

QUANTUM COMPUTATIONS IN LOOP QUANTUM GRAVITY

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arXiv:2304.03559

arXiv:2003.13124 (with Jakub Mielczarek)



Quantum Cosmos Lab



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IN KRAKÓW

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QUANTUM COMPUTERS

- qubits \equiv 2 level quantum systems

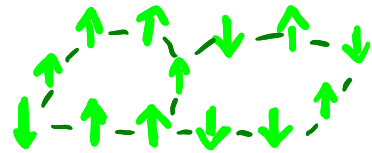
$$\alpha |\uparrow\rangle + \beta |\downarrow\rangle$$

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$$\alpha |\uparrow\rangle + \beta |\downarrow\rangle$$

- Hilbert space dimension = $2^{\text{\# qubits}}$

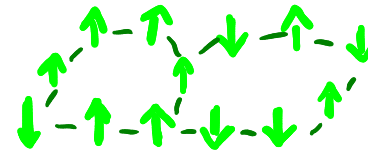


QUANTUM COMPUTERS

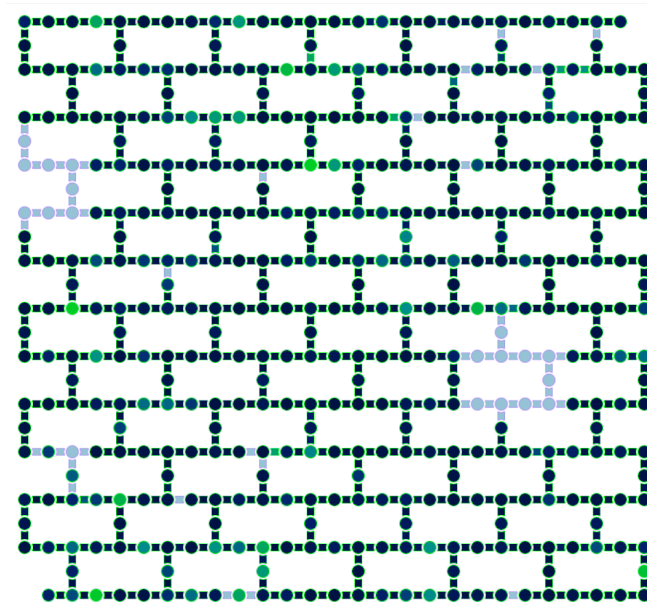
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- e.g.: IBM SEATTLE, 433 qubits



QUANTUM COMPUTERS

- exponentially hard to simulate on a classical computer

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- there are known quantum algorithms that are faster than classical
- e.g.: Shor's algorithm factorizes integer N with complexity
 $O((\log N)^2 (\log \log N))$ vs fastest classical algorithm
 $O(e^{1.9 (\log N)^{2/3} (\log \log N)^{2/3}})$

QUANTUM COMPUTERS

PROBLEMS:

- quantum computations errors
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SOLUTIONS:

- error correction codes
- each logical qubit is encoded in many physical qubits
- similarly logical gates

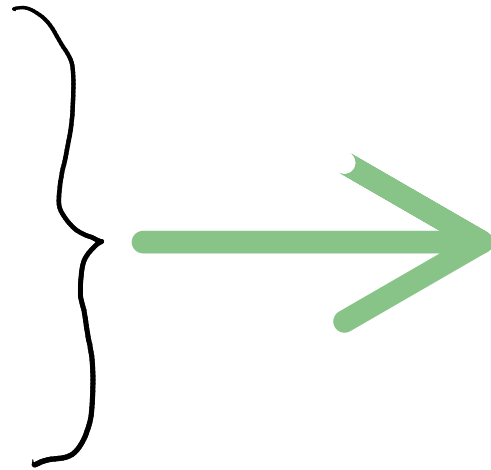
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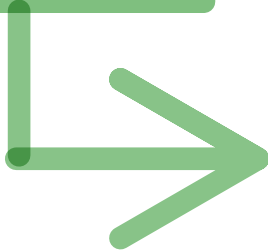
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fault tolerant
quantum computers

QUANTUM COMPUTERS

NISQ era

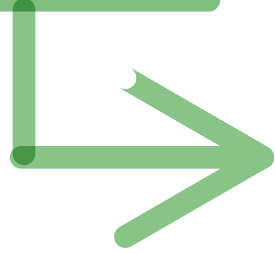


≡ noisy intermediate-scale quantum era

- short circuits
- variational algorithms

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PROBLEM:

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NISQ era

≡ noisy intermediate-scale quantum era

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PROBLEM:

- transition between quantum \leftrightarrow classical data

SOLUTION:

- work as long as possible
only on quantum data

simulations of quantum
systems in physics

LOOP QUANTUM GRAVITY

- background independent, non-perturbative

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- background independent, non-perturbative
- canonical quantization of gravity

$$S = \frac{1}{16\pi G} \int d^3x \sqrt{-g} R \longrightarrow S = \frac{1}{16\pi G} \int d^3x (N^a H_a + NH)$$

LOOP QUANTUM GRAVITY

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- Ashtekar variables

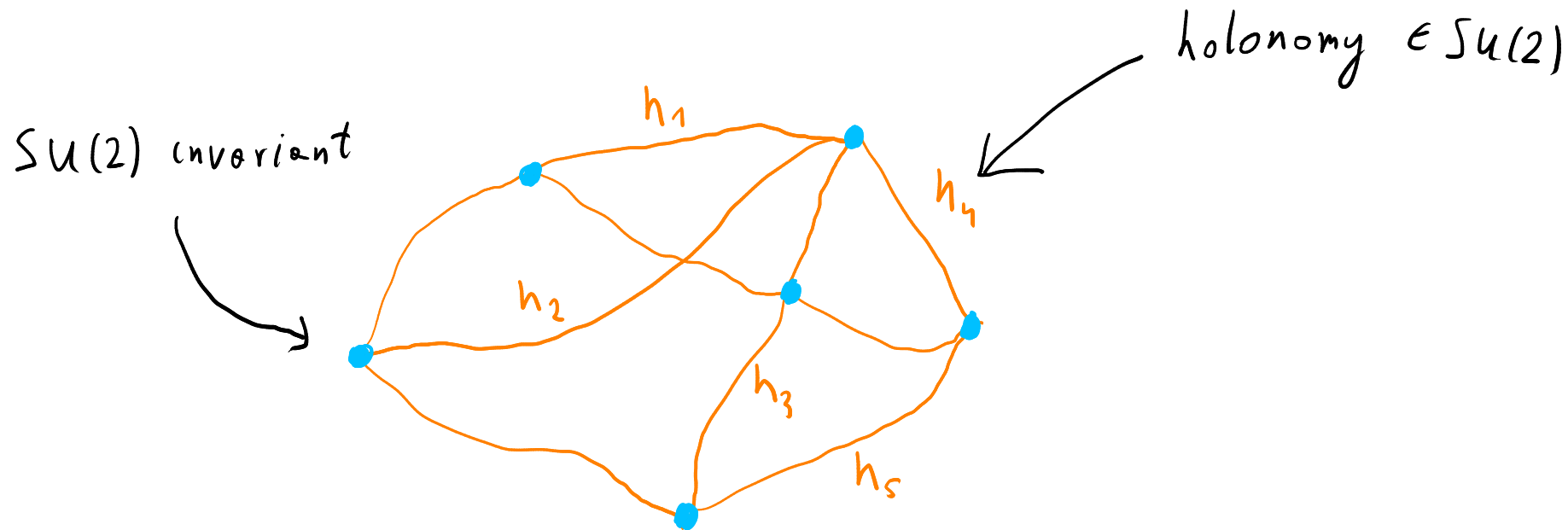
$$\begin{cases} E_i^a = \frac{1}{2} \epsilon_{ijk} \epsilon^{abc} e_b^j e_c^k \\ A_a^i = \gamma \omega_a^{0i} + \frac{1}{2} \epsilon_{ijk} \omega_a^{jk} \end{cases}$$

