: With or Without Discreteness

• Different Physics at different scales:

- Quantum Gravity DOF QG: Discrete or Continuum
- $\mathbb{QG} \sim (M,g)$  at  $V > V_p$

• How do we recognise the continuum approximation when we see it?

- GR:  $V \gg V_p$ , Planck scale:  $V \sim V_p$ , Trans-Planckian scale:  $V \ll V_p$  (?)
- GR emergent from Planck scale : what Planck scale physics should one forget to get GR?

### Continuum or Discrete DOF

- If QG is a quantum spacetime geometry (or any other classically equivalent set of continuum quantities), with some continuity requirements:  $C^2, C^1, C^0...$  OR is discrete (graph, triangulation, network, causal set, ...)
- What distinguishing features does  $\mathbb{QG}$  have at the Planck scale  $V_p$ ? (Is there physics) below  $V_p$  for continuum theories?)
- Say your theory spits out a (coherent state of) continuum geometry  $(M,g) \in \mathcal{M}_r$  but r < 2 or some pre-geometric **Q** (example: piece-wise continuous, orbifolds, extra dimensions, etc.)
- Does (M,g) or **Q** approximate  $(\mathbf{M},\mathbf{g}) \in \mathcal{M}_2$  (or desired differentiability) at a scale  $V \gg V_p?$

## Uniqueness

Is this approximation unique? 

• What does it mean for two spacetimes to be close at a given scale?





- Bombelli, 2000, Bombelli and Noldus, 2004 -Burtscher and Allen, 2021, --Kunzinger and Steinbauer, 2021, ...

A

#### $\mathbb{QG} \sim_V (\mathbf{M}, \mathbf{g}) \text{ AND } \mathbb{QG} \sim_V (\mathbf{M}', \mathbf{g}') \text{ are } (\mathbf{M}, \mathbf{g}) \text{ and } (\mathbf{M}', \mathbf{g}') \text{ "close" at } V \gg V_p$ ?



 $\boldsymbol{B}$ 

## The Causal Set Way: Quantising the Causal Structure

Lorentzian Spacetime (Causal, Distinguishing) =  $(M, \prec) + \epsilon$ 

 $(M, \prec)$  is a poset :(i) Acyclic:  $x \prec y, y \prec x \Rightarrow x = y$  (ii) Transitive:  $x \prec y, y \prec z \Rightarrow x \prec z$ 

The Causal Set Hypothesis:

*-Myrheim*, 1978 -Bombelli, Lee, Meyer and Sorkin, 1987

1. Locally finite posets or Causal Sets are the fine grained structure of spacetime

2. Continuum Approximation: *Order + Number ~ Spacetime* (counting replaces volume)



-Hawking, Hawking-King-MacCarthy, Malament, Kronheimer-Penrose

...Robb, Zeeman, Penrose, Kronheimer, Finkelstein, Myrheim, Hemion, 't Hooft ....





### The Continuum Approximation

- $n \sim \rho V$  correspondence has to be diffeo invariant
- Random discretisation via a Poisson sprinkling process:

• 
$$P_V(n) = \frac{(\rho V)^n}{n!} e^{-\rho V}, \langle n \rangle = \rho$$

- For every causal spacetime (M, g) there is a kinematic ensemble  $\{C\}_{\rho}$  (first quantisation)
- $C \sim_{\rho} (M, g)$  is a faithful embedding at density  $\rho$  if :
  - $C \hookrightarrow (M, g)$  is order preserving



 $\rho$  sets a scale

•  $n_V$ : number of points in spacetime volume V is a random variable,  $P_V(n) = \frac{(\rho v)}{n!} e^{-\rho V}$ 



#### Causal Set Non-locality

- The nearest neighbours lie all along the light cone
- A continuum-like causal set is a graph without a fixed valency •

Preserving Local Lorentz invariance <u>Theorem</u> : — Bombelli, Henson and Sorkin,2006

There is no measurable map  $D: \Omega \to H$  which is equivariant, i.e.,  $D \circ \Lambda = \Lambda \circ D$ .

<u>Proof</u>: If such a map existed, then  $\mu_D = \mu \circ D^{-1}$  is a Lorentz invariant probability measure on H which is not possible since is H non-compact.





- Measure triple  $(\Omega, \Sigma, \mu)$
- unit hyperbola  $H \subset \mathbb{M}^d$  (fdtl directions)



### Calibrating the Continuum Approximation

• Is this approximation unique?

