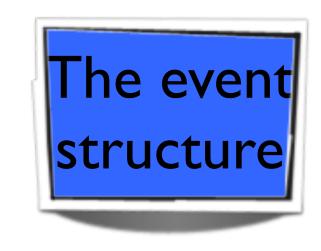
# I, 2 or 3 Higgses at the LHC and in the future

Pamela Ferrari (Nikhef and Radboud University)





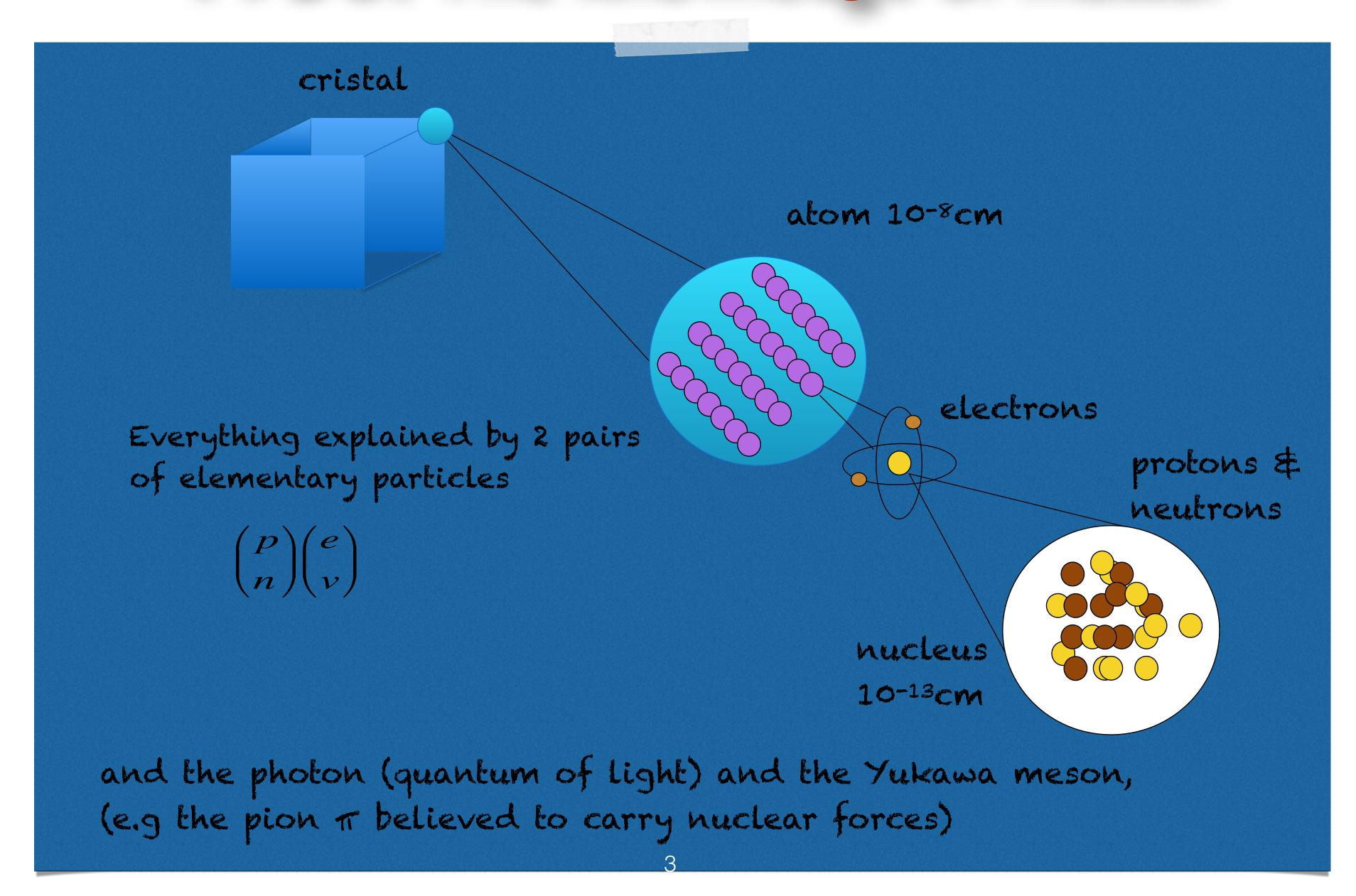


# Historical background

Shortly introducing the Standard Model and its shortcomings to motivate our studies at the LHC

