Particle physics & quantum black holes from asymptotically safe correlation functions

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Sant'Elmo beach hotel, October 4th 2023

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STRUCTURES CLUSTER OF EXCELLENCE



Bundesministerium für Bildung und Forschung









• Asymptotically safe correlation functions

• Applications I: asymptotically safe Standard Model

• Applications II: asymptotically black holes



Outline

Asymptotically safe correlation functions

JMP, Reichert, Front.in Phys. 8 (2021) 527

2309.10785

Background (in)dependence in gravity



Effective action

 $\Gamma_k[\bar{g},\bar{h}] = \Gamma_k[\bar{g}] + \Gamma_k^{(0,1)}[\bar{g}] * \bar{h} \cdot$

aka background and fluctuations fields, modified STIs and their importance

$$+\frac{1}{2}\Gamma_{k}^{(0,2)}[\bar{g}] * \bar{h}^{2} + \frac{1}{6}\Gamma_{k}^{(0,3)}[\bar{g}] * \bar{h}^{3} + \cdots \qquad (\bar{h} = \langle h \rangle)$$

JMP, Reichert, Front.in Phys. 8 (2021) 527 2309.10785



Background (in)dependence in gravity



Effective action

 $\Gamma_k[\bar{g},\bar{h}] = \Gamma_k[\bar{g}] + \Gamma_k^{(0,1)}[\bar{g}] * \bar{h} - \Gamma_k^{(0,1)}[\bar{g}] = \Gamma_k[\bar{g}] + \Gamma_k^{(0,1)}[\bar{g}] + \Gamma_k^{(0,1)}[\bar{g}$

$$\left\{\Gamma_k[\bar{g}]\right.$$

aka background and fluctuations fields, modified STIs and their importance

$$+\frac{1}{2}\Gamma_{k}^{(0,2)}[\bar{g}] * \bar{h}^{2} + \frac{1}{6}\Gamma_{k}^{(0,3)}[\bar{g}] * \bar{h}^{3} + \cdots \qquad (\bar{h} = \langle h \rangle)$$

$$\left\{ \Gamma_{k}^{(0,1)}[\bar{g}], \Gamma_{k}^{(0,2)}[\bar{g}], \Gamma_{k}^{(0,3)}[\bar{g}], \ldots \right\}$$



Background (in)dependence in gravity



$$\Gamma_k[\bar{g},\bar{h}] = \Gamma_k[\bar{g}] + \Gamma_k^{(0,1)}[\bar{g}] * \bar{h}$$

$$\left\{ \Gamma_k[\bar{g}], \Gamma_k^{(0,1)}[\bar{g}], \Gamma_k^{(0,2)}[\bar{g}], \Gamma_k^{(0,3)}[\bar{g}], \ldots \right\}$$

Se vogliamo che tutto rimanga come è, bisogna che tutto cambi.

II Gattopardo

aka background and fluctuations fields, modified STIs and their importance



From vertex dressings/distribution functions to physics

Effective action

$$\Gamma[\bar{g}, h, c_{\mu}, \bar{c}_{\mu}] = \int_{x} \left[\frac{2\Lambda - R}{16\pi G_{N}} + R f_{R}(\Delta) R + C f_{C}(\Delta) C + \cdots \right]_{\text{BRST-inv}} + S_{\text{gf}} + S_{\text{gh}}$$

Background effective action

$$\Gamma[g_{\mu\nu}] = \frac{1}{16\pi} \int_x \sqrt{g} \left\{ \mathcal{R}(\Delta, R) + Rf_{R^2}(\Delta)R + R_{\mu\nu}f_{R^2_{\mu\nu}}(\Delta)R^{\mu\nu} + \cdots \right\}$$



Se vogliamo che tutto rimanga come è, bisogna che tutto cambi.

II Gattopardo

aka form factors

Enforced by IR-UV consistence $R f_{R^2}(\Delta, R) R = \mathcal{R}(\Delta, R) + R f_{R^2}(\Delta) R$



From vertex dressings/distribution functions to physics

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$$\Gamma[\bar{g}, h, c_{\mu}, \bar{c}_{\mu}] = \int_{x} \left[\frac{2\Lambda - R}{16\pi G_{N}} + Rf_{R}(\Delta) \right] dt$$

Background effective action

$$\Gamma[g_{\mu\nu}] = \frac{1}{16\pi} \int_x \sqrt{g} \left\{ \mathcal{R}(\Delta, R) + R \right\}$$



Se vogliamo che tutto rimanga come è, bisogna che tutto cambi.

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 $Rf_{R^2}(\Delta)R + R_{\mu\nu}f_{R^2_{\mu\nu}}(\Delta)R^{\mu\nu} + \cdots$

gauge independent

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A lesson from graviton spectral functions



Fehre, Litim, JMP, Reichert, PRL 130 (2023) 081501

Se vogliamo che tutto rimanga come è, bisogna che tutto cambi.

