# Impact of matter fields

CANA ROLLA

#### on quantum geometry

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# Introduction to Causal Dynamical Triangulations

#### What is Causal Dynamical Triangulations?

**Causal Dynamical Triangulations** is a background-independent and diffeomorphism-invariant approach to quantum gravity. It provides a lattice regularization of the formal gravitational path integral via a sum over causal triangulations.



# **Causal Dynamical Triangulations**

The partition function and action

$$\int \mathcal{D}[g] e^{iS^{EH}[g]} \longrightarrow \sum_{\mathcal{T}} e^{-S^{R}[\mathcal{T}]}$$

The Regge action  $S^{R}[T]$  is equal to the Einstein-Hilbert action  $S^{EH}$  evaluated on a triangulation T built of *four-simplices*.

- Global time foliation of the spacetime manifold (Causal DT)
- Fixed spatial and global topology ( $\mathcal{M} = T^3 \times T^1$ )
- Monte Carlo simulations expectation values of observables



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#### **Dynamical scalar fields**

- We introduce dynamical (backreacting) scalar fields φ<sub>σ</sub>, σ = x, y, z, t, taking values in a circle of circumference δ and winding around S<sup>1</sup> once in the direction x, y, z, and t, respectively.
- The continuous Euclidean action for a massless scalar field:

$$\mathcal{S}[g,arphi] = rac{1}{2}\int \mathrm{d}^4x\,\sqrt{g(x)}\;\partial^\muarphi(x)\partial_\muarphi(x),\quad arphi(x)\in \mathcal{S}^1(\delta)$$

The discrete counterpart of the matter action decomposes into the quantum and the classical parts:

$$\begin{split} S[\mathcal{T}, \varphi &= \bar{\varphi} + \eta] = \sum_{i \leftrightarrow j} (\varphi_i - \varphi_j - \delta \mathbf{B}_{ij})^2 = \eta^\top \mathbf{L} \eta + \delta^2 \tilde{S}^{\text{clas}}[\mathcal{T}] \\ \tilde{S}^{\text{clas}}[\mathcal{T}] &= \tilde{\varphi}^\top \mathbf{L} \tilde{\varphi} - 2b^\top \tilde{\varphi} + \|\mathbf{B}\|^2, \quad L \tilde{\varphi} = b \end{split}$$

L is the discrete Laplacian matrix and B is an antisymmetric *boundary* matrix.

#### Single field winding in time direction

Single dynamical scalar field winding once around a circle of circumference  $\delta$  in the time direction. For large  $\delta$ , the volume profile is described by a *pinched* cosine function.



#### Minisuperspace model with scalar field

The behavior  $\langle n_t \rangle \propto \cos\left(\frac{2\pi}{\tau}t\right)$ , can be explained by assuming spatial homogeneity and isotropy.

Minisuperspace action ( $v = v(t) \propto a^3(t), \ \varphi = \varphi(t)$ )

$$S[v,\varphi] = \int \mathrm{d}^4 x \sqrt{g} \left( \# R - \Lambda + (\partial \varphi)^2 \right) = \int_{-T/2}^{T/2} \mathrm{d}t \, \frac{\dot{v}^2}{v} + v \dot{\varphi}^2$$

Constraints

$$V = \int \mathrm{d}t \, \mathbf{v}, \quad \delta = \int \mathrm{d}t \, \dot{\varphi} = \varphi\left(\frac{T}{2}\right) - \varphi\left(-\frac{T}{2}\right), \quad \mathbf{v}(t) \ge \varepsilon$$

Constant solution for  $\delta \leq 2\pi$ 

$$v(t) = rac{V}{T}, \qquad \dot{\varphi}(t) = rac{\delta}{T}, \qquad S = rac{V}{T^2} \cdot \delta^2$$

Cosine solution for  $\delta \geq 2\pi$ 

$$egin{aligned} \mathsf{v}(t) = egin{cases} c \cdot \cos\left(rac{2\pi}{ au}t
ight) + c + arepsilon & |t| \leq rac{ au}{2} \ arepsilon & rac{ au}{2} \leq |t| \leq rac{ au}{2} \ , & \dot{arphi}(t) = rac{eta}{ extsf{v}(t)} \end{aligned}$$