#### **FUNDAMENTAL CONJECTURE:**

- Order Invariants  $\mathcal{O}$  as Geometric Invariants  $\mathbb{G}$
- WEAK FORM OF FUNDAMENTAL CONJECTURE:

If  $C \sim_{(\rho, \mathcal{O})} (\mathbf{M}, \mathbf{g})$  AND  $C \sim_{(\rho, \mathcal{O})} (\mathbf{M}', \mathbf{g}')$  then  $(\mathbf{M}, \mathbf{g}) \sim_{(\rho, \mathbb{G})} (\mathbf{M}', \mathbf{g}')$ 

• Which order invariants  $\mathcal{O}$  correspond to Geometric Observables  $\mathbb{G}$ ?

If  $C \sim_{\rho} (\mathbf{M}, \mathbf{g})$  AND  $C \sim_{\rho} (\mathbf{M}', \mathbf{g}')$  then  $(\mathbf{M}, \mathbf{g})$  and  $(\mathbf{M}', \mathbf{g}')$  are "close" at some  $\rho^{-1} \gg V_p$ .



### Geometric Reconstruction: geometry from counting

- Dimension Estimators Myrheim, Myer, Glaser & Surya, ...
- Timelike Distance — Brightwell & Gregory
- Spatial Homology Major, Rideout & Surya
- Spatial and Spacelike Distance Rideout & Wallden, Eichhorn, Mizera & Surya, Eichhorn, Surya & Versteegen
- Sorkin, Henson, Benincasa & Dowker, Dowker & Glaser • D'Alembertian
- Benincasa-Dowker-Glaser Action Benincasa & Dowker, Dowker & Glaser
- GHY terms in the Action Buck, Dowker, Jubb & Surya
- Locality and Interval Abundance Glaser & Surya
- Horizon Molecules Barton, Counsell, Dowker, Gould & Jubb, Machet and Wang
- Scalar Field Greens functions Johnston, Dowker, Surya & Nomaan X
- Scalar Field SJ vacuum Johnston, Sorkin, Yazdi, Nomaan X, Surya
- Entanglement Entropy Dou & Sorkin, Sorkin & Yazdi, Yazdi, Nomaan X, Surya
- Null Geodesics from Ladder molecules -- with A. Bhattacharya and A. Mathur, 2022







#### **Example: Dimension Estimator**

• If  $C \sim_{\rho} (\mathbb{M}^d, \eta)$ 

• If 
$$C \sim_{\rho} (\mathbf{M}', \mathbf{g}')$$
, then  $(\mathbf{M}', \mathbf{g}') \sim_{(\rho, d)} (\mathbb{M}^d, \eta)$ 

- Example  $\mathbf{M}' = \mathbb{R}^d \times S^1$ ,  $\mathbf{g} = \eta \oplus l$ : For  $\rho^{-1} > V_c \times |S^1|$ ,  $C \sim_{\rho,d} (\mathbf{M}', \mathbf{g}')$
- dimension  $d' \neq d$ .

# • Myrheim-Myer dimension estimator $\langle r \rangle = \frac{2\langle R \rangle}{\langle n \rangle^2} = \frac{\Gamma(d+1)\Gamma(d/2)}{4\Gamma(3d/2)}$

• Therefore if C has Myrheim-Myer dimension d, then it cannot approximate a spacetime of

## But is the continuum emergent from dynamics?

#### Lorentzian Path Sum over $\Omega_n$

•  $\Omega_n$  : sample space of all n-element causal sets

• 
$$|\Omega_n| \sim 2^{\frac{n^2}{4} + \frac{3n}{2} + o(n)}$$

- Typical causal sets are Kleitmann-Rothschild:  $|\Omega_{KR}| \sim 2^{\frac{n^2}{4} + \frac{3n}{2} + o(n)}$
- Other layered Posets are subdominant: ~  $2^{c(d)n^2 + o(n^2)}$ ,







- Kleitman and Rothschild, Trans AMS, 1975 - J. Henson, D. Rideout, R. Sorkin and S.Surya, JEM, 2015

 $c(d) \le 1/4$ 

-D. Dhar, JMP, 1978 - Promel, Steger, Taraz 2001









### Do manifold-like causal sets stand a chance?

#### Layered Posets are not manifold-like



#### The Benincasa-Dowker-Glaser Action

$$S_{BDG}^{(d)}(C) = \mu \left( n + \sum_{j=0}^{j_{max}} \lambda_j N_j \right)$$
$$S_{BDG}^{(4)} = \frac{4}{\sqrt{6}} \left( n - N_0 + 9N_1 - 16N_2 + 8N_3 \right)$$

For the KR poset:



— Benincasa & Dowker 2010, — Dowker & Glaser 2011 — Glaser 2012

#### Do manifold-like causal sets stand a chance? Yes!

Bilayer Posets :

Suppression for:

/

$$\tan(-\frac{\mu\lambda_0}{2}) > \sqrt{3}$$

$$d = 4, \quad \mu = \left(\frac{l}{l_p}\right)^2 \Rightarrow l \approx 1.452 l_p$$

The discrete Einstein Hilbert action in any dimension suppresses all k-layer orders for k < < n: Action wins over Entropy



—Loomis and Carlip, 2017 — A.Anand Singh, A.Mathur and Surya, 2021 —P. Carlip, S. Carlip and S. Surya, 2022 P. Carlip, S. Carlip and S. Surya, in preparation



 $\sim n/4$ 







Why d = 4 spacetime?