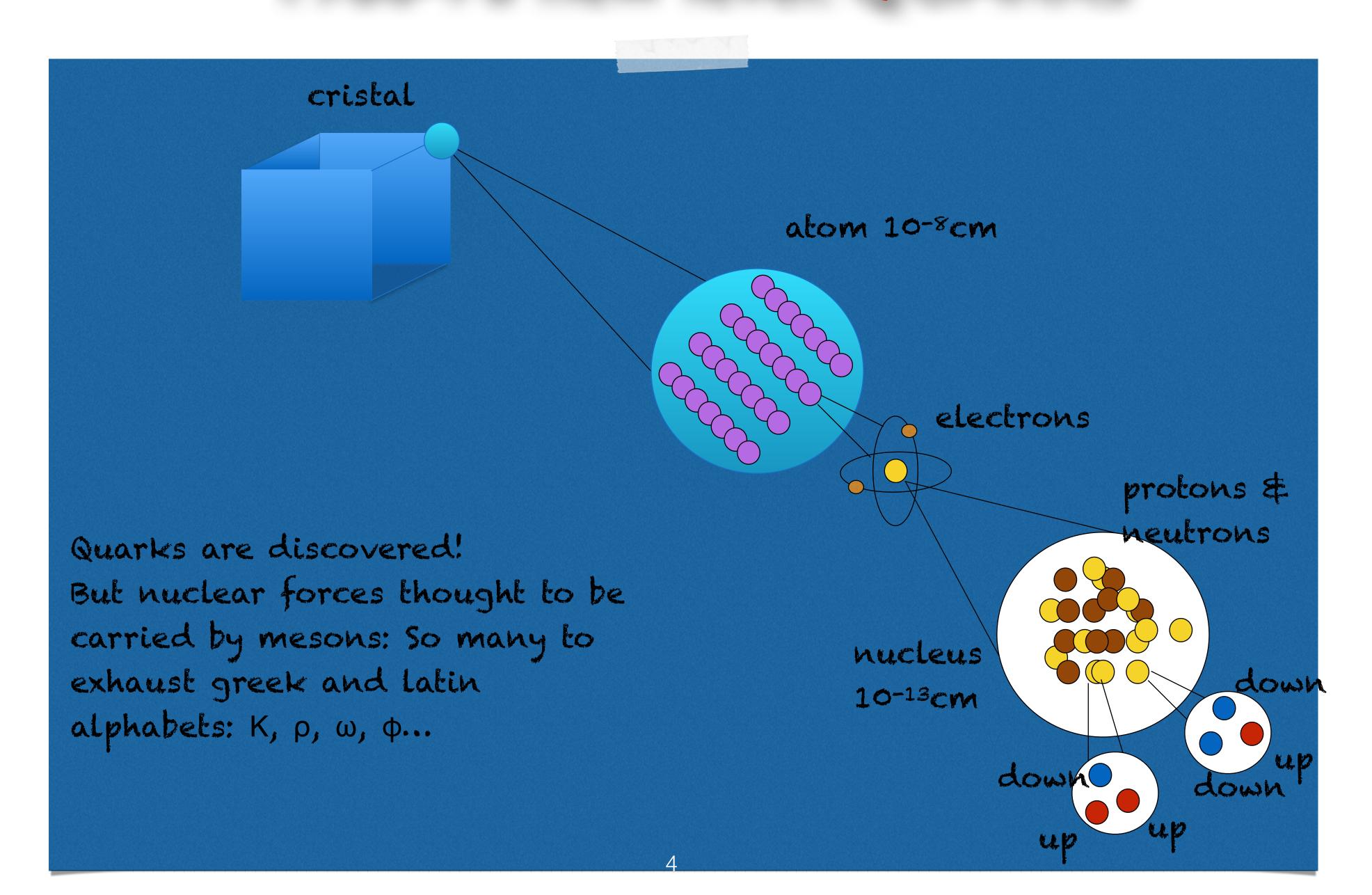


# 1930: The knowledge of matter

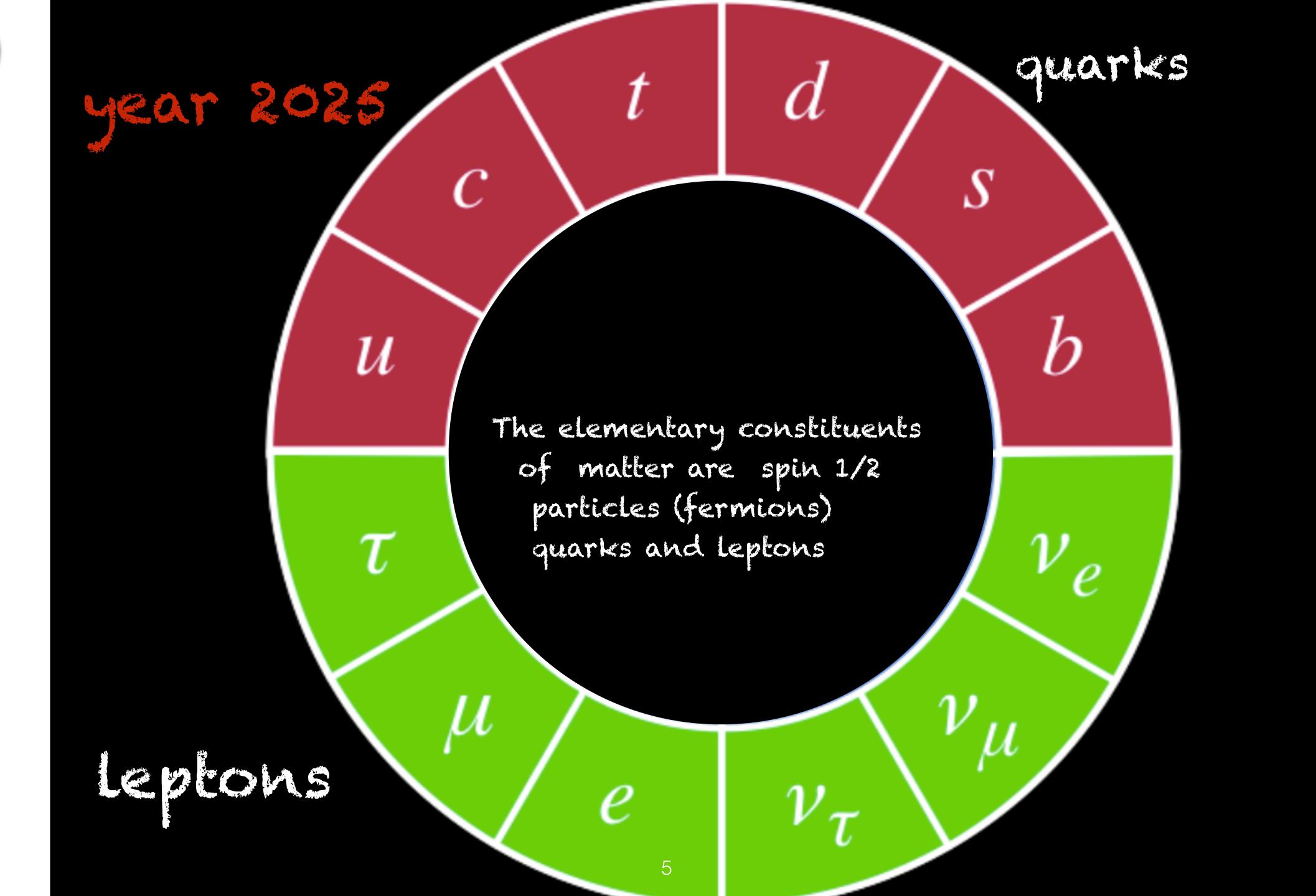




#### 1960-70 new level: QUARKs



