

General Relativistic Decoherence with Applications to Dark Matter

Mark Hertzberg, Tufts

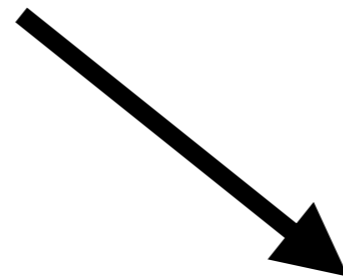
Quantum Gravity 2023, July 14 2023

Based on work with Itamar Allali

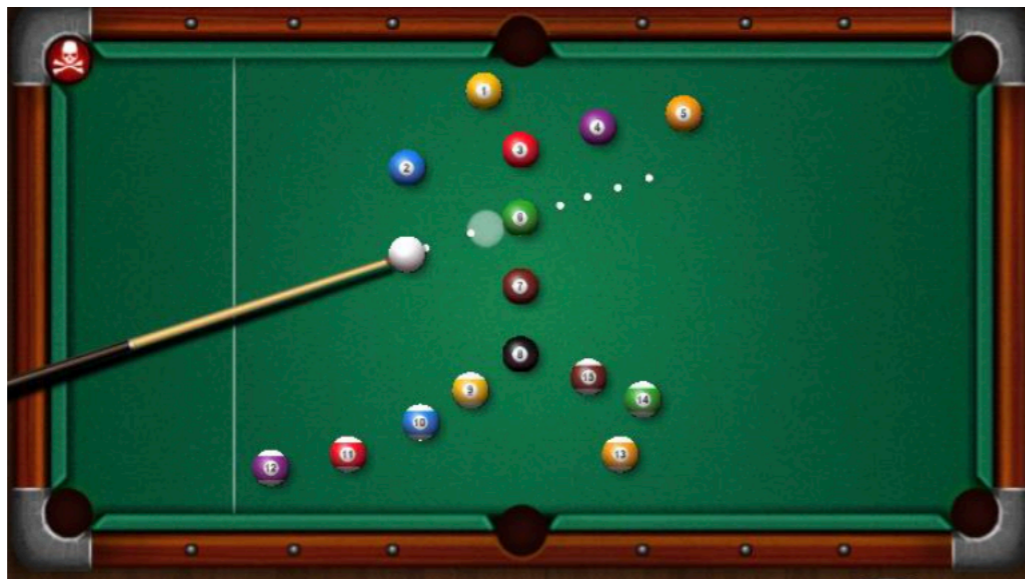


Allali, Hertzberg 2005.12287 (JCAP), 2012.12903 (PRD), 2103.15892 (PRL)

Dynamics can launch states into Schrodinger cat-like states



Schrodinger Cat Billiards

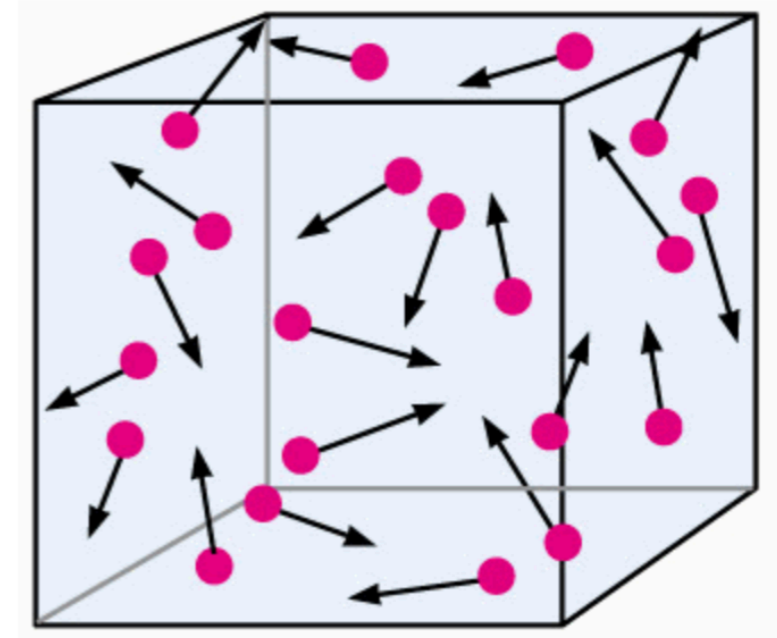
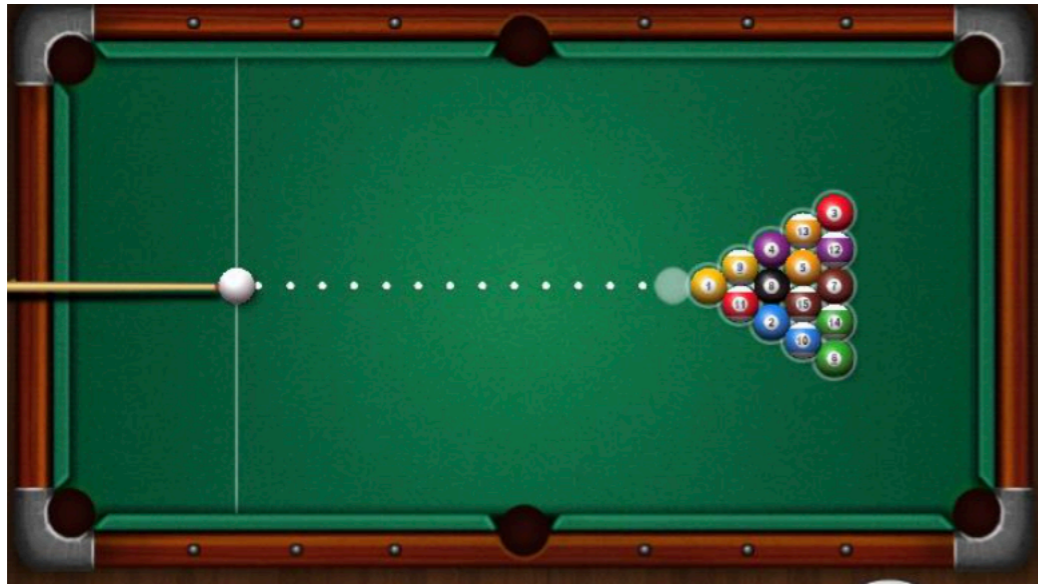