$$\{A_a^i(x), E_j^b(y)\} = 8\pi G \gamma \delta_a^b \delta_j^i \delta(x-y)$$

- gravity as $SU(2)$ gauge theory

SPIN NETWORKS

- in LQG represent quantum geometry of space



$$\Psi_\Gamma: SU(2)^{\times E} \rightarrow \mathbb{C}$$

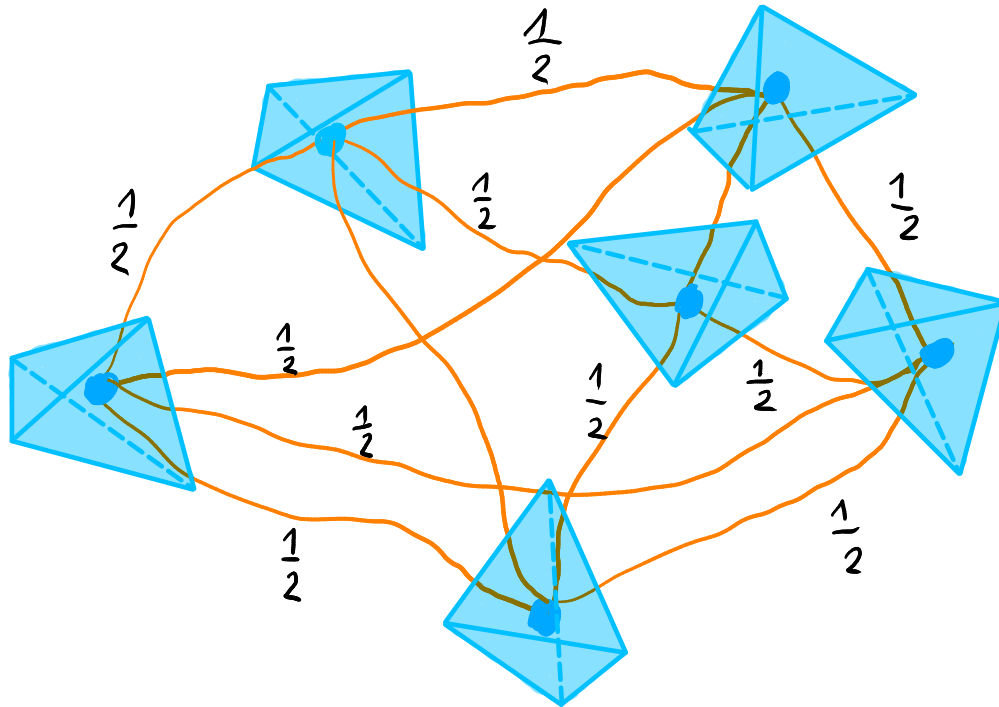
$$\{g_e\}_{e \in \Gamma} \longmapsto \Psi_\Gamma(\{g_e\}_{e \in \Gamma}) = \Psi_\Gamma(\{h_{s(e)} g_e h_{t(e)}^{-1}\})$$

SPIN NETWORKS

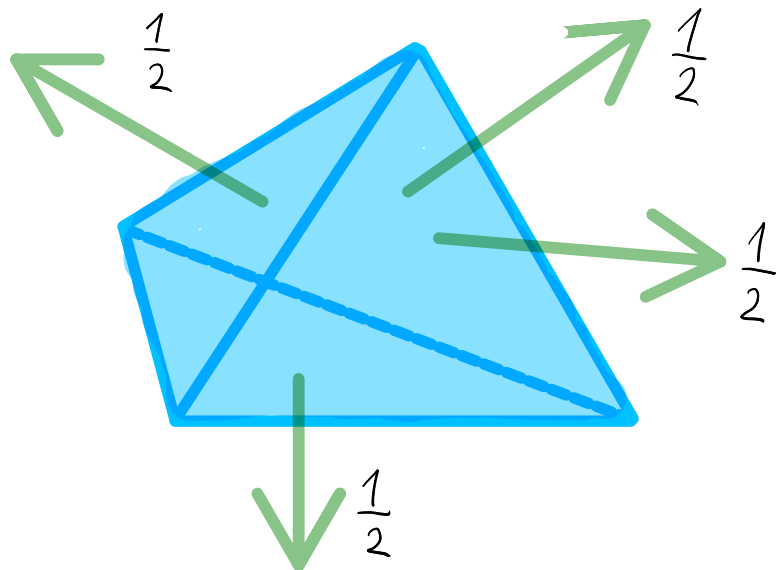
- nodes are associated with quanta of volume
- links are associated with areas around given volume
- set of all spin networks (up to diffeomorphisms) constitutes basis of kinematical Hilbert space
- spin networks can be also used for general gauge theories with a compact Lie group and a connection form

SPIN NETWORKS IN OUR CASE

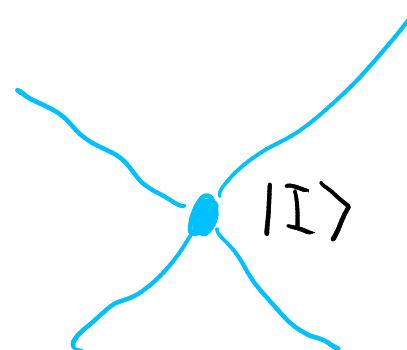
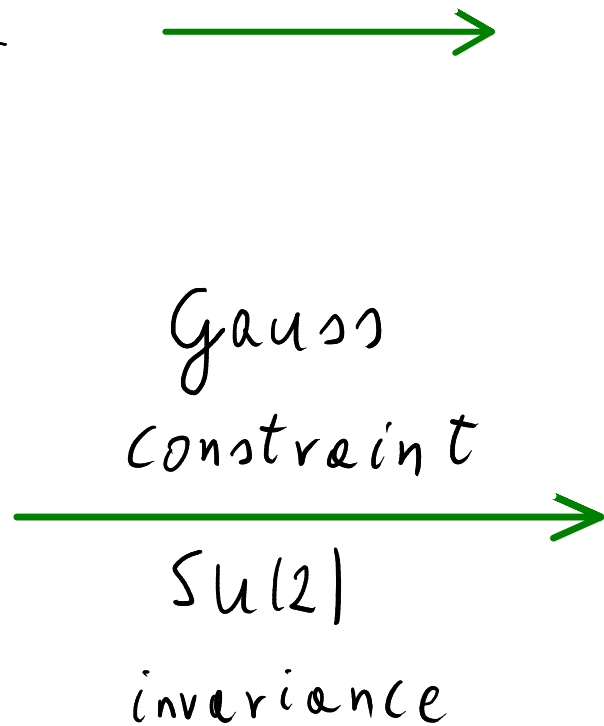
- 4-valent
- all holonomies in fundamental $SU(2)$ representation