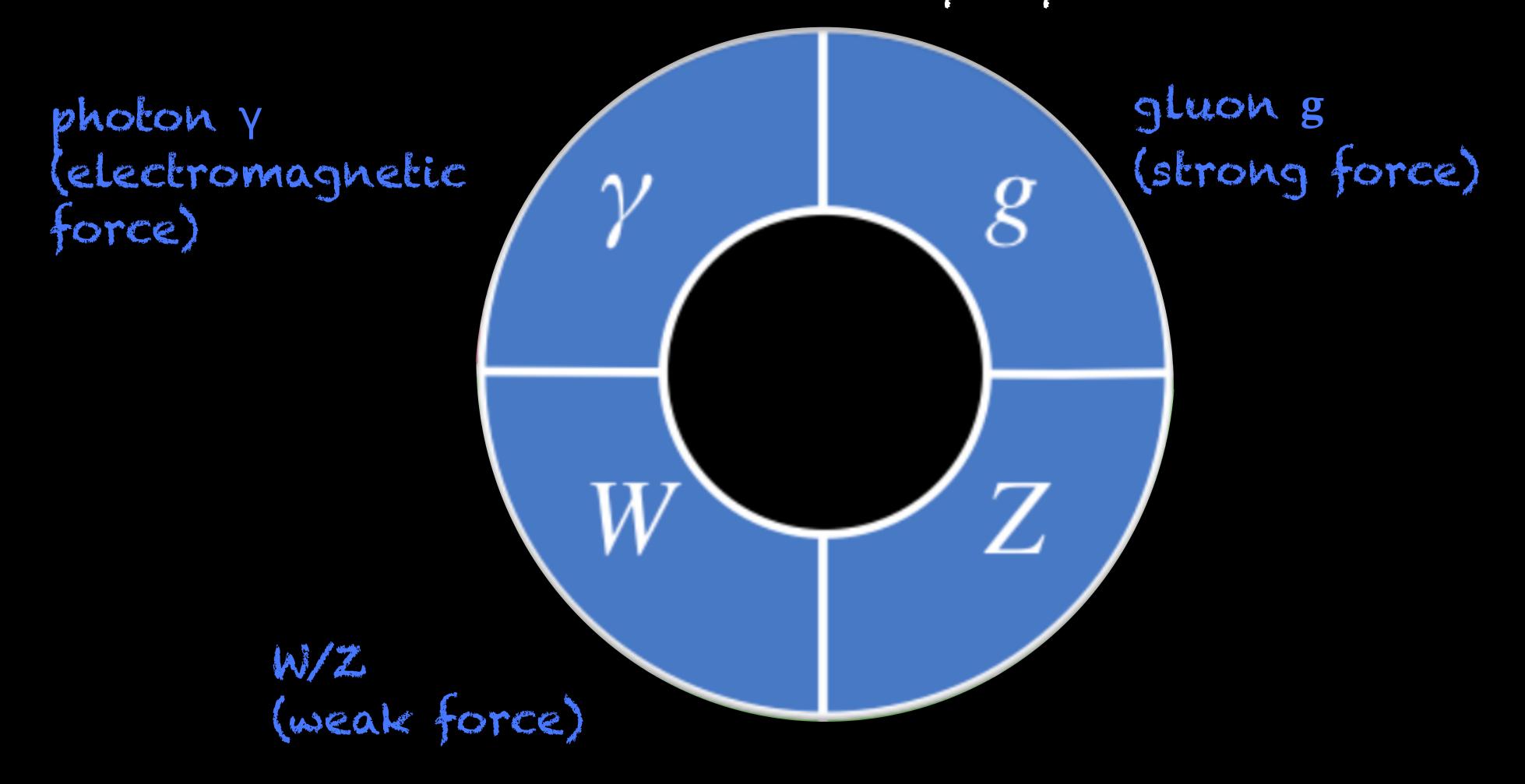
History



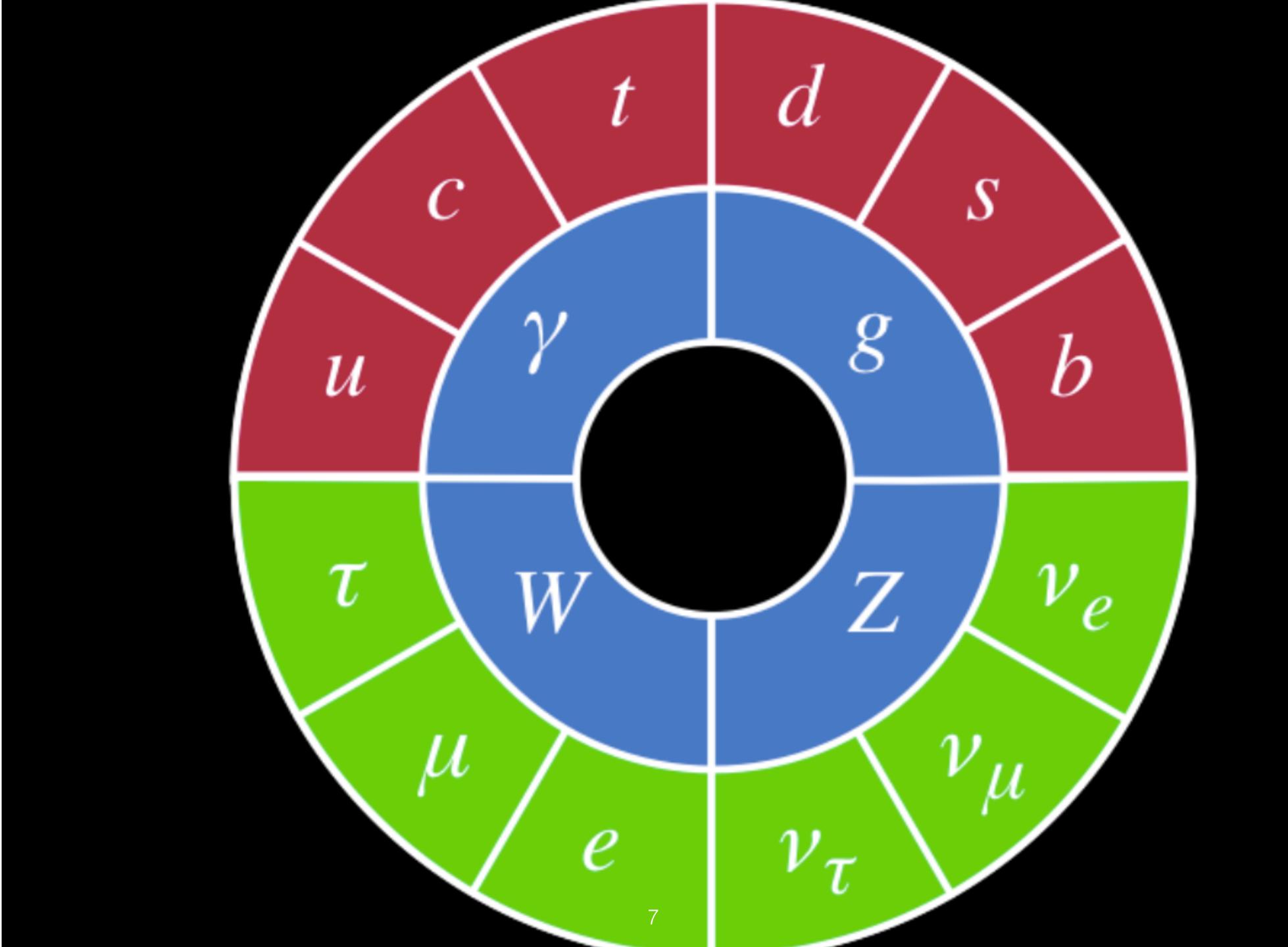


# Mhat about forces?

The force carriers are integer spin particles (bosons)