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Albrecht, Phillips 2012

Dynamics can launch states into Schrodinger cat-like states



Quantumness destroyed
due to **DECOHERENCE**

Schrodinger Cat Billiards

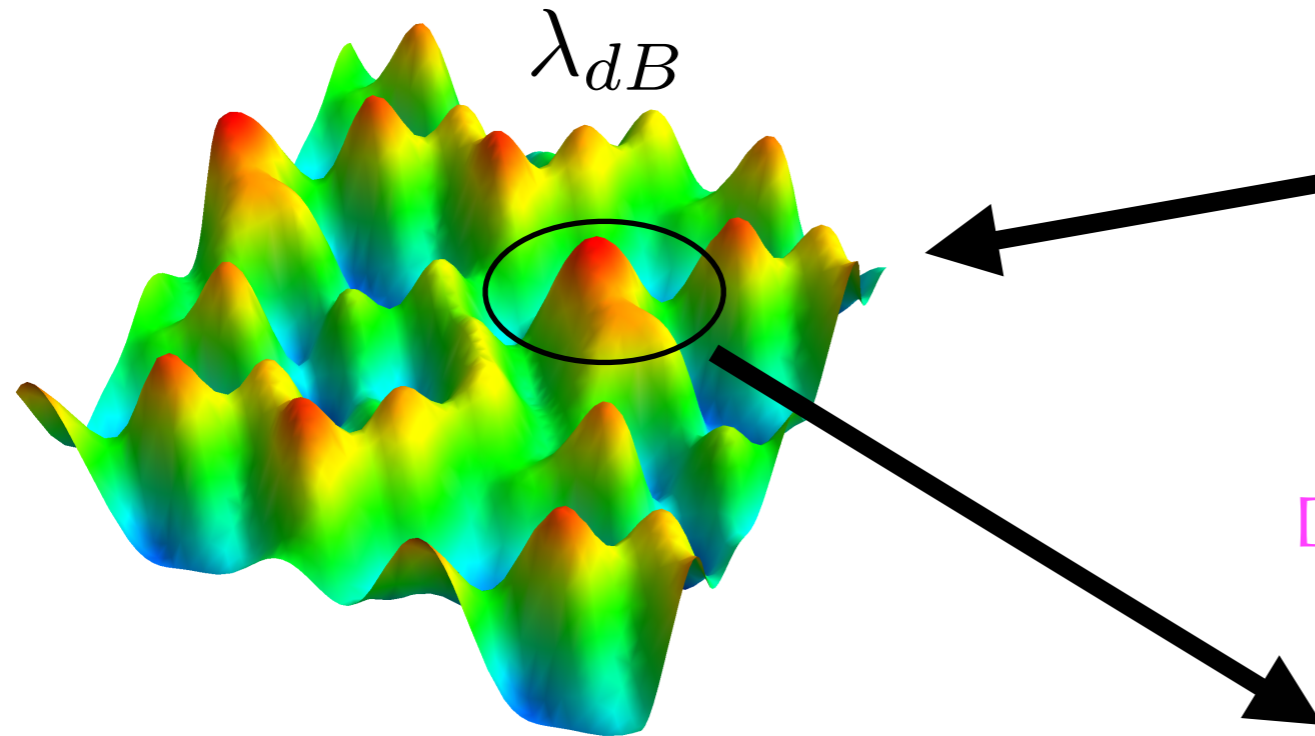
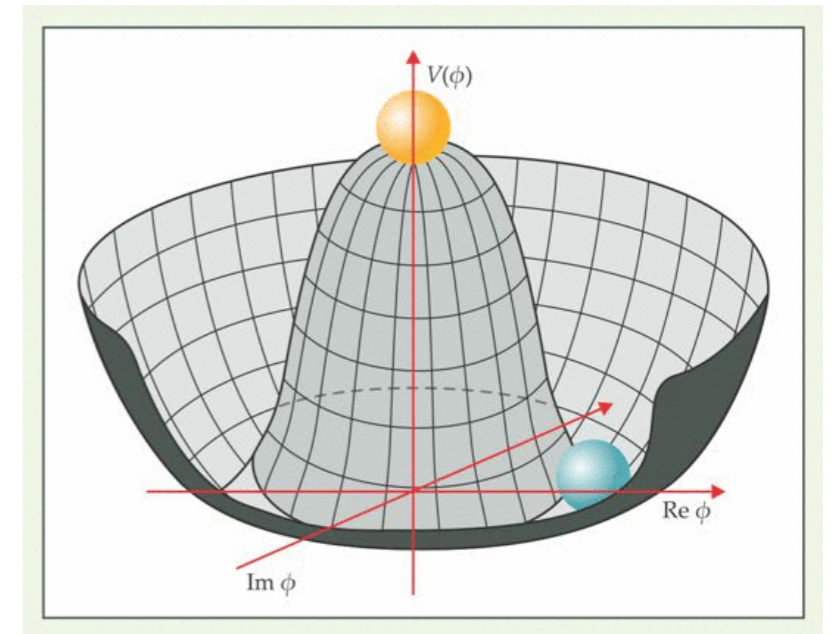