SPIN NETWORKS IN OUR CASE

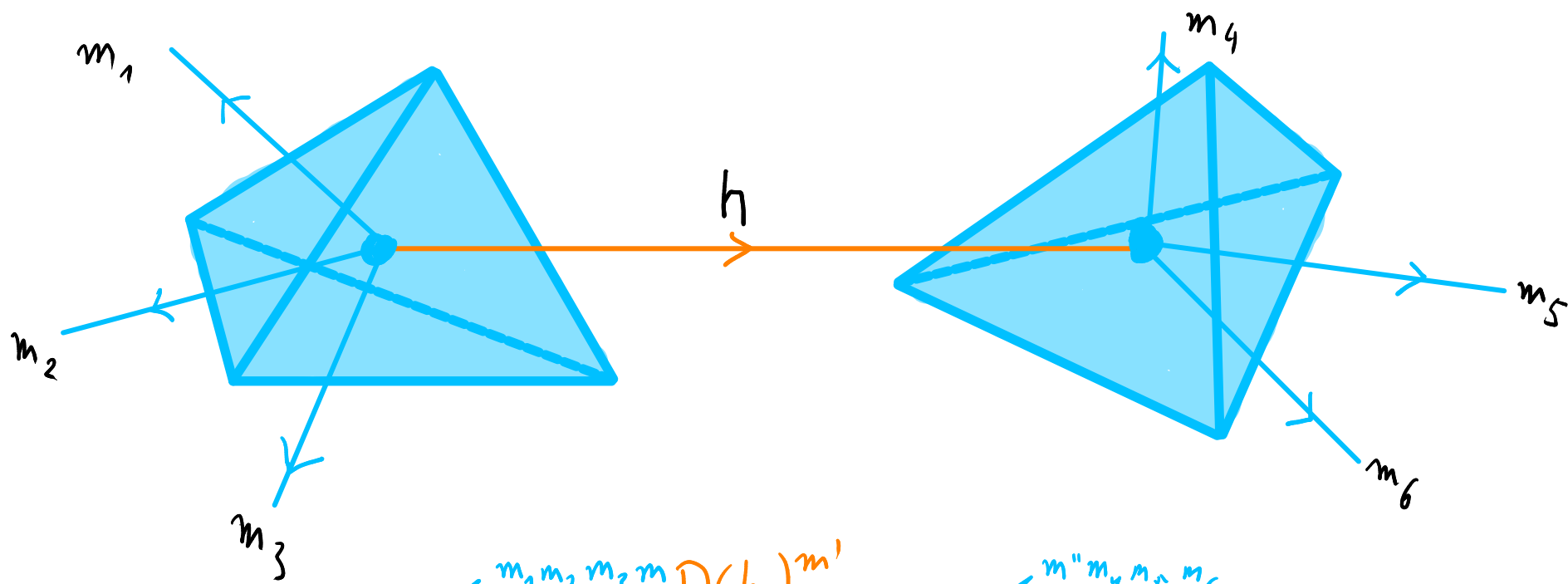


$$\mathcal{H}_{\frac{1}{2}} \otimes \mathcal{H}_{\frac{1}{2}} \otimes \mathcal{H}_{\frac{1}{2}} \otimes \mathcal{H}_{\frac{1}{2}}$$



$$|I\rangle = \alpha|0\rangle + \beta|1\rangle$$

STATE OF SPIN NETWORK



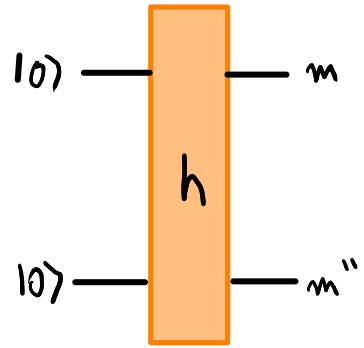
$$\begin{array}{c}
 \leftarrow m_1 m_2 m_3 m \\
 D(h)^{m'} \\
 m \rho_{m' m''} \leftarrow m'' m_4 m_5 m_6
 \end{array}$$

$$\rho_{m m'} = \delta_{m, -m'} (-1)^{j-m}$$

Wigner 4j-symbol

Wigner matrix

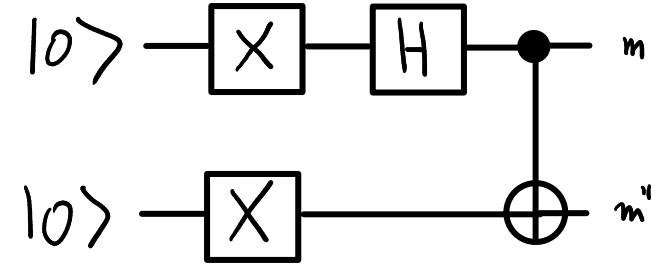
LINKS AND GAUSS CONSTRAINT AS QUANTUM CIRCUITS



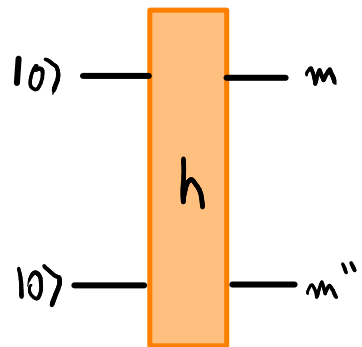
e.g.: $h=11$

$$D(h)_{m'}^{m''} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\leadsto \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle)$$



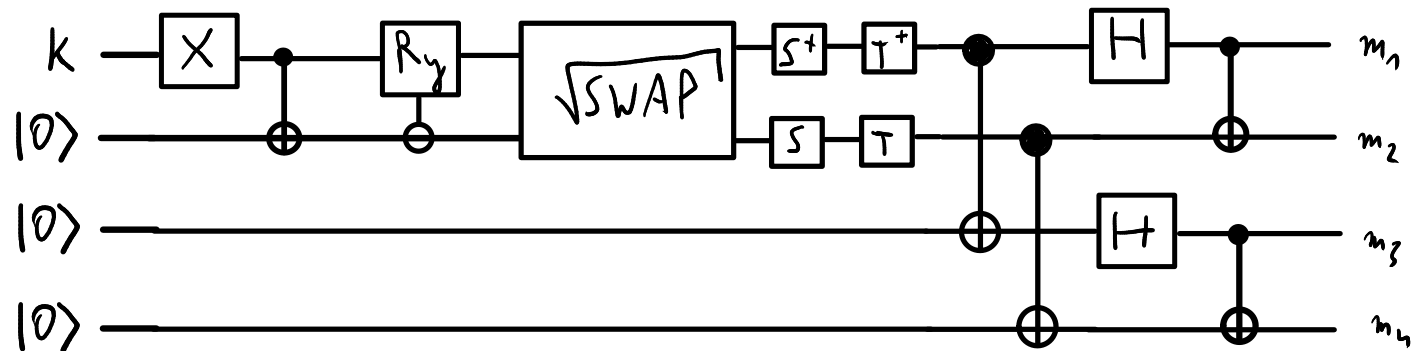
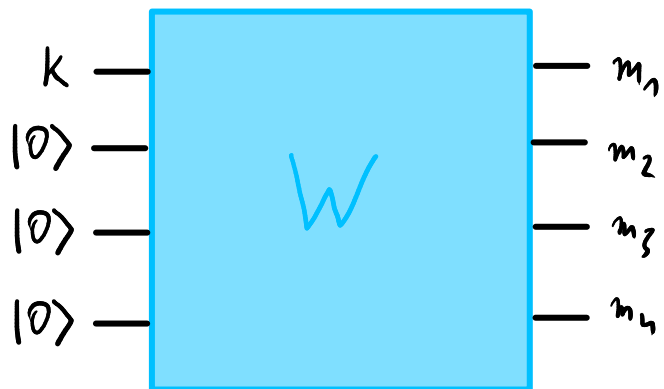
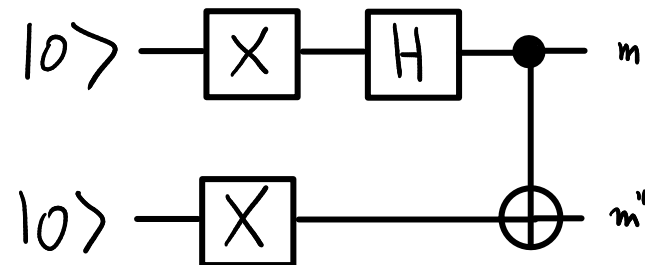
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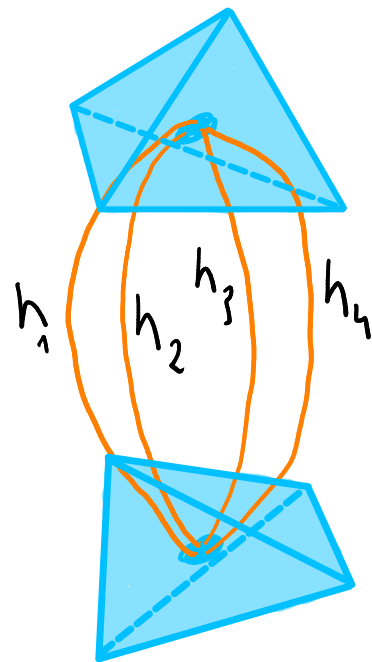
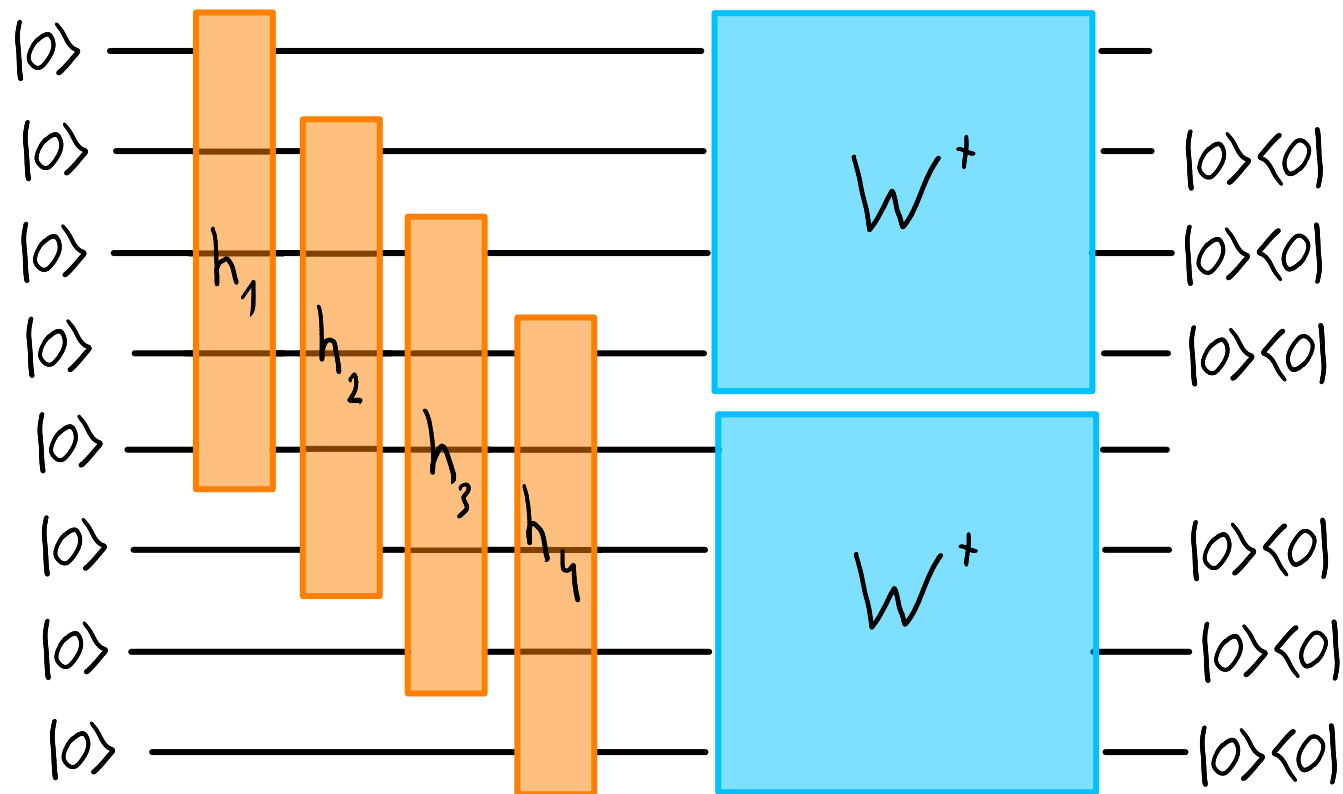
$$D(h)_{m'}^m f_{m''}^{m'} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\sim \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle)$$