II Gattopardo



Bonanno, Denz, JMP, Reichert, SciPost Phys. 12 (2022) 1, 001



2309.10785

A lesson from graviton spectral functions



Fehre, Litim, JMP, Reichert, PRL 130 (2023) 081501

 $\rho_h(\lambda) \in \mathbb{R}^+$

Spectral properties 'resemble' that of an asymptotic

$$\int_{\mathbb{R}} \frac{d\lambda}{2\pi} \,\lambda\rho_h(\lambda) = \infty$$

Se vogliamo che tutto rimanga come è, bisogna che tutto cambi. Il Gattopardo



Bonanno, Denz, JMP, Reichert, SciPost Phys. 12 (2022) 1, 001

Spectral properties of an unphysical mode

$$\rho_{\bar{g}}(\lambda) \in \mathbb{R} \qquad \qquad \int_{\mathbb{R}} \frac{d\lambda}{2\pi} \,\lambda \rho_{\bar{g}}(\lambda) = 0$$



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Spectral properties of an unphysical mode





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Graviton-graviton scattering



Bonanno, Denz, JMP, Reichert, SciPost Phys. 12 (2022) 1, 001



RG-invariant vertex



aka RG-invariant coupling /form factor

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Fluctuation approach: 2012 ...

Graviton-graviton scattering

RG-invariant vertex

 $\Gamma_{hhh}^{(3)}(p_1, p_2, p_3)$ $\overline{Z_{h}^{\frac{1}{2}}(p_{1})}Z_{h}^{\frac{1}{2}}(p_{2})Z_{h}^{\frac{1}{2}}(p_{3})$

aka RG-invariant coupling /form factor

Form factor approach: 2018 •••

Knorr, Ripken, Saueressig, ...



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Bonanno, Denz, JMP, Reichert, SciPost Phys. 12 (2022) 1, 001

Fluctuation approach: 2012 ...

Suggestive educated guess $\bar{\Gamma}_{\bar{g}^n}^{(n)}(p_1,...,p_n) \approx \frac{\Gamma_{h^n}^{(n)}(p_1,...,p_n)}{\pi^{\frac{1}{2}}}$ $Z_h^{\overline{2}}(p_1)\cdots Z_h^{\overline{2}}(p_n)$

Graviton-graviton scattering

RG-invariant vertex

 $\frac{\Gamma_{hhh}^{(3)}(p_1, p_2, p_3)}{\frac{1}{1}}$ $\overline{Z_{h}^{\frac{1}{2}}(p_{1})Z_{h}^{\frac{1}{2}}(p_{2})Z_{h}^{\frac{1}{2}}(p_{3})}$

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Some thoughts on vertex expansion schemes/fluctuation approach and covariant expansion schemes/background approximation

A bit of gauge invariance and all that

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Some thoughts on vertex expansion schemes/fluctuation approach and covariant expansion schemes/background approximation

Clear cut case: Einstein-Hilbert truncation & minimally coupled gravity matter systems

Exact map with additional truncation

Fluctuation approach with flat vertex expansion

Additional truncation: background approximation

A bit of gauge invariance and all that

Background approximation with heat kernel methods

Background approximation flows can be derived also in the

flat vertex expansion







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Clear cut case: Einstein-Hilbert truncation & minimally coupled gravity matter systems

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Fluctuation approach with flat vertex expansion

Additional truncation: background approximation

claim also tested in QCD



where there are enough benchmark results from the lattice

aka: if there are qualitative differences between results in the fluc. approach and the background approx., the former is conceptually more trustworthy

- (a) better account of dynamics
- (c) better -tested- systematics

Background approximation with heat kernel methods

Background approximation flows can be derived also in the flat vertex expansion

Relation to background approach via pinch technique

(b) better account of physical diffeomorphism invariance









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there are qualitative differences between flows in the fluc. approach & the background approximation in pure gravity and minimally coupled systems

A bit of gauge invariance and all that



Background approximation with heat kernel methods

Background approximation flows can be derived also in the flat vertex expansion



Relation to background approach via pinch technique

Beware













Bonanno et al., Critical reflections on asymptotically safe gravity, Front.in Phys. 8 (2020) 269

QCD & SM thresholds in the RG since (many) decades

QCD with the fRG since a decade

Bonanno et al., Critical reflections on asymptotically safe gravity, Front.in Phys. 8 (2020) 269