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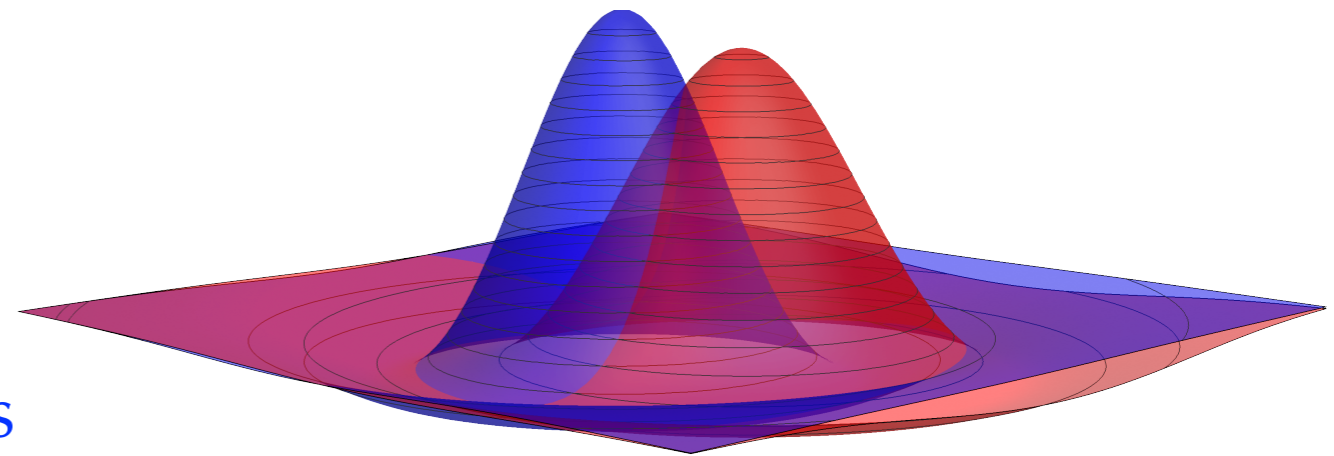
What about Dark Matter? Conceivably in Schrodinger cat-like states too

Claims: 0901.1106, 1111.1157, 1607.00949, 1710.02195, 1712.08219



Dark Matter Schrodinger Cat (Axions)

$$|DM_1\rangle + |DM_2\rangle$$



Quantumness destroyed
due to **DECOHERENCE???**

Less clear because dark matter has
tiny (non-gravitational) interactions

Claim it affect axion experiments: 2211.13602 Marsh

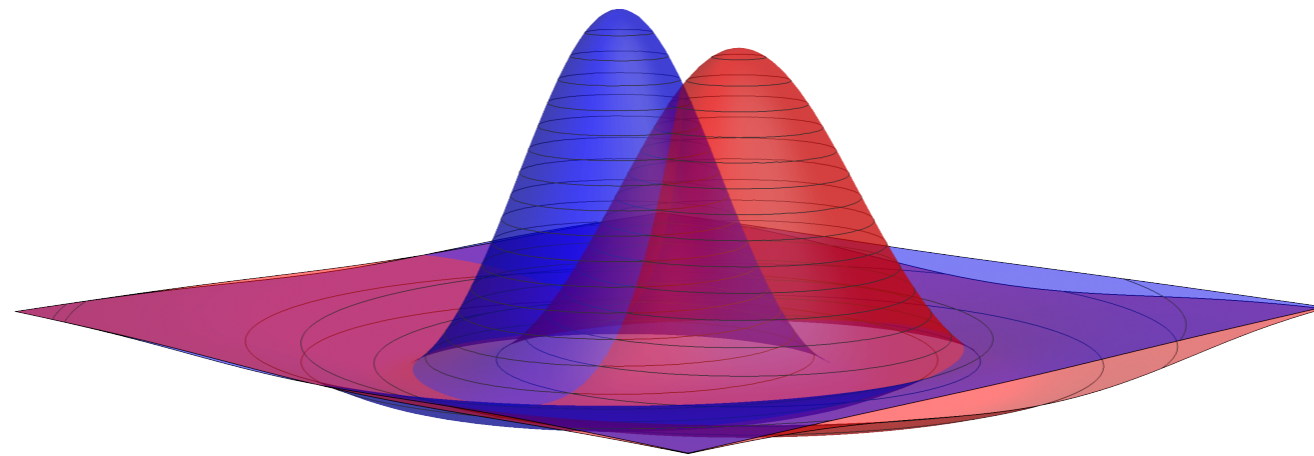
Environmental Entanglement from Gravitational Scattering

Probe particle $|\psi\rangle$



(e.g., baryon,...)