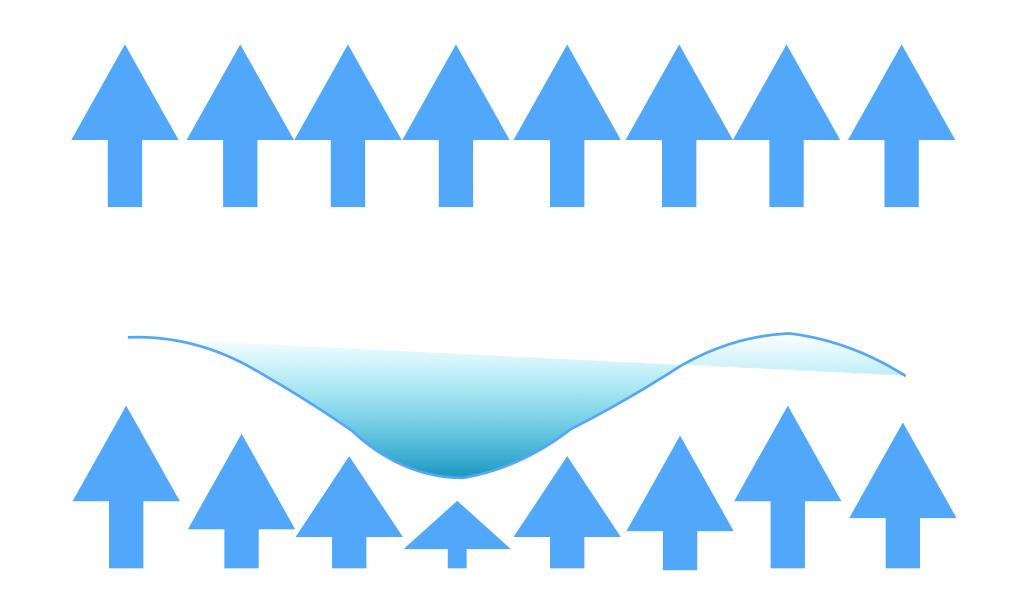
History





### What is responsible of particles masses?

The Higgs field that fills the space. Particles get mass by interacting with it.



The vacuum is like the surface of still lake

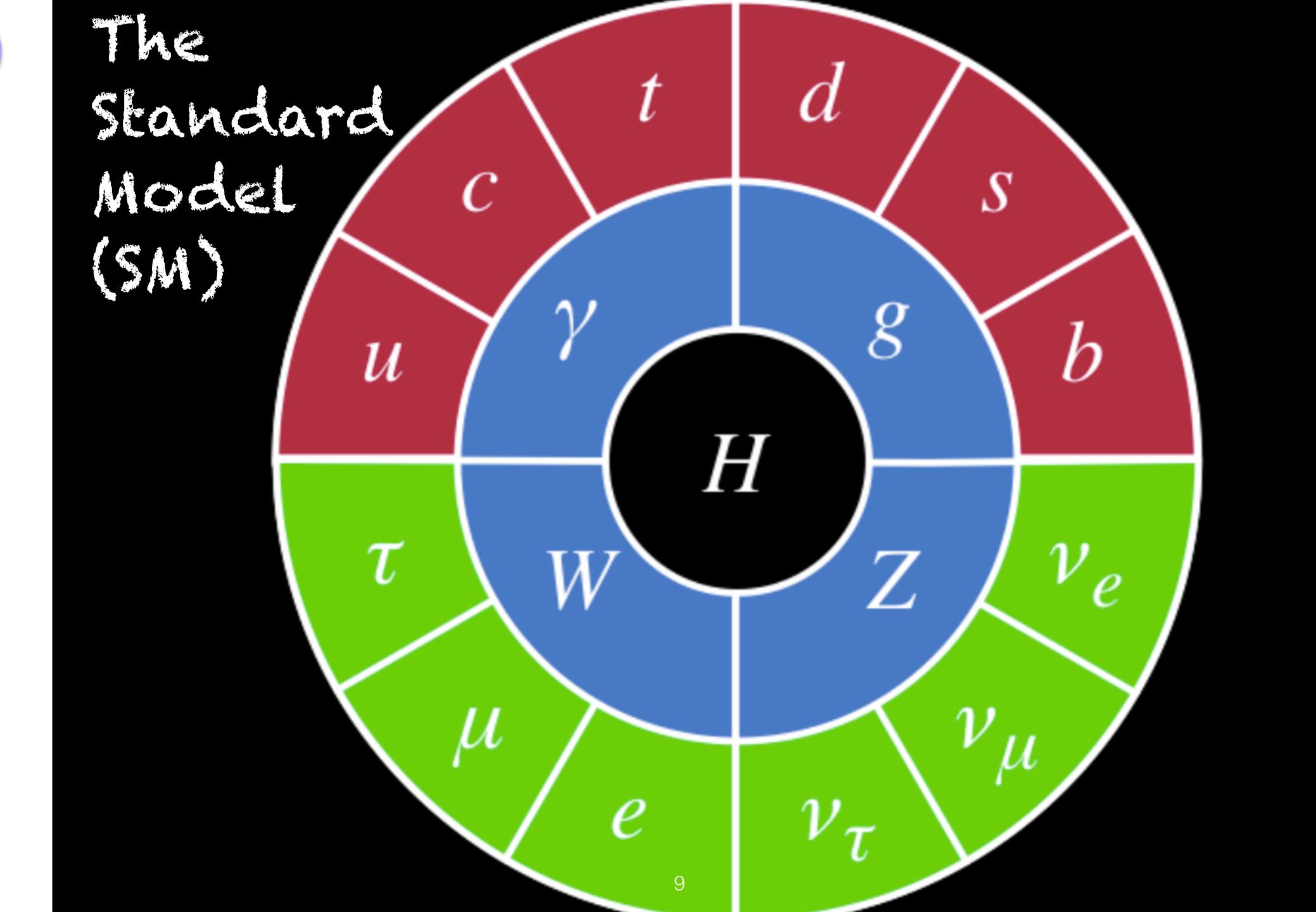
collisions produce waves (oscillation of field =particle)



Higgs boson spin zero particle, m~125 GeV

Vacuum is filled!