Pastor-Gutiérrez, JMP, Reichert, SciPost Phys. 15 (2023) 105



k-mirrors of physical thresholds

QCD & SM thresholds in the RG since (many) decades

Example: asymptotically safe Standard Model

QCD with the fRG since a decade

JMP, Reichert, Front.in Phys. 8 (2021) 527 2309.10785

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Towards apparent convergence in quantum gravity

vertex expansion



Aiming at apparent convergence



Towards apparent convergence in quantum gravity

vertex expansion



Aiming at apparent convergence





Aiming at apparent convergence





Aiming at apparent convergence

geometrical approach: Donkin, JMP, arXiv:1203.4207 flat expansion: Christiansen, Litim, JMP, Rodigast, PLB 728 (2014) 114

level 2:
$$\Gamma^{(m,n)} \approx \Gamma^{(m+n-2,2)}$$

$$Z_h(p), \ Z_c(p), \ \mu = -2\lambda_2$$





Aiming at apparent convergence





Aiming at apparent convergence



























 ${
m R}^2$ - tensor structure generated

${f R}$ - tensor structure sustained











 ${
m R}^2$ - tensor structure generated

${f R}$ - tensor structure sustained









Typically diagrams with higher order vertices are strongly suppressed

(a) couplings stay finite

(b) combinatorical suppression of diagrams with higher vertices

(c) phase space (angular) suppression of diagrams with higher vertices

turns out to be very efficient!

Towards apparent convergence in quantum gravity



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Why does/could it fail?

Resonant interaction channels and their interactions circumvent (b) and make (a) irrelevant

(a) couplings diverge

(b) hadrons, diquarks, glueballs, ... in QCD

Towards apparent convergence in quantum gravity

Emergent composites, BSE

Gies, Wetterich, PRD 65 (2002) 0650016 JMP, AP 322 (2007) 2831 Flörchinger, Wetterich, PLB 680 (2009) 371



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QG as perturbative as possible & apparently converging

... slight oversimplification for the sake of this talk ...



Applications I: asymptotically safe Standard Model

Dona, Eichhorn, Percacci, PRD 89 (2014) 084035

Meibohm, JMP, Reichert, EPJC 76 (2016) 285 Christiansen, Litim, JMP, Reichert PRD 97 (2018) 4, 046007

> Shaposhnikov, Wetterich, PLB 683 (2010) 196 Eichhorn, Versteegen, JHEP 1801 (2018) 030 Eichhorn, Held, PRL 121 (2018) 151302

Latest 'status report': Eichhorn, Schiffer, 2212.07456

Pastor-Gutiérrez, JMP, Reichert, SciPost Phys. 15 (2023) 105





Experimental value (PDG)

top pole mass (getting real)

 $M_{t,\text{pole}}^{(\text{exp})} = 172.5 \pm 0.7 \,\text{GeV}$







Experimental value (PDG)

top pole mass (getting real)



Euclidean curvature mass







Experimental value (PDG)

 $\Gamma_{t,\mathrm{po}}^{(\mathrm{the})}$

top pole mass (getting real)



Euclidean curvature mass

Prediction of decay width

$$_{\text{ole}}^{\text{eo})} = 1.72^{+0.09}_{-0.41} \,\text{GeV}$$

$$\Gamma_{t,\text{pole}}^{(\text{exp})} = 1.42^{+0.19}_{-0.15} \,\text{GeV}$$

Experimental value (PDG)









flat Higgs FP potential 2.0 N_{\max} : 2 _ 10 1.53 _ 11 $(u^* \left(ar{
ho}
ight) - u_0^*
ight) \cdot 10^3$ 1.04 - 125**—** 13 0.56 - 140.0 7 — 15 8 _ 16 -0.59 -1.00.020.010.03 0.04 0.050.00 0.060.07 0.08 0

Two relevant directions

 $\Gamma_{t,\mathrm{po}}^{(\mathrm{the})}$

Asymptotically safe Standard Model

top pole mass (getting real)



Experimental value (PDG)

Euclidean curvature mass

Prediction of decay width

$$_{\text{ole}}^{\text{eo})} = 1.72^{+0.09}_{-0.41} \,\text{GeV}$$

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Experimental value (PDG)









Applications II: asymptotically black holes

Platania, 2309.17043 Black Holes in Asymptotically Safe Gravity: and beyond: Held, Eichhorn, 2212.09495



Unfolding the background effective action

$$\begin{split} \Gamma[g_{\mu\nu}] &= \frac{1}{16\pi} \int_x \sqrt{g} \left\{ \mathcal{R}(\Delta, R) \right. \end{split}$$

$$\end{split}$$

$$\begin{split} & \left[\bar{\Gamma}_{hh}^{(2)}(p) \right] \\ & \left[\bar{\Gamma}_{hh}^{(3)}(p) \right] \\ & \left[\bar{\Gamma}_{h^3}^{(4)}(p) \right] \end{split}$$

 $) + Rf_{R^2}(\Delta)R + R_{\mu\nu}f_{R^2_{\mu\nu}}(\Delta)R^{\mu\nu} + \cdots \bigg\}$