Does the *d* dimensional discrete Einstein Hilbert action suppress  $d' \neq d$ ? (Some hints..)

#### State Sum Models: Lorentzian Statistical Geometry

$$Z_{\beta} = \sum_{c \in \Omega} \exp(i\beta S_{BDG}(c)),$$

Inverse "temp":  $\beta \rightarrow i\beta$  $Z_{\beta} \rightarrow \tilde{Z}_{\beta} = \sum \exp(-\beta S_{BDG}(c))$  $c \in \Omega$ 

![](_page_17_Figure_3.jpeg)

 $S^1 \times \mathbb{R}$ 

![](_page_17_Figure_5.jpeg)

### Track Observables O for continuum-non-continuum phases

<ul> <li>Myrheim-Myer dimension estimator</li> </ul>	0
<ul> <li>Interval Abundance</li> </ul>	-100 Action -200
<ul> <li>Action</li> </ul>	-300
• Height	25
Finite size scaling to know if $\langle \mathcal{O} \rangle \lim_{n \to \infty} \mathbb{G}$	20 Height 15 10
	5

Computationally very expensive!!

![](_page_18_Figure_3.jpeg)

- Cunningham & Surya, 2020

![](_page_18_Picture_5.jpeg)

## Is the Continuum Unique?

### Calibrating the Continuum Approximation: new techniques

- $(M,g) \sim_{\rho} (M',g')$ : what does this mean?
- Lorentzian spacetimes via GH convergence
- Can we define  $\rho$ -closeness?
- $\bullet$  Convergence in  $\rho \rightarrow \infty$  limit

Bombelli, 2000, Bombelli and Noldus, 2004
 Burtscher and Allen, 2021,
 -Kunzinger and Steinbauer, 2021

-- Bombelli and Meyer, 1989 -- Minguzzi and Suhr , 2022 -- Muller, 2022

#### **Uniform Random Sampling Method**

![](_page_21_Figure_1.jpeg)

![](_page_21_Picture_2.jpeg)

#### Lorentzian Uniform Random Sampling - Causal Sets

- Consider two spacetimes (M, g), (M', g') of volume ~ V
- Random sampling produces a causal set c by using causality relation  $\prec$  in (M, g)
- $\Omega_n$  : ensemble of *n*-element causal sets
- $P_n(c | M)$ ,  $c \in \Omega_n$  is a probability distribution.
- Since  $\sum P_n(c \mid M) = 1$ ,  $\sqrt{P_n(c \mid M)}$  form coordinates on positive part of the sphere in  $\mathbb{R}^{|\Omega_n|}$  $c \in \Omega_{n}$

• 
$$d_n(M, M') = \frac{2}{\pi} \cos^{-1} \left( \sum_{c \in \Omega_n} \sqrt{P_n(c \mid M)} \sqrt{P_n(c \mid M)} \right)$$

Closeness function but not a distance function

![](_page_22_Figure_8.jpeg)

![](_page_22_Figure_9.jpeg)

![](_page_22_Figure_10.jpeg)

![](_page_22_Picture_11.jpeg)

#### $(B, d_R)$ Distance between Abstract Metric Spaces $(A, d_A)$ ,

 $d_{GH} \equiv \inf_{(M,d),\Phi,\Psi} d_H(\Phi(A),\Psi(B))$ 

• Calculating this quantity explicitly is very hard!

![](_page_23_Figure_5.jpeg)

$$d_{GH}(\mathbb{S}^m, \mathbb{S}^n) \leq \frac{\pi}{2}$$
$$d_{GH}(\mathbb{S}^0, \mathbb{S}^n) = \frac{\pi}{2}$$
$$d_{GH}(\mathbb{S}^m, \mathbb{S}^\infty) = \frac{\pi}{2}$$

Lim, Memoli & Smith, 2022

 $p_2$ 

#### • Let (M, d) be a metric space such that $\Phi : A \hookrightarrow M$ , $\Psi : B \hookrightarrow M$ are two **isometric embeddings**

• Gromov-Hausdorff Distance: **shortest** Hausdorff distance over all possible isometric embeddings