History

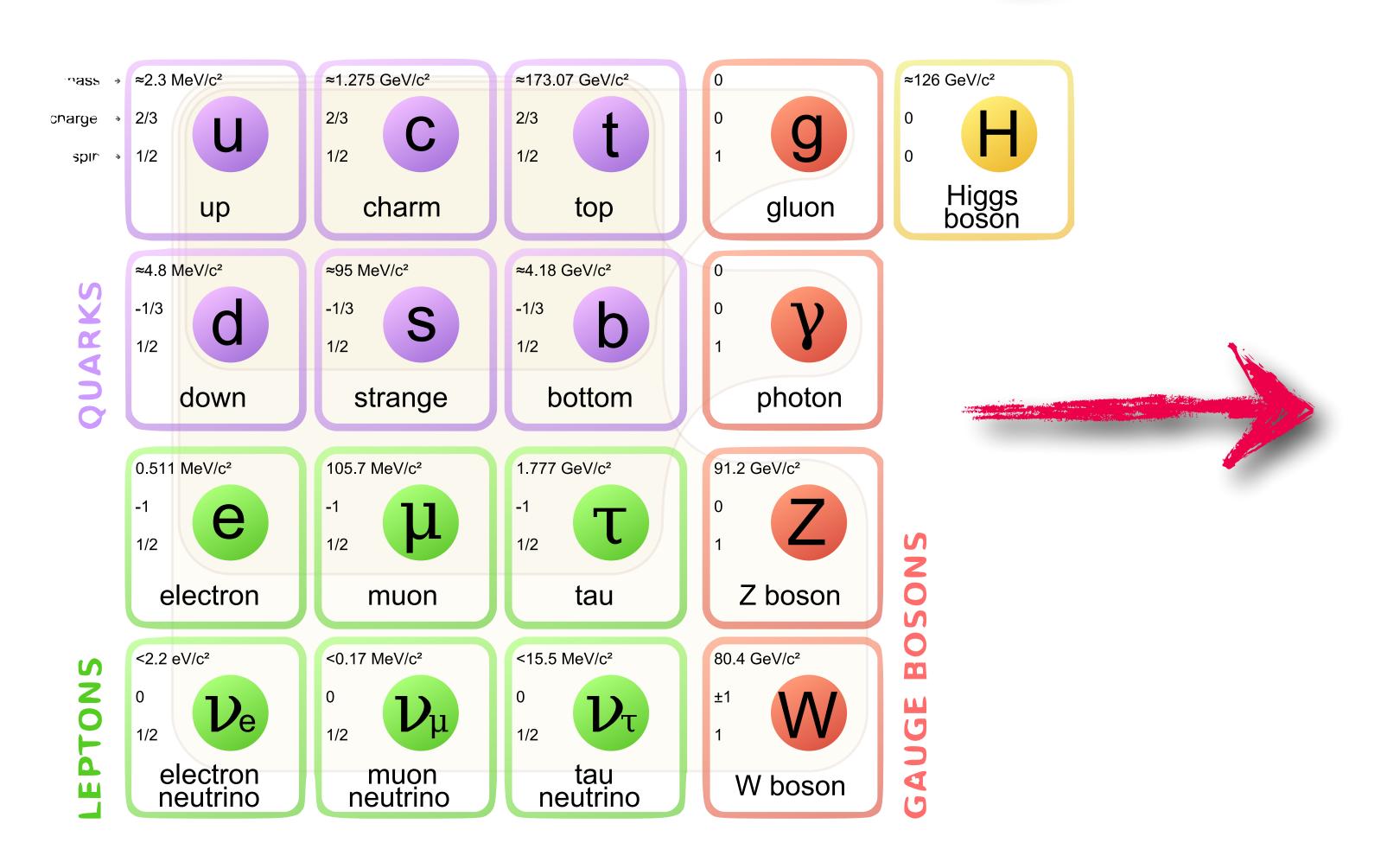




- •The Standard Model predicts that the Higgs boson and field acts as a sole player in the game of Electroweak symmetry breaking.
- •This is a strong prediction that has yet to be verified experimentally.
- •Answering this question is one of the pressing goals for the future



#### What is the mass/energy scale we are talking about?



New physics may appear at higher scales

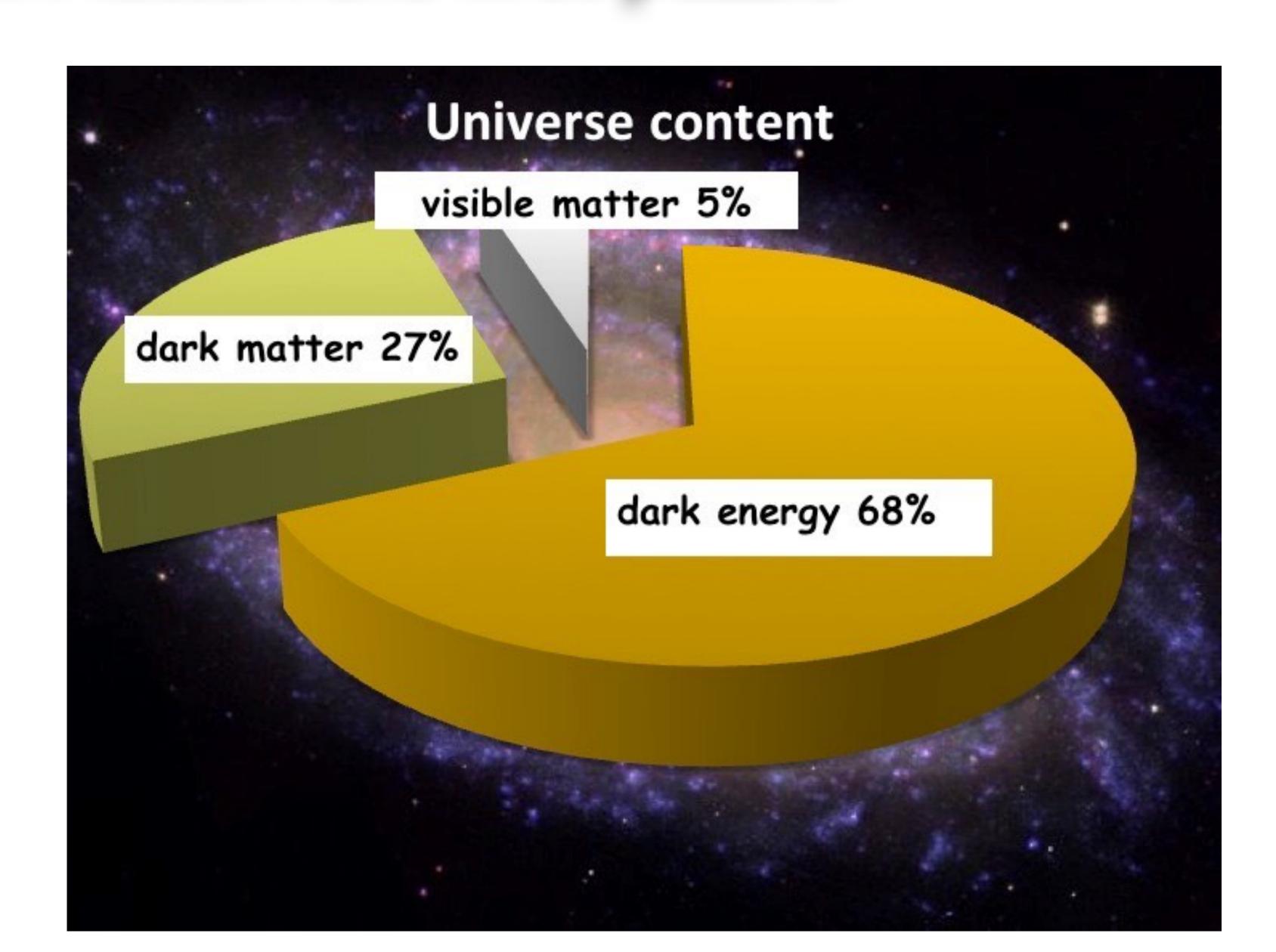
below 200 Giga-electronVolt = GeV



# Dark Matter: the first puzzle

Stable Ordinary particles
Can they explain everything?