#### Minisuperspace model with scalar field

#### **Classical solutions**



Phase transition at  $\delta = \frac{2\pi}{\sqrt{G}}$ 

#### **Density maps**

▶ The matter action decomposes into quantum and classical parts:

$$S[\mathcal{T}, \varphi = \bar{\varphi} + \eta] = \eta^{\mathsf{T}} \mathbf{L} \eta + \delta^2 \tilde{S}^{\mathrm{clas}}[\mathcal{T}]$$

- Quantum fluctuations can be integrated out and are negligible.
- The nontrivial classical solution contributes to the classical action δ<sup>2</sup> Š<sup>clas</sup>[T] which depends in a crucial way on the triangulation T.
- For small  $\delta$  the geometric part of the action dominates.
- For large δ the matter action dominates and the total action is the lowest for *pinched* configurations.
- To visualize a triangulation and a field configuration, we reintroduce four classical scalar fields. Each four-simplex is assigned coordinates (φ<sub>x</sub>, φ<sub>y</sub>, φ<sub>z</sub>, φ<sub>t</sub>) and is visible as a single pixel on the following images.
- Dense fibers correspond to evolving outgrowths.



Dynamical fields  $\varphi^{\mathbf{x}}, \varphi^{\mathbf{y}}, \varphi^{\mathbf{z}}.$ Projection on  $\varphi^t - \varphi^x$ . Circumference  $\delta = 0.0.$ 







Dynamical fields  $\varphi^{\mathbf{x}}, \varphi^{\mathbf{y}}, \varphi^{\mathbf{z}}.$ Projection on  $\varphi^t - \varphi^x$ . Circumference  $\delta = 3.0.$ 3000



Dynamical fields  $\varphi^{\mathbf{x}}, \varphi^{\mathbf{y}}, \varphi^{\mathbf{z}}.$ Projection on  $\varphi^t - \varphi^x$ . Circumference  $\delta = 4.0.$ 2000 200000 250000 200000 150000 100000



Dynamical fields  $\varphi^{\mathbf{x}}, \varphi^{\mathbf{y}}, \varphi^{\mathbf{z}}.$ Projection on  $\varphi^t - \varphi^x$ . Circumference  $\delta = 7.0.$ mus in directions x, u, z  $\langle n_l \rangle =$ 40000 30000 2000 1000 40000 200000 200000 100000



Dynamical fields  $\varphi^{\mathbf{x}}, \varphi^{\mathbf{y}}, \varphi^{\mathbf{z}}.$ Projection on  $\varphi^t - \varphi^x$ . Circumference  $\delta = 8.0.$ Jumps in directions x, u, z  $\langle n_l \rangle =$ 40000 30000 400000 200000

200000



Dynamical fields  $\varphi^{\mathbf{x}}, \varphi^{\mathbf{y}}, \varphi^{\mathbf{z}}.$ Projection on  $\varphi^t - \varphi^x$ . Circumference  $\delta = 10.0.$ Jumps in directions z. u. z 1,400.00 12000 100000 80000 60000 40000 20000 400000 300000

200000

#### Volume profile for circumference $\delta = 7.0$













#### Spherical vs toroidal spatial topology



The difference between **spherical** and **toroidal spatial topology** is visible in the volume profile - as predicted by the minisuperspace model.

# Summary

- Causal dynamical triangulations is a model of generic geometry fluctuations at the Planck scale.
- Introduction of dynamical scalar fields with matching topological boundary conditions has a dramatic effect on the geometries that dominate the CDT path integral.
- Scalar fields induce a new type of phase transition, where the effective spacetime topology changes from toroidal to a simply connected one.
- This new kind of coupling between the topology of the matter fields and the topology of spacetime is likely to result in a **phase transition** for sufficiently strong coupling.
- The simple minisuperspace model predicts a *phase transition* when changing the circumference  $\delta$ .
- ► The classical limit agrees with *the minisuperspace model*.





#### In collaboration with

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