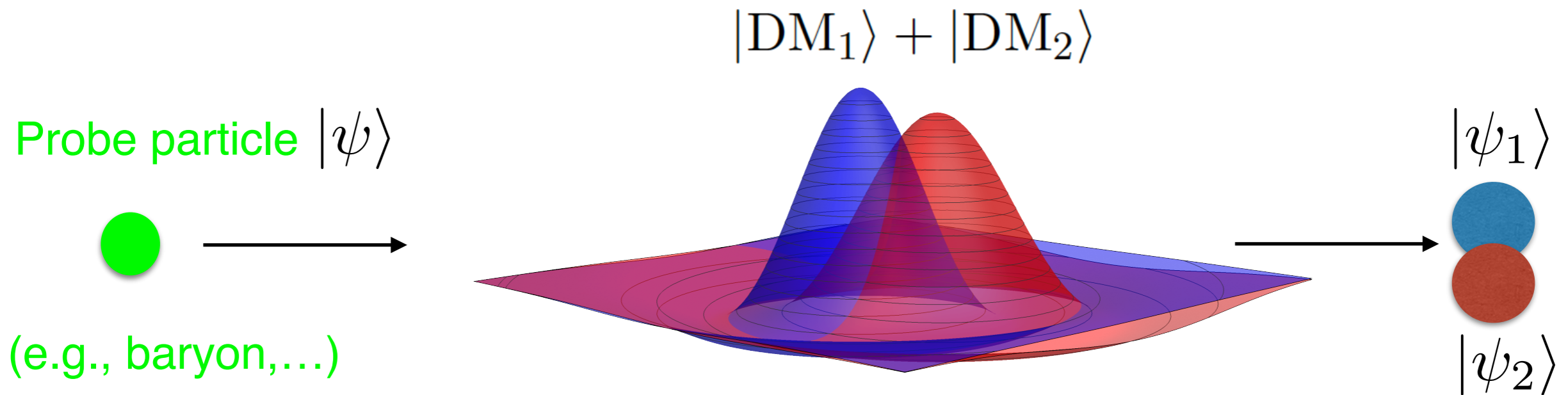
$|\text{DM}_1\rangle + |\text{DM}_2\rangle$



$$|\Psi_{\text{ini}}\rangle = (|\text{DM}_1\rangle + |\text{DM}_2\rangle) |\psi\rangle$$

Product State with probe particle

Environmental Entanglement from Gravitational Scattering



$$|\Psi_{\text{ini}}\rangle = (|\text{DM}_1\rangle + |\text{DM}_2\rangle) |\psi\rangle$$

Product State with probe particle

$$|\Psi_{\text{fin}}\rangle = |\text{DM}_1\rangle |\psi_1\rangle + |\text{DM}_2\rangle |\psi_2\rangle$$

Entangled State

Trace Out Environmental Probe Particles

$$\hat{\rho} \equiv |\Psi\rangle \langle \Psi|$$

Full Density Matrix

$$\hat{\rho}_{\text{red}} = \text{Tr}_{|\psi\rangle} [\hat{\rho}]$$

Reduced Density Matrix

$$= |\text{DM}_1\rangle \langle \text{DM}_1| + \langle \psi_2 | \psi_1 \rangle |\text{DM}_1\rangle \langle \text{DM}_2| + \langle \psi_1 | \psi_2 \rangle |\text{DM}_2\rangle \langle \text{DM}_1| + |\text{DM}_2\rangle \langle \text{DM}_2|$$

Off diagonal elements; controlling true quantum effects

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Off diagonal elements; controlling true quantum effects