No!





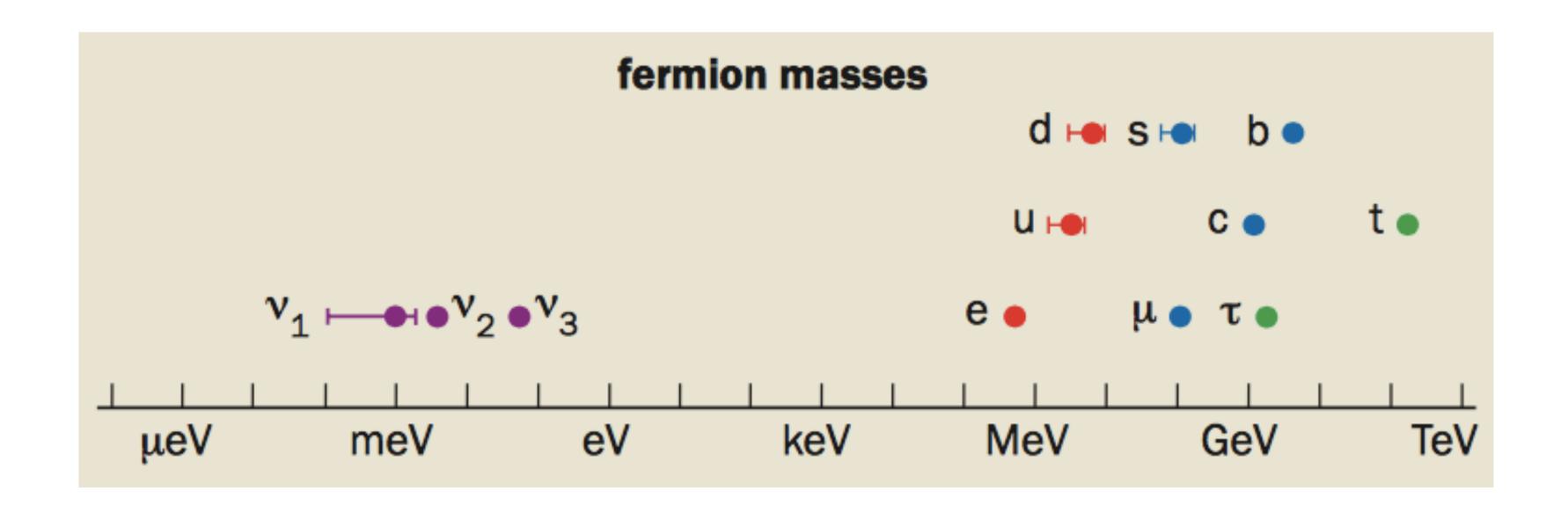
#### Why is our universe made of matter and not anti-matter?

- 2. One of the major shortcomings of our understanding of particle physics is the matter over antimatter asymmetry in the universe.
- 3. While the Standard Model does predict a matter vs. anti-matter asymmetry, it is much too small compared to what we observe.
- 4. Moreover, the thermal history of electroweak symmetry breaking is important for particle physics and cosmology. If in the early universe, there was a first order electroweak phase transition (think boiling water), this could explain the matter vs. anti-matter asymmetry as well as sources for potentially observable gravitational radiation.
- 5. The Standard Model's prediction is again clear no first- order transition. Therefore if such a transition took place, the Higgs doesn't act alone and some new physics is present.



#### The neutrino masses....

2. How to incorporate in SM neutrino masses? Why are they so small?



New physics should appear in the worst case at Planck scale where quantum gravity effects become important and quantum filed theory breaking down. Very large scale  $1.22 \times 10^{19}$  GeV!



## Hierarchy Problem

If SM is a complete description of Nature



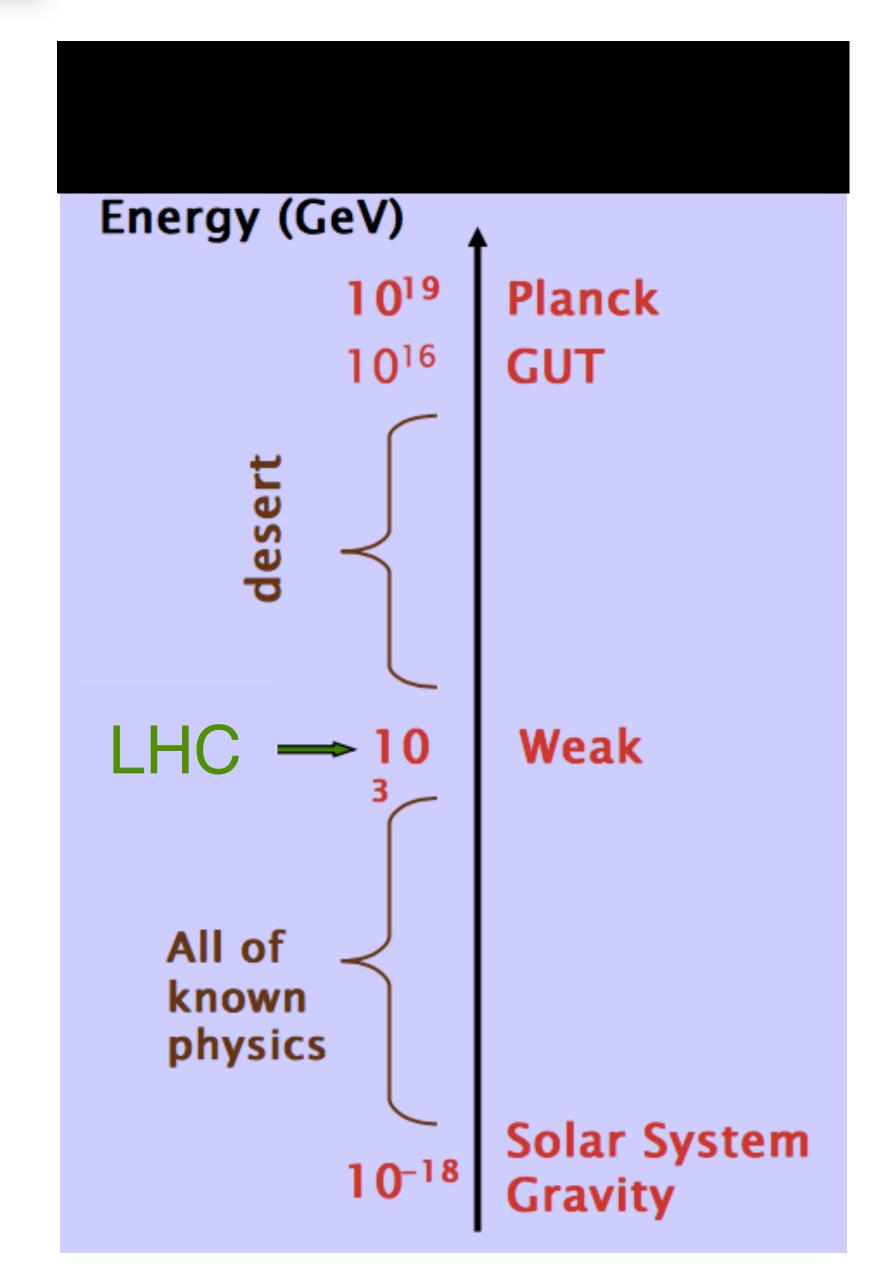
no hierarchy problem.