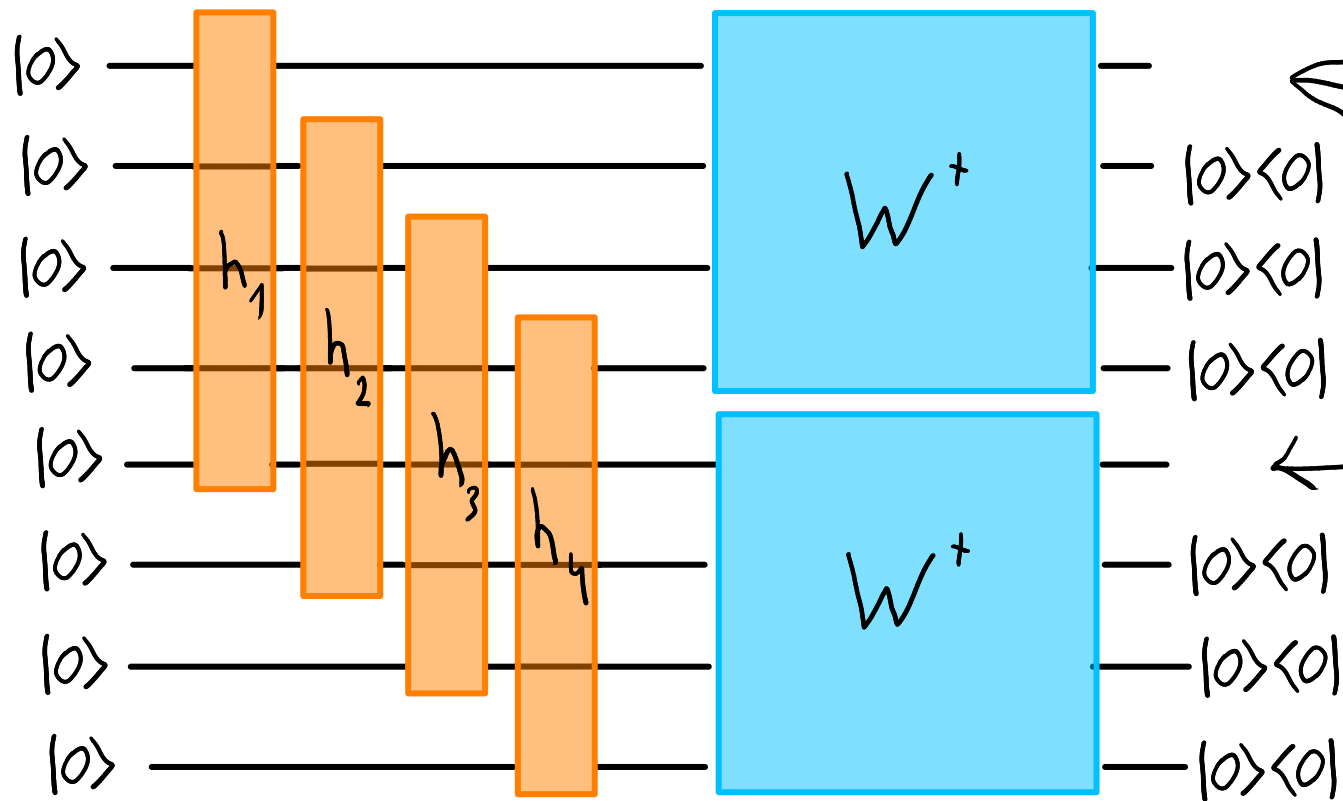


$$W|k\rangle|0000\rangle = \sum_{m_i} c_{(k)}^{m_1 m_2 m_3 m_4} |m_1 m_2 m_3 m_4\rangle$$

1ST METHOD: PROJECTION



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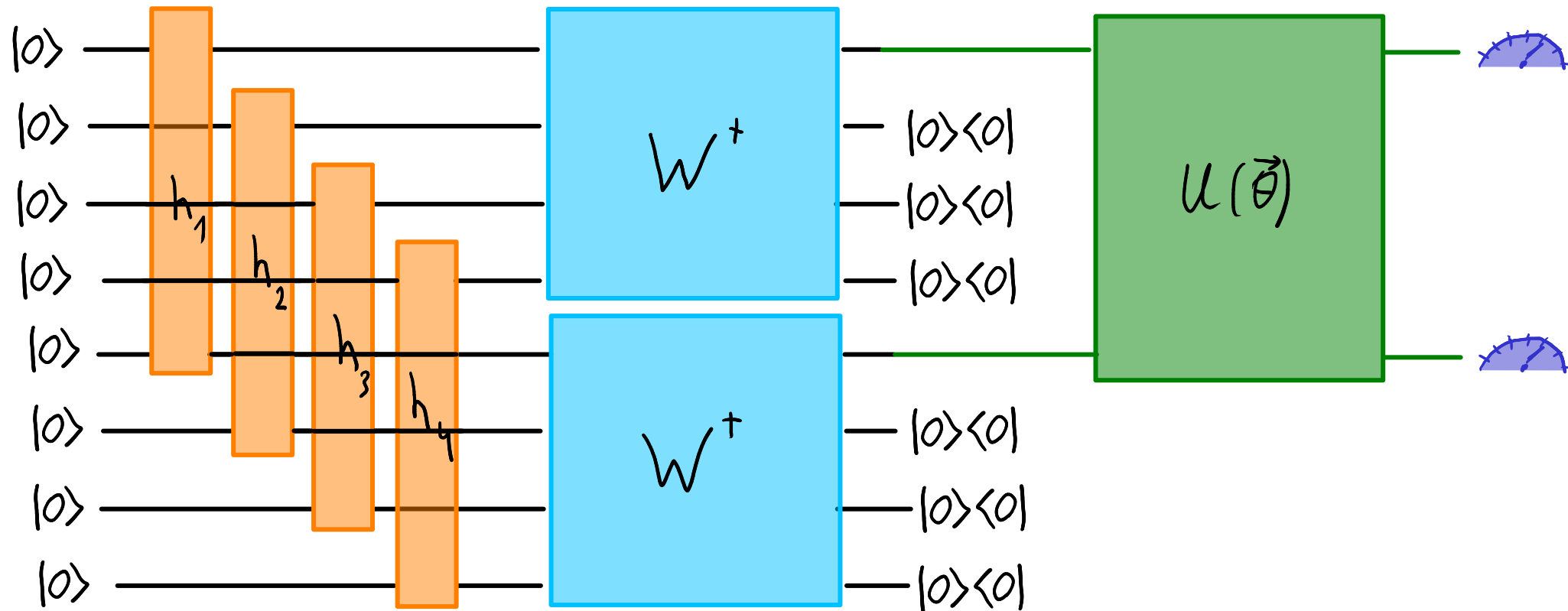