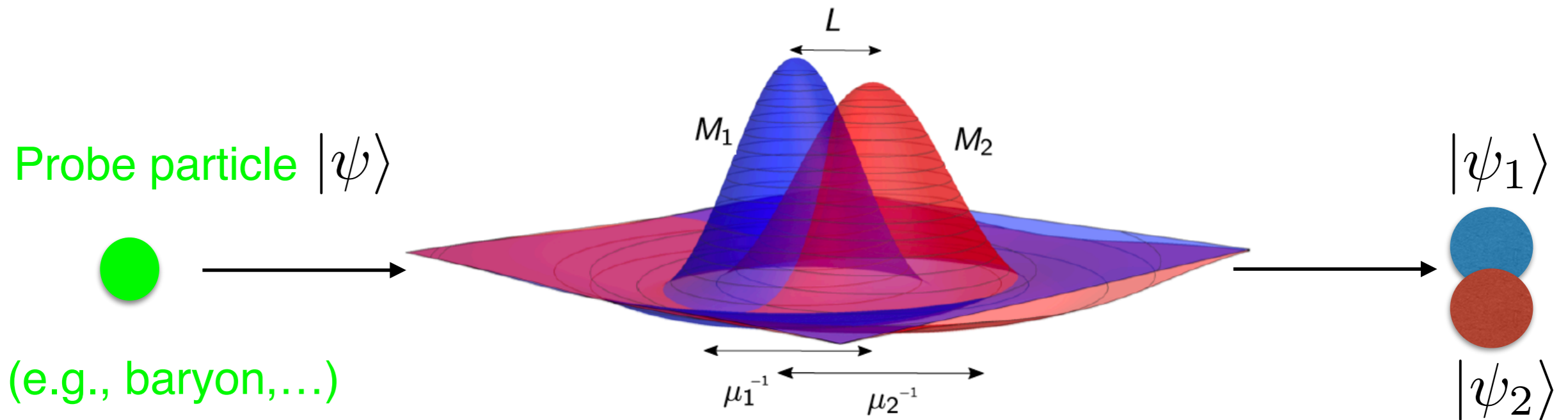
N-particles

$$\prod_{n=1}^N |\langle \psi_1 | \psi_2 \rangle|_n = \prod_{n=1}^N (1 - \Delta_b) \sim e^{-\sum_{n=1}^N \Delta_b}$$

Decoherence rate

$$\Gamma_{\text{dec}} = n v \int d^2b \Delta_b$$

Decoherence Rate from Generalized Cross-sections



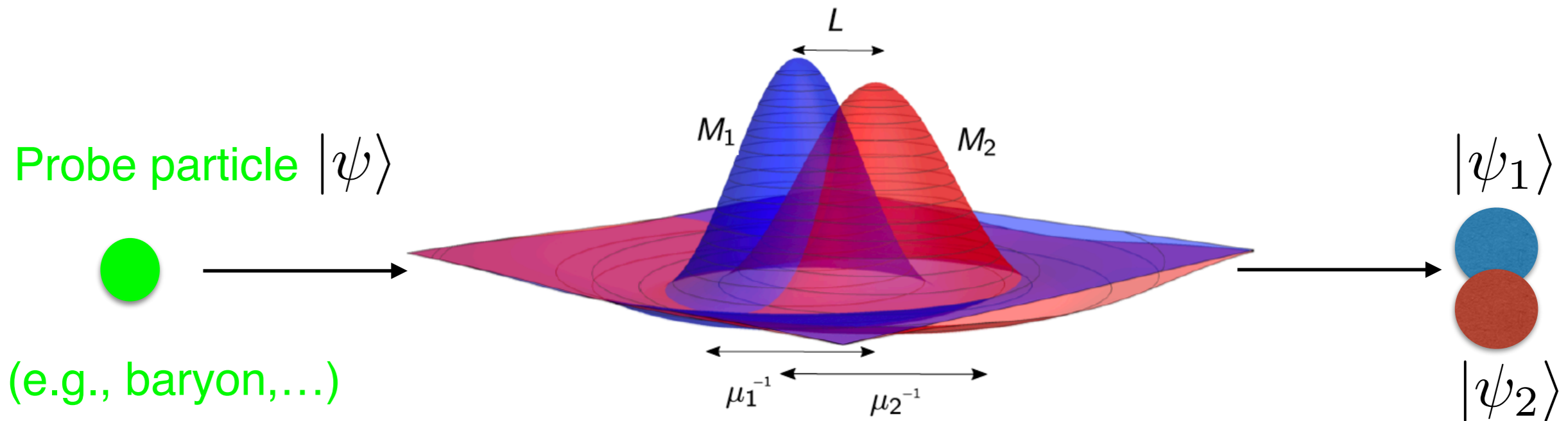
Scattering Amplitude

$$f(\vec{q}', \vec{q}) \equiv -\frac{1}{2\pi} \int d^3x' e^{i(\vec{q}-\vec{q}')\cdot\vec{x}'} \Phi(\vec{x}') m^2$$

Cross-section

$$\tilde{\sigma}_{ij}(q) \equiv \int d^2\Omega f_i^*(\vec{q}', \vec{q}) f_j(\vec{q}', \vec{q}) j_0(2qL_{ij} \sin \theta/2)$$

Decoherence Rate from Generalized Cross-sections



Decoherence Rate

$$\Gamma_{\text{dec}} \approx nv(\tilde{\sigma}_{1,1} + \tilde{\sigma}_{2,2} - 2\Re[\tilde{\sigma}_{1,2}])/2$$

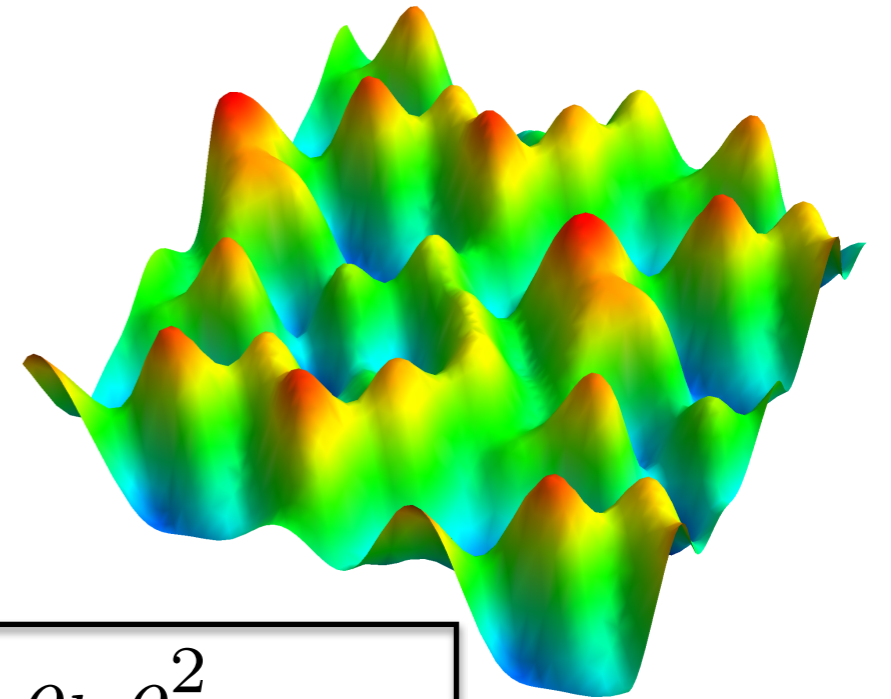
$$\Gamma_{\text{dec}} = \frac{4\pi G^2 m^4 n v}{k^2} \left[\frac{M_1^2}{\mu_1^2} \chi_{11} + \frac{M_2^2}{\mu_2^2} \chi_{22} - 2 \frac{M_1 M_2}{\mu_1 \mu_2} \chi_{12} \right]$$