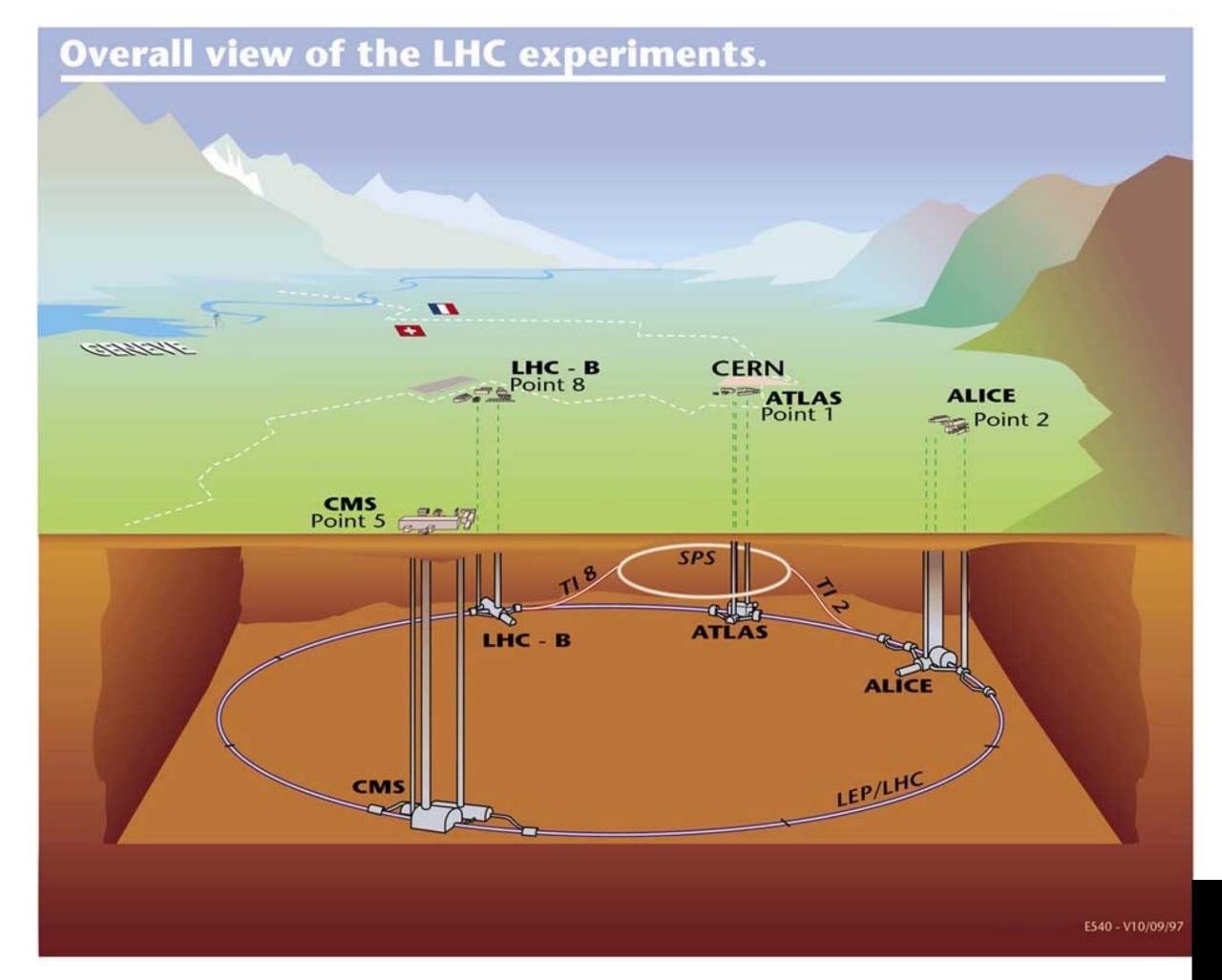
But the SM has unresolved issues which point to New Physics (NP)
If NP appears at Planck scale unnatural large difference wrt EW Scale

Most accredited models predict NP @TeV scale.

ITeV=1000GeV





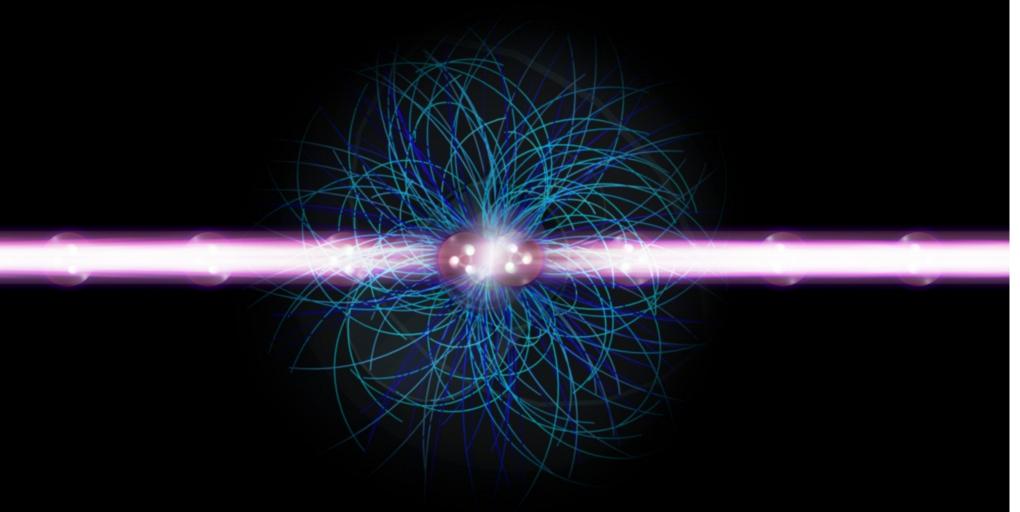




LHC accelerates protons in opposite directions along a 28 km ring.