QUANTUM
TOMOGRAPHY

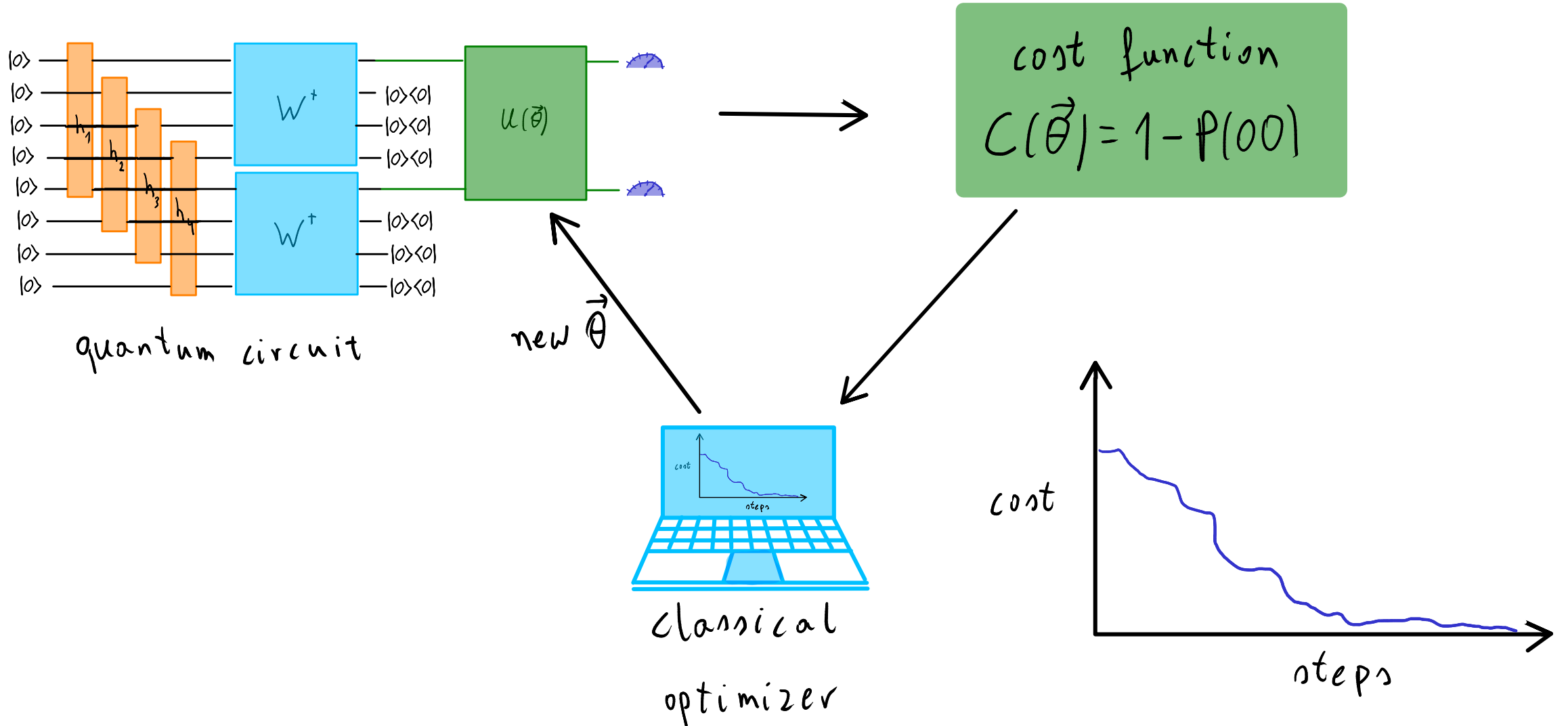


$\sim \left. \begin{array}{l} \# \text{ qubits} \\ \text{exp. val. measurements} \end{array} \right\}$