![](_page_23_Figure_12.jpeg)

## Null Distance Function for GH Convergence

- Time function  $T: (M, d) \rightarrow \mathbb{R}$  -- monotonic, etc.
- $\gamma(p,q)$  : piece-wise causal curve: allow it to go backward and forward in time!
- The length along the curve is defined as  $L_T(\gamma(p,q)) = \sum |T(s_i) T(s_{i-1})|$
- Null Distance function:  $d_T(p,q) \equiv \inf_{T} L_T(\gamma(p,q))$  IS a genuine distance function  $\gamma(p,q)$
- Can be used to study convergence of FRW-type spacetimes

Workshop on Non-regular spacetimes, ESI, Vienna, March 2023

![](_page_24_Picture_7.jpeg)

Sormani and Vega, 2016

Allen and Burtscher, 2020

![](_page_24_Figure_10.jpeg)

Allen and Burtscher, 2020

#### GH-like distance for 2d orders using a lattice embedding

- $S = (1, 2, ..., n), U = (u_1, ..., u_n), V = (v_1, ..., v_n), u_i \in S, v_i \in S$
- 2d order  $C = U \cap V$ :  $e_i = (u_i, v_i) \prec e_i = (u_i, v_i) \Leftrightarrow u_i < u_i, v_i < v_i$
- Examples:
  - $u_1 < u_2 \dots < u_n, v_1 < v_2 \dots < v_n \Rightarrow C$  is a chain
  - $u_1 < u_2 \dots < u_n, v_n < v_{n-1} \dots < v_1 \Rightarrow C$  is an antichain
  - U, V randomly sampled : random 2d order ~  $(\mathbb{D}^2, \eta)$
- ullet Every 2d order can be embedded as a 2d order into the light cone lattice  $\,\mathscr{L}$
- The null distance function on  $\mathscr{L}: d_t(a, b) = \frac{1}{2}(|u_b|)$

• 
$$A, B \subseteq \mathcal{L}, \quad d_H(A, B) = \sup \inf_{a \in A} \inf_{b \in B} d_t(a, b)$$

• Let  $c_1, c_2 \in \Omega_{2d}$ ,  $\mathscr{E}_i : c_i \hookrightarrow \mathscr{L}$ ,  $d_{GH}(c_1, c_2) \equiv \inf_{\mathscr{E}_i} d_H^{\leftrightarrow}(\mathscr{E}_1(c_1), \mathscr{E}_2(c_2))$ 

-Work in progress with Alan Daniel Santhosh

$$-u_{a}|+|v_{b}-v_{a}|)$$

![](_page_25_Figure_18.jpeg)

![](_page_25_Picture_20.jpeg)

#### Preliminary calculations...

- $d_{GH}(a_n, a_{n+1}) = 1$ ,  $d_{GH}(c_n, c_{n+1}) = 1$
- $d_{GH}(a_n, c_n) = m$ , n = 2m or n = 2m + 1

• 
$$d_{GH}(B_2, c_n) = \frac{n}{4}$$
,  $d_{GH}(B_2, a_n) = \frac{n}{4} + \frac{1}{2}$ 

• 
$$d_{GH}(KR, c_n) \le \frac{n}{2}$$
,  $d_{GH}(KR, a_n) = \frac{n}{4}$ 

• 
$$d_{GH}(L_4, c_n) \le \frac{n}{8}$$
,  $d_{GH}(L_4, a_n) \le \frac{3n}{8}$ 

- Distance between Antichain  $a_n$  and Chain  $c_n$  grows the fastest
- Distance between the K-layer poset -- does it get closer to  $c_n$  than  $a_n$  as K increases?
- Measuring distance between random orders : challenging and may need numerical work

![](_page_26_Figure_9.jpeg)

## In Conclusion...

- - QG ~  $(\mathbf{M}, \mathbf{g})$
- In CST, the continuum approximation is recognised using geometric order invariants O
- New Lorentzian geometric tools to calibrate how close Lorentzian QG is to a smooth

spacetime..

• Can they help us determine this up to a scale:  $(\mathbf{M}, \mathbf{g}) \sim_V (\mathbf{M}', \mathbf{g}')$ ?

![](_page_27_Picture_8.jpeg)

• The Continuum Approximation could be relevant to many approaches to quantum gravity:

• Uniqueness: If  $\mathbb{QG} \sim_V (\mathbf{M}, \mathbf{g})$  AND  $\mathbb{QG} \sim_V (\mathbf{M}', \mathbf{g}')$  then  $(\mathbf{M}, \mathbf{g}) \sim_V (\mathbf{M}', \mathbf{g}')$ 

Thank you!