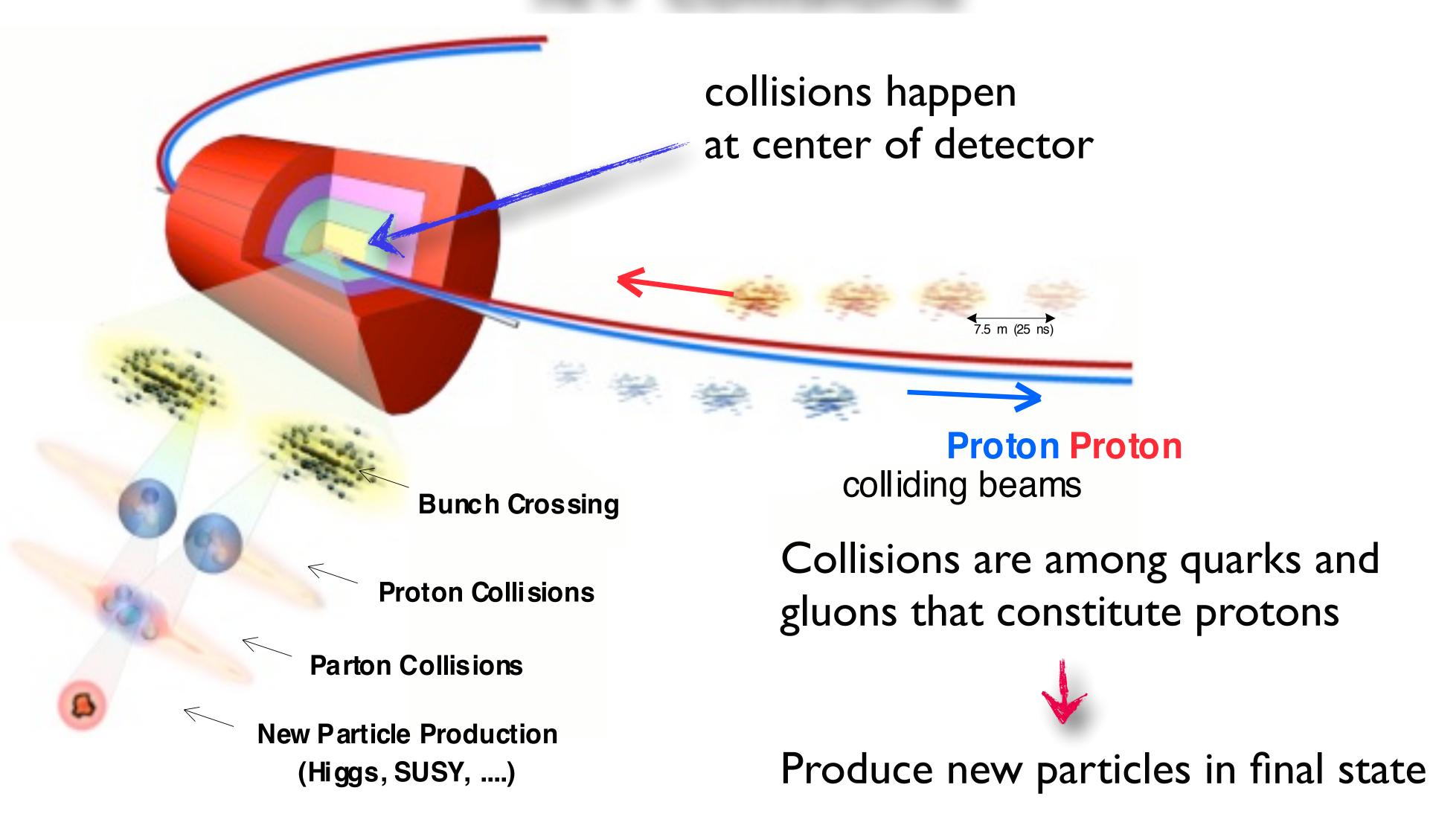
Protons collide at experiments @13.6 TeV

Collisions bring us back to Big Bang producing particles abundant at that time





#### TeV collisions



# particle jet parti

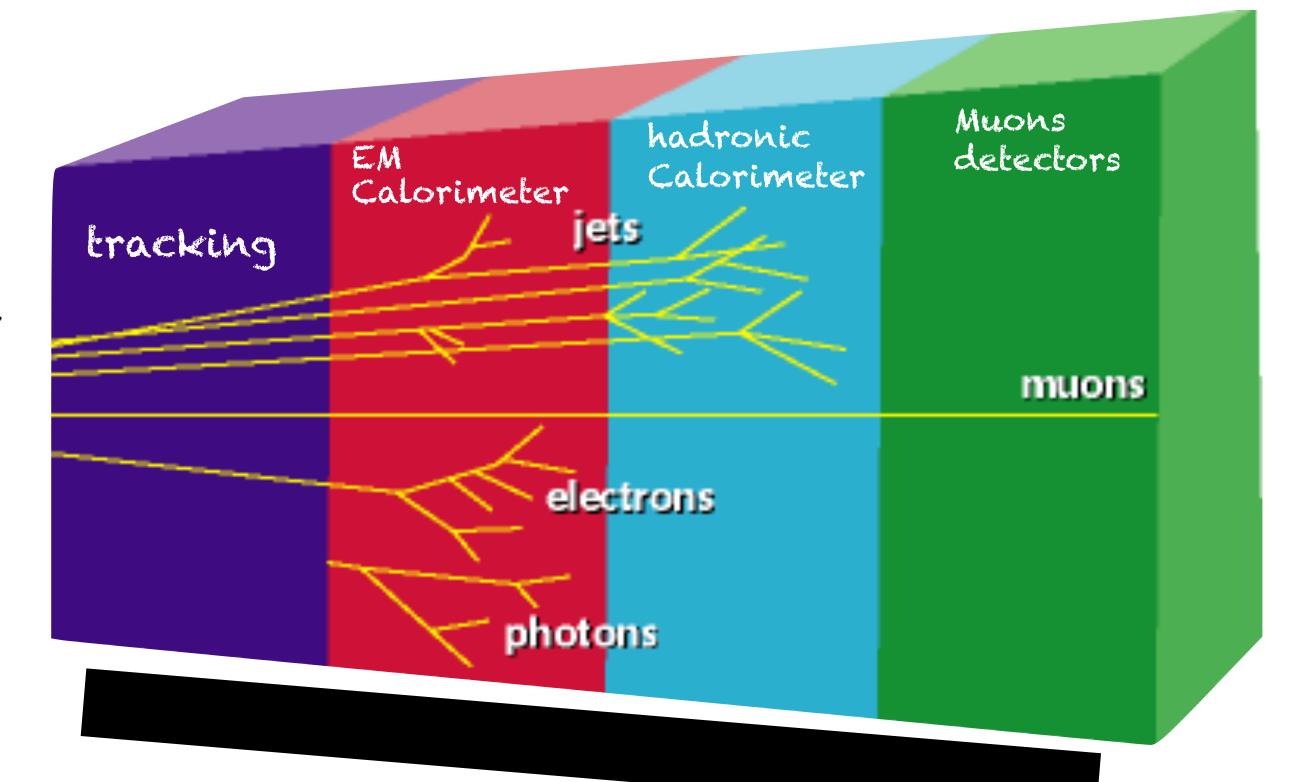
#### Typical detector