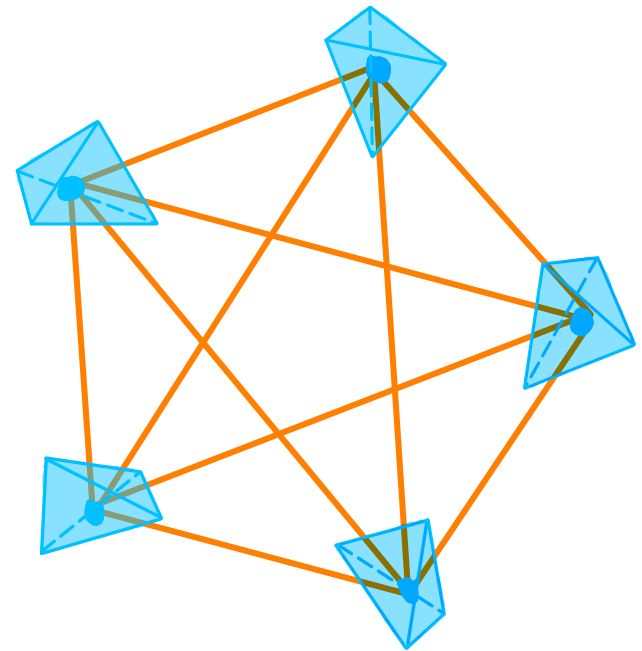
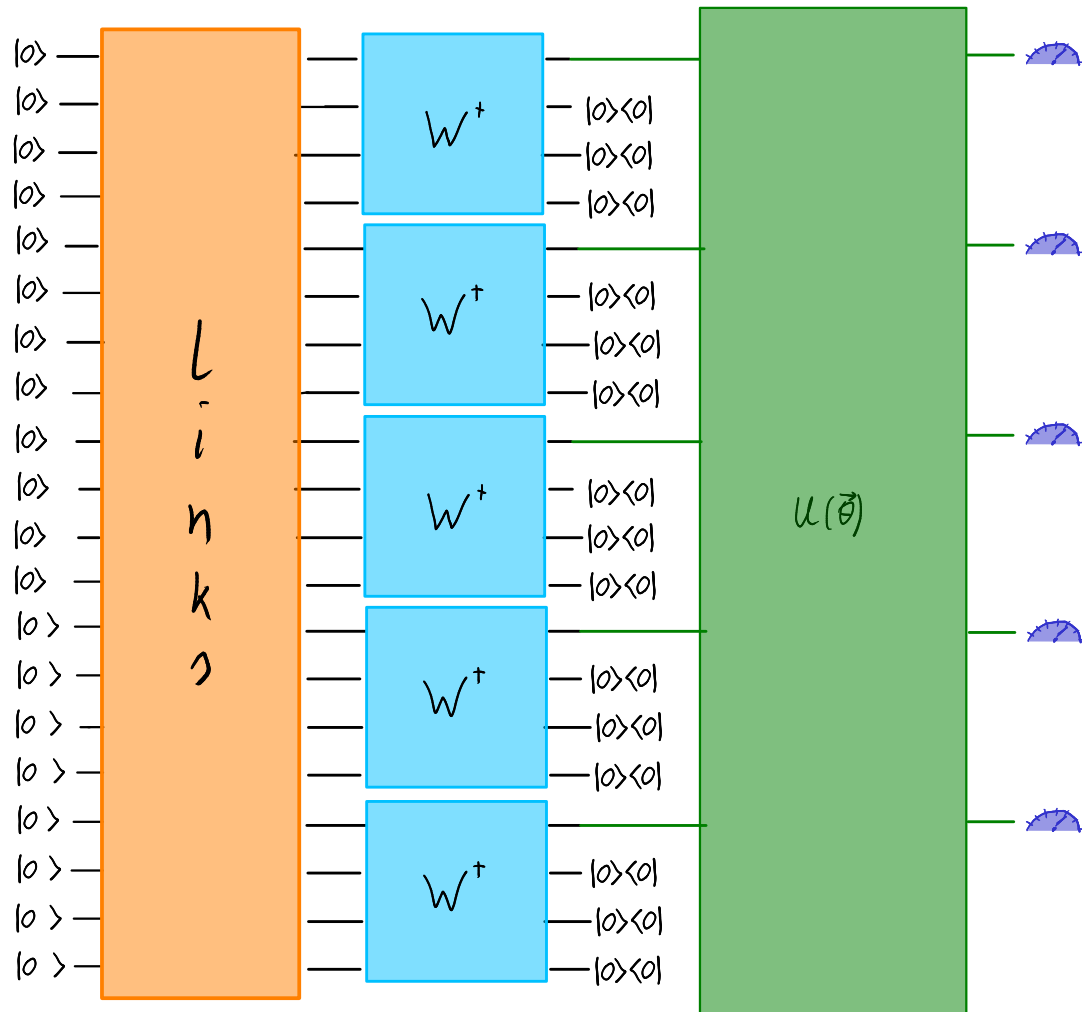
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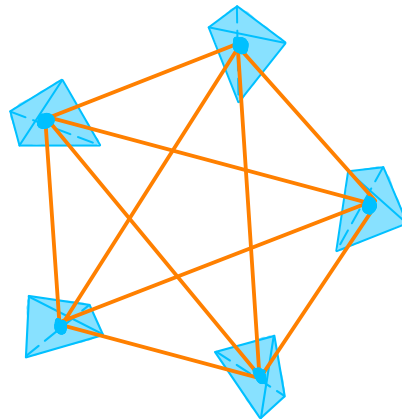
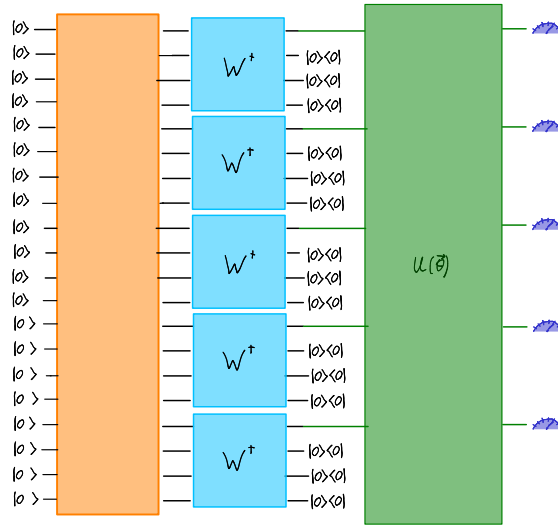


1ST METHOD: PROJECTION

qubits

=

4 · # nodes

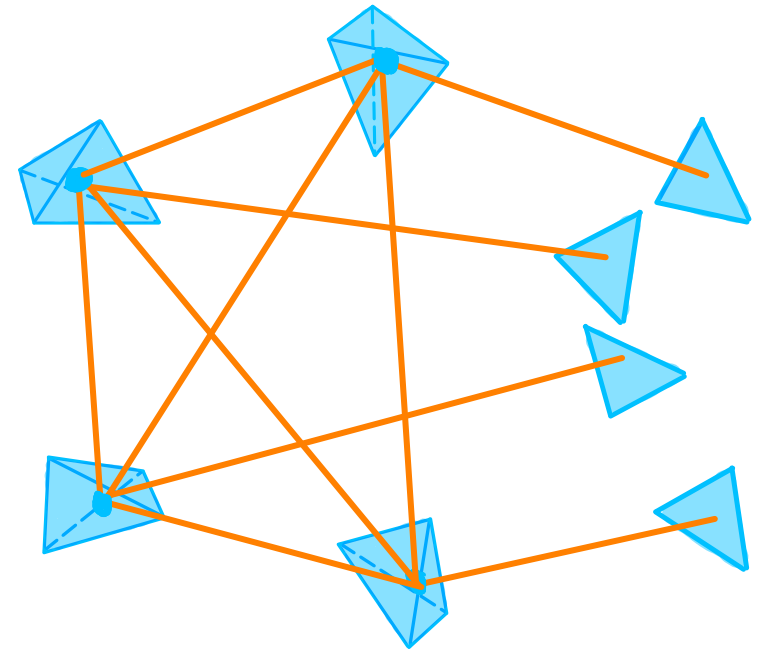
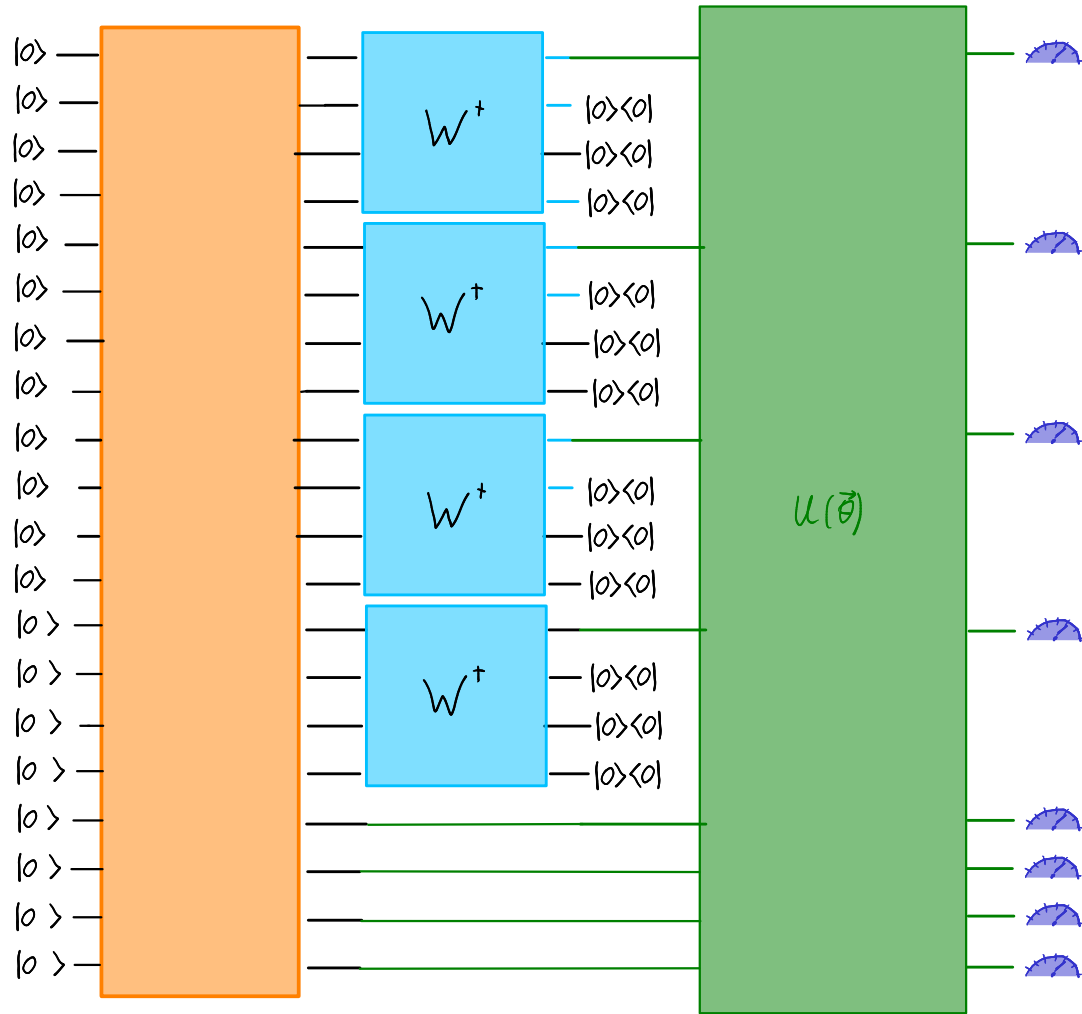