RG-invariant

 $\bar{\Gamma}_{\bar{g}^3}^{(3)}(p)$

 $\bar{\Gamma}_{\bar{g}^4}^{(4)}(p)$



Unfolding the background effective action

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$$\begin{split} \mathbf{g}_{auge \ dependent} & \mathbf{F}_{hh}^{(2)}(p) \\ & \bar{\Gamma}_{hh}^{(3)}(p) & \bullet \\ & \bar{\Gamma}_{h^3}^{(4)}(p) \\ & \bar{\Gamma}_{h^4}^{(4)}(p) \end{split}$$



 $\bar{\Gamma}_{\bar{g}^3}^{(3)}(p)$

 $\bar{\Gamma}_{\bar{g}^4}^{(4)}(p)$

 $\mathcal{R}(\Delta, R)$

 $Rf_{R^2}(\Delta)R$

 $R_{\mu\nu}f_{R^2_{\mu\nu}}(\Delta)R^{\mu\nu}$



Unfolding the background effective action

$$\begin{split} \Gamma[g_{\mu\nu}] &= \frac{1}{16\pi} \int_{x} \sqrt{g} \left\{ \mathcal{R}(\Delta, R) + Rf_{R^{2}}(\Delta)R + R_{\mu\nu}f_{R^{2}_{\mu\nu}}(\Delta)R^{\mu\nu} + \cdots \right\} \\ \end{split}$$



Unfolding the background effective action

$$\begin{split} \Gamma[g_{\mu\nu}] &= \frac{1}{16\pi} \int_{x} \sqrt{g} \left\{ \mathcal{R}(\Delta, R) + Rf_{R^{2}}(\Delta)R + R_{\mu\nu}f_{R^{2}_{\mu\nu}}(\Delta)R^{\mu\nu} + \cdots \right\} \end{split}$$
gauge dependent RG-invariant gauge independent R(\Delta, R)
Unfolding
 $\bar{\Gamma}^{(3)}_{h^{3}}(p)$
 $\bar{\Gamma}^{(3)}_{h^{3}}(p)$
 $\bar{\Gamma}^{(4)}_{h^{4}}(p)$
 $\bar{\Gamma}^{(4)}_{h^{4}}(p)$
R(\Delta, R)
Unfolding R(A)
R(A)

Educated guess



Unfolding the background effective action





10

$$+ Rf_{R^{2}}(\Delta)R + R_{\mu\nu}f_{R^{2}_{\mu\nu}}(\Delta)R^{\mu\nu} + \cdots \bigg\}$$

RG-invariant

gauge independent

$$\bar{\Gamma}^{(3)}_{\bar{g}^{3}}(p)$$

R $f_{R^{2}}(\Delta)R$

Guated guess

R $f_{R^{2}}(\Delta)R$

Guated guess

R $f_{R^{2}}(\Delta)R$

R $f_{R^{2}}(\Delta)R$

R $f_{R^{2}}(\Delta)R$

Results for form factors

$$\mathcal{R}(\Delta, R) = R \frac{\gamma_g^{(3)}(\Delta) - \bar{\gamma}_3 \Delta}{\Delta + R} R$$



Infrared asymptotic effective action

$$\Gamma_{\rm IR}[g_{\mu\nu}] = \frac{1}{16\pi} \int_x \sqrt{-g} \left(G_N^{-1}R + g_{R_{\mu\nu}^2} R_{\mu\nu} R^{\mu\nu} + g_{R^2} R^2 + c_1 R_{\mu\nu} \Box R^{\mu\nu} + c_2 R \Box R \right)$$

Spherical symmetric solution

$$ds^{2} = -f(r)dt^{2} + \frac{1}{g(r)}dr^{2} + r^{2}d\Omega^{2}$$

Weak field solutions

$$f(r) = 1 - \frac{2M}{r} + S_0 \frac{e^{-m_0 r}}{r} + S_2 \frac{e^{-m_2 r}}{r}$$
$$g(r) = 1 - \frac{2M}{r} - S_0 \frac{e^{-m_0 r}}{r} (1 + m_0 r) + \frac{1}{2} S_2 \frac{e^{-m_2 r}}{r} (1 + m_2 r)$$



Weak field solutions







Weak field solutions





$$T = \frac{1}{4\pi} \sqrt{|f'(r_h)|}$$

3.0

2.5

2.0

1.5

1.0

0.5

'Hawking temperature' see also Borissova, Held, Afshordi, CQuant.Grav. 40 (2023) 7, 075011



$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \left(\begin{array}{c} & & & \\ & &$$

Asymptotically safe SM



Summary



Quantum black holes

Type III 6 5 4 S_0 3 Type II 2 1 0 2.5 3.0 0.5 1.0 1.5 2.0 M

Phase structure

3.0

2.5

2.0

1.5

1.0

0.5