Application to Light Diffuse scalar DM (axions)

Diffuse scalars (axions) $\mu_i \sim \frac{1}{\lambda_{dB,a}} \sim \frac{m_a v_{vir}}{2\pi}$ $M_i \sim \frac{4\pi}{3} \rho_{DM} \lambda_{dB,a}^3$

Probe: Diffuse baryons

$$k_p \sim m_p v_{vir}$$



Decoherence Rate

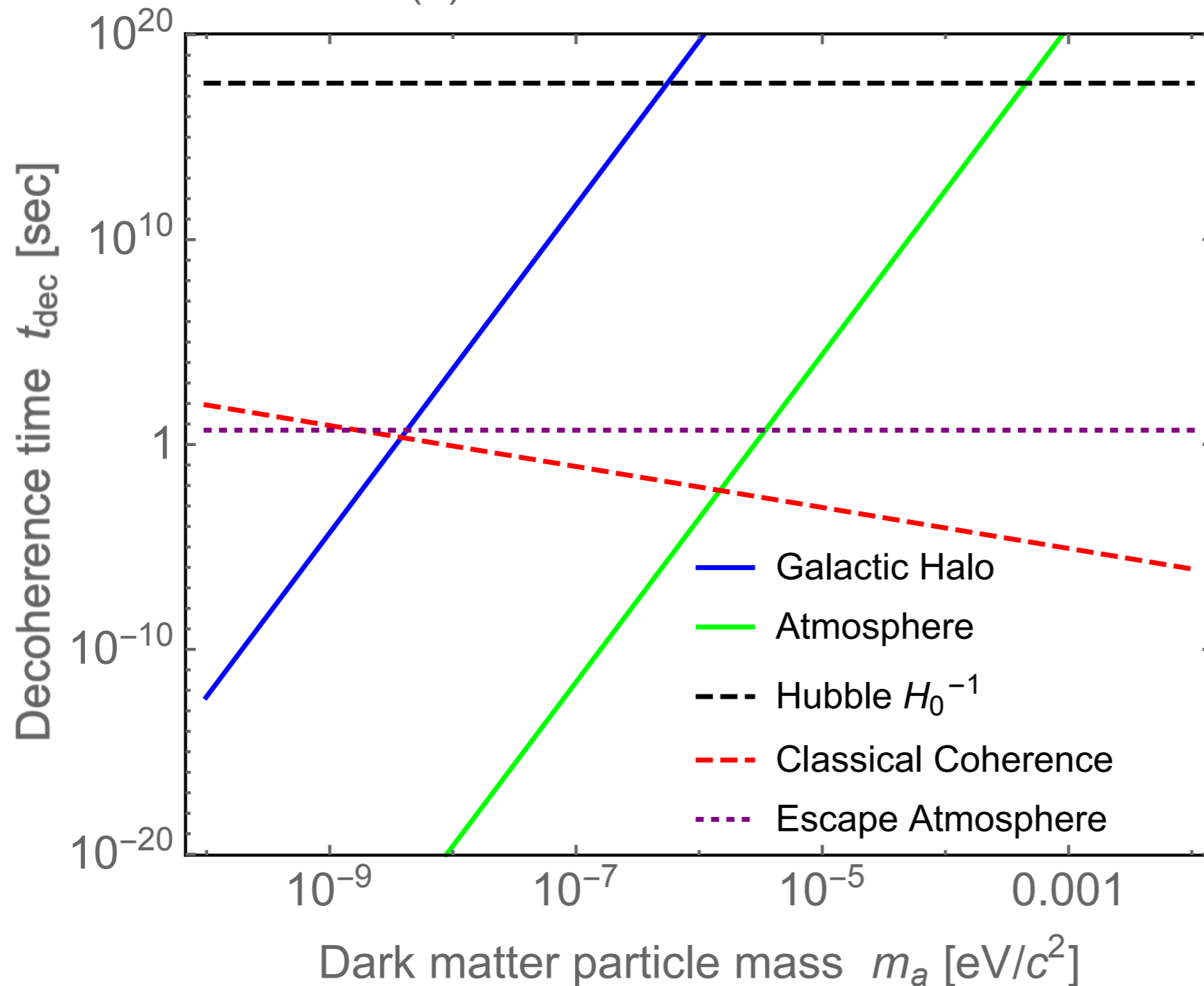
$$\Gamma_{dec} \sim \frac{G^2 m_p \rho_b \rho_{DM}^2}{m_a^8 v_{vir}^9}$$

Decoherence Time

$$t_{dec} = 1/\Gamma_{dec} \propto m_a^8$$

Application to Light Diffuse scalar DM (axions)

O(1) Mass and Size Parameters

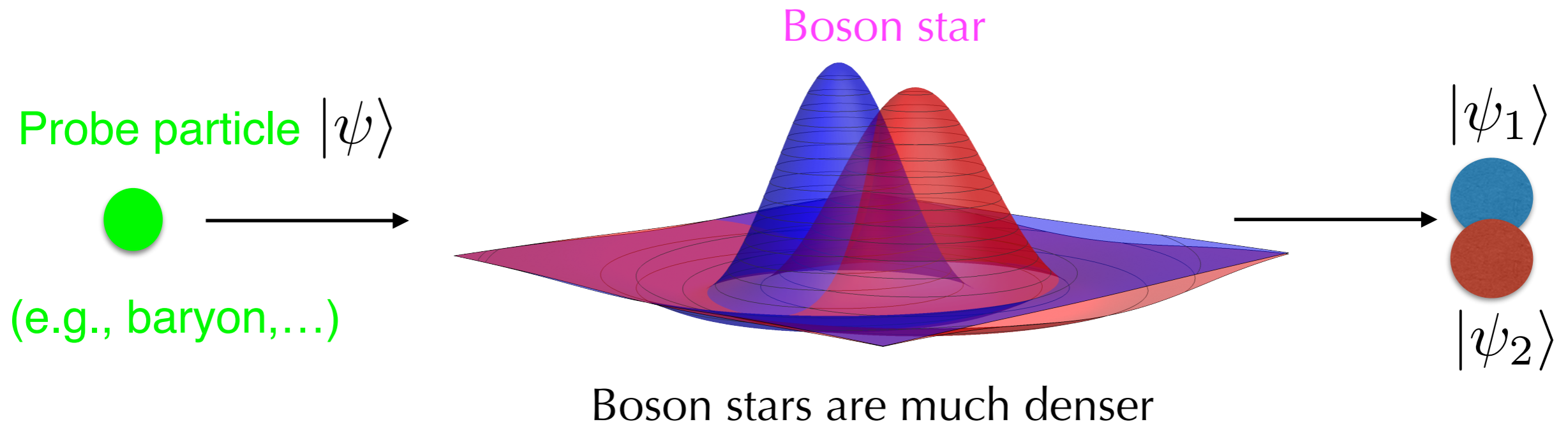


$$\Omega_{QCD,a} \sim 0.25 \left(\frac{10^{-5} \text{ eV}}{m_a} \right)^{1.2}$$

$$\lambda_{dB,a} \sim \text{km} \left(\frac{10^{-5} \text{ eV}}{m_a} \right)$$

$$M \sim \text{ng} \left(\frac{10^{-5} \text{ eV}}{m_a} \right)^3$$

Application to Boson Stars

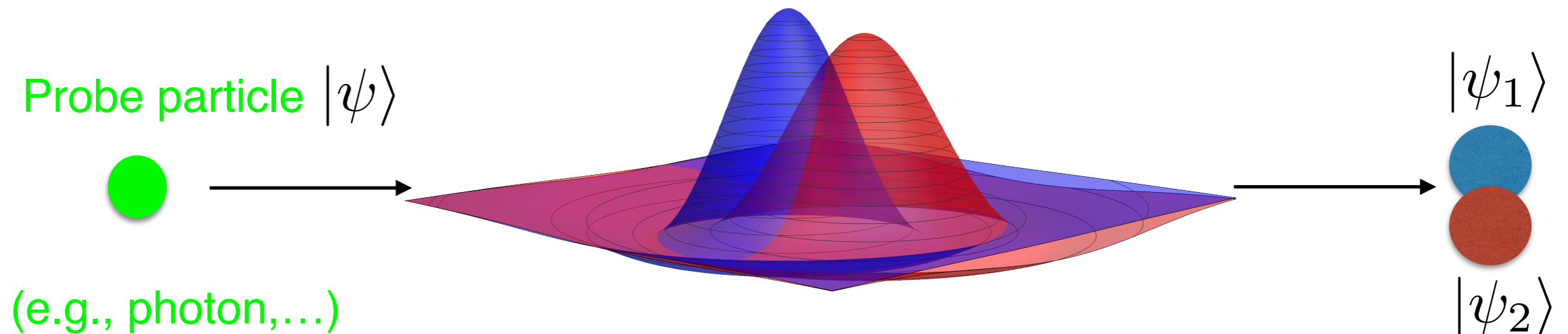


Decoherence Rate

$$\Gamma_{\text{dec}} \gtrsim \frac{\hbar^2 m_p \rho_p}{v_p m_a^4} \sim 10^{21} \text{ sec}^{-1} \left(\frac{1 \text{ eV}}{m_a c^2} \right)^4$$

Extremely rapid decoherence \rightarrow Very classical

General Relativistic Extension



Robust quantum gravity calculation;
General Relativity treated as quantum effective theory

Decoherence Rate for Static Source

Metric - Newton gauge $g_{\mu\nu} = \text{diag}[(1 + 2\Phi), -(1 - 2\Psi), -(1 - 2\Psi), -(1 - 2\Psi)]$

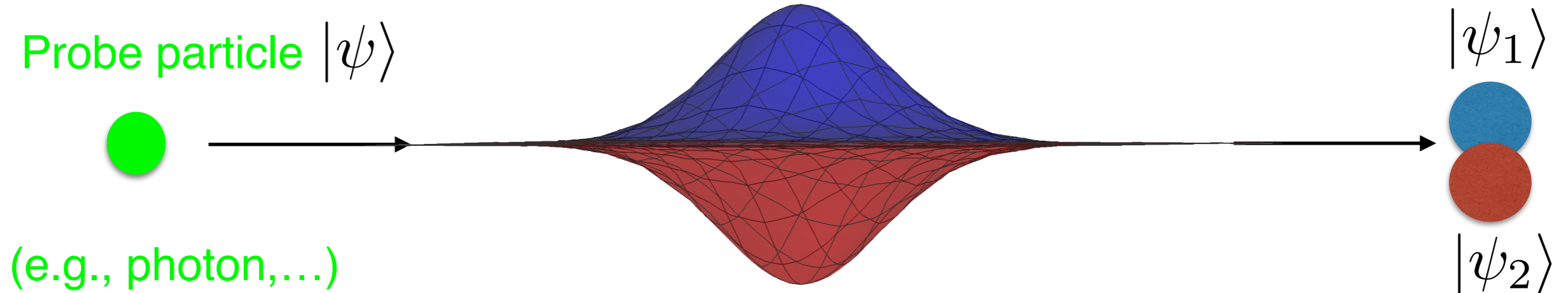
Amplitude $f(\vec{q}', \vec{q}) \equiv -\frac{1}{2\pi} \int d^3x' e^{i(\vec{q}-\vec{q}')\cdot\vec{x}'} \left[\Phi(\vec{x}') E_q^2 + \Psi(\vec{x}') q^2 \right]$

Rate

$$\Gamma_{\text{dec}} \approx 4\pi G_N^2 n_p v_p \frac{(m_p^2 + 2k^2)^2}{k^2} \left(\frac{M_1^2}{\mu_1^2} \chi_{11} + \frac{M_2^2}{\mu_2^2} \chi_{22} - \frac{2M_1 M_2}{\mu_1 \mu_2} \chi_{12} \right)$$

In galaxy, the baryons provide a bigger environment than photons/neutrinos

Decoherence Rate for Oscillating Source



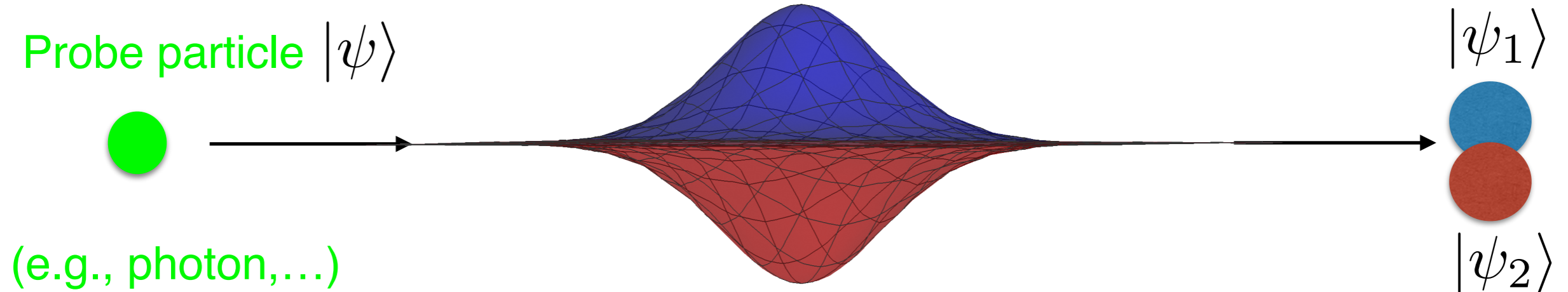
To “learn” about phase; inelastic scattering

Kinematic **mismatch**

$$E_p \rightarrow E_p \pm \omega_a \qquad \delta p_T \sim \frac{E_p}{p_p} \omega_a \gg p_a$$

Decoherence of phase is suppressed (unless all states are relativistic)

Decoherence Rate for Oscillating Source



Conceivably, relevant to direct detection,
which is sensitive to axion phase

$$|\mathbf{DM}\rangle \sim \sum_i c_i |\cos(\omega t - \mathbf{k}_a \cdot \mathbf{x} + \varphi_i)\rangle$$

Open question: are there observables related to this?

Conclusions

- Macroscopic quantum states (Schrodinger cats) of light scalar dark matter might exist, and could potentially have slow decoherence
- We studied the decoherence of such states due to gravitational scattering from probe particles; a robust quantum gravity calculation
- We found that superpositions of spatial profiles decohere rapidly for very light DM (axions), and boson stars decohere extremely rapidly
- We found that superpositions of phases live much longer, may launch detectors into superpositions (non-grav interactions can be considered in future work).
- Relativistic states (near black holes) decohere quickly

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