projection on $|0\rangle\langle 0|$ ^{⊗3}
by postselection of outcomes

Exponentially many
results rejected

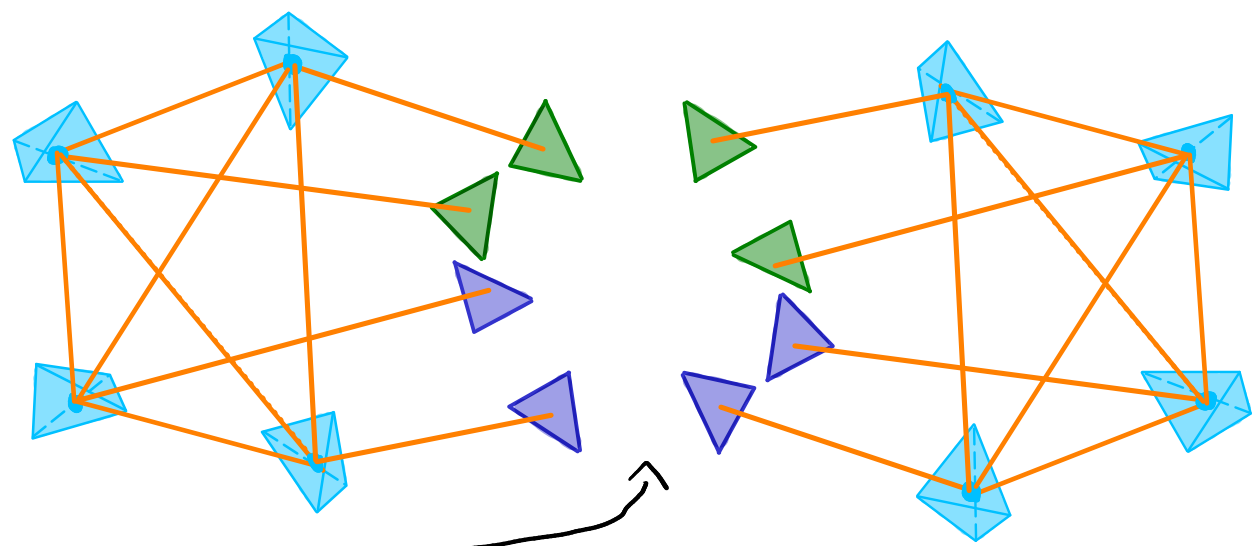
exponentially many
measurements required
to estimate cost function

1ST METHOD: PROJECTION, PARTIAL PROJECTION



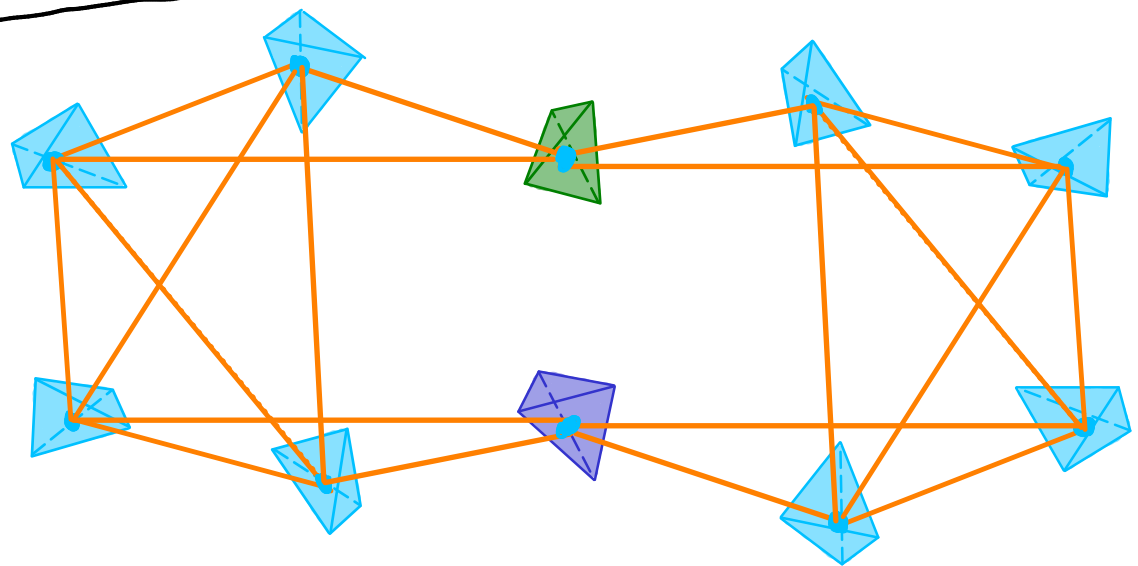
1ST METHOD: PROJECTION. GLUING.

partially
projected
network



gluing by applying

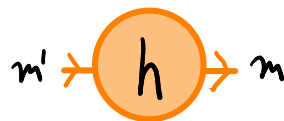
W^+	-
	$ 0\rangle\langle 0 $
	$ 0\rangle\langle 0 $
	$ 0\rangle\langle 0 $



2ND METHOD: TENSOR NETWORK



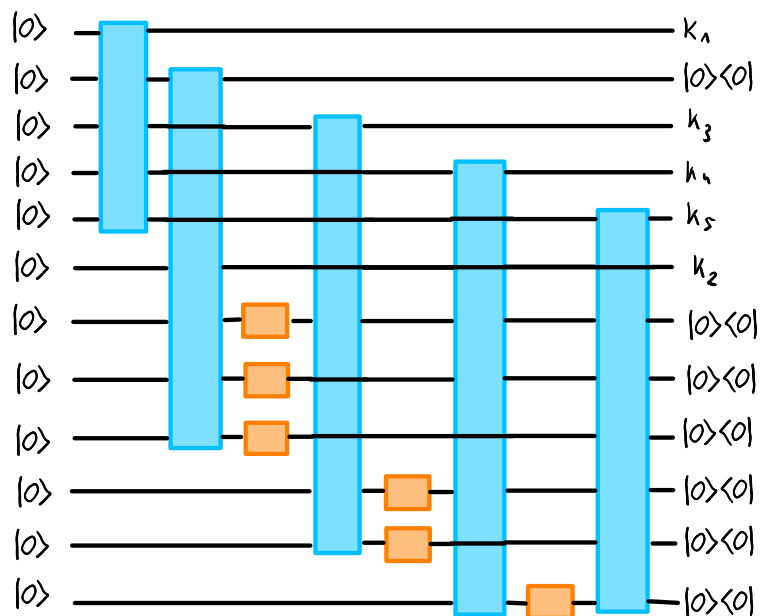
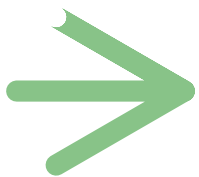
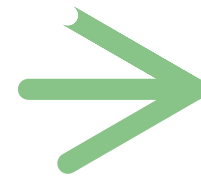
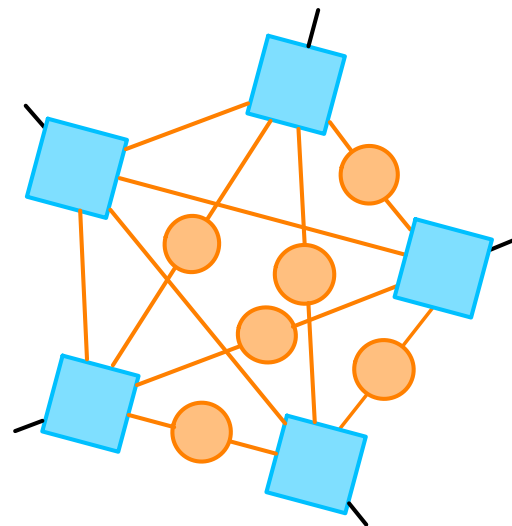
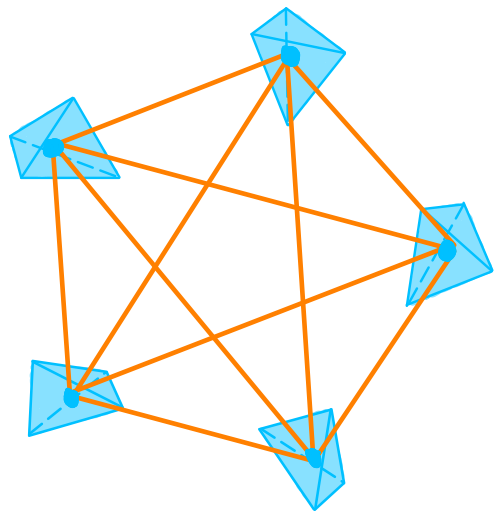
$$D(h)_{m'}^m$$



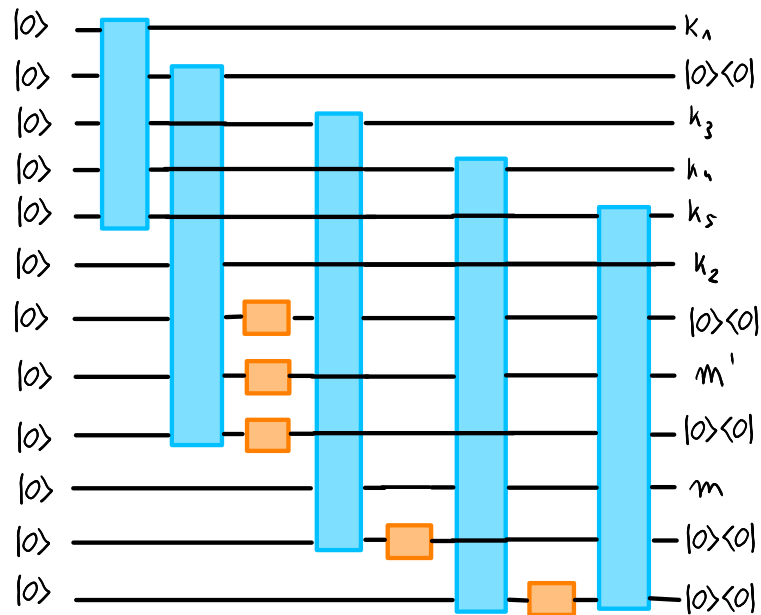
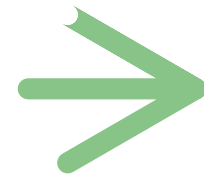
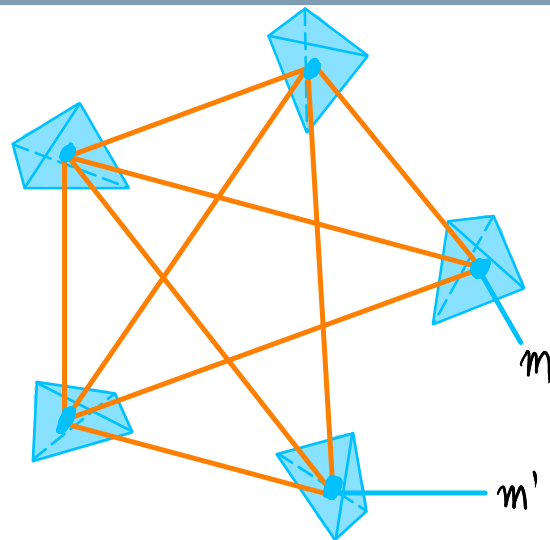
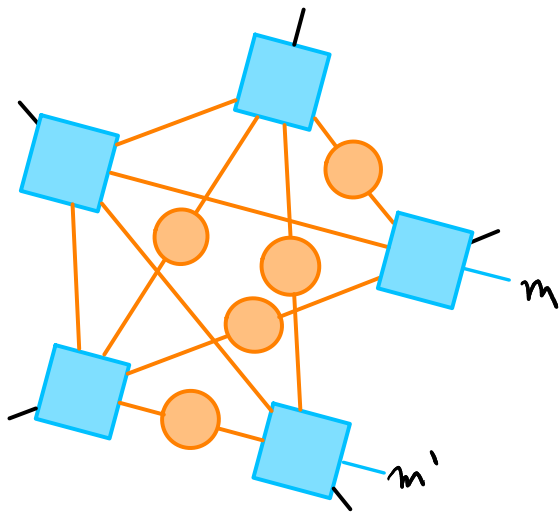
$$h = 1 \quad m' \text{ --- } m$$

$$h = \sigma_x \quad m' \text{ --- } [X] \text{ --- } m$$

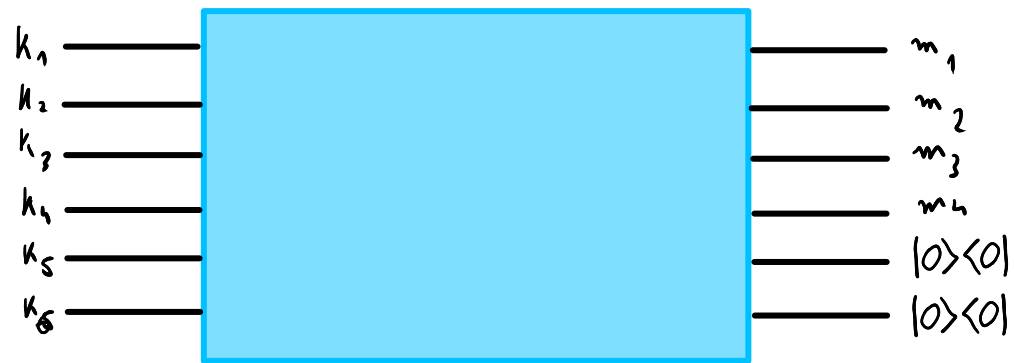
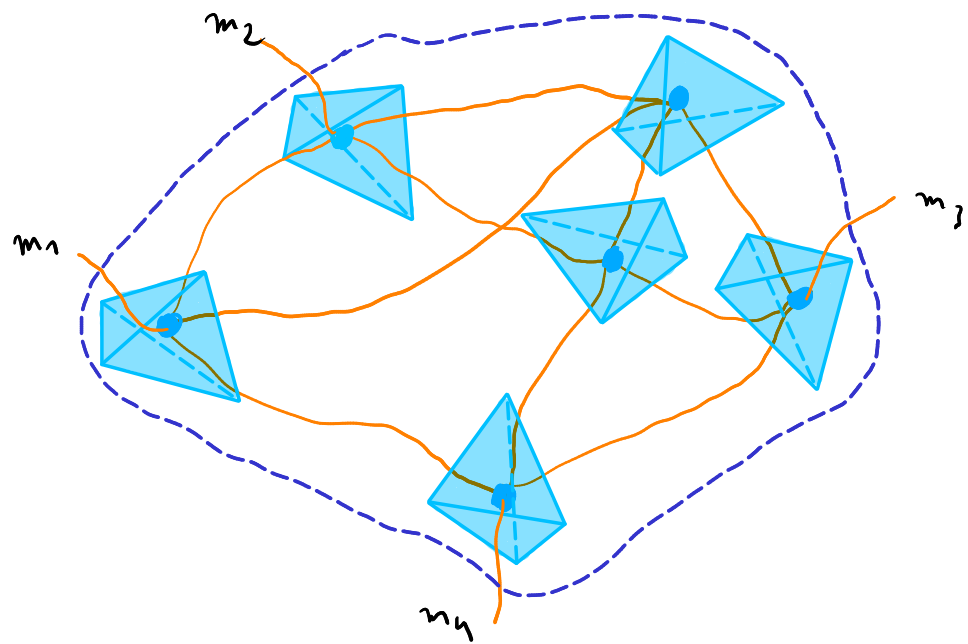
2ND METHOD: TENSOR NETWORKS



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BULK-BOUNDARY MAP



SUMMARY

- We can construct **explicit quantum circuit** for arbitrary 4-valent spin network with all spins $\frac{1}{2}$
- it can be used as a **tool for computations** that are classically intractable
- and as a **quantum information characteristics** of quantum gravity states, e.g.: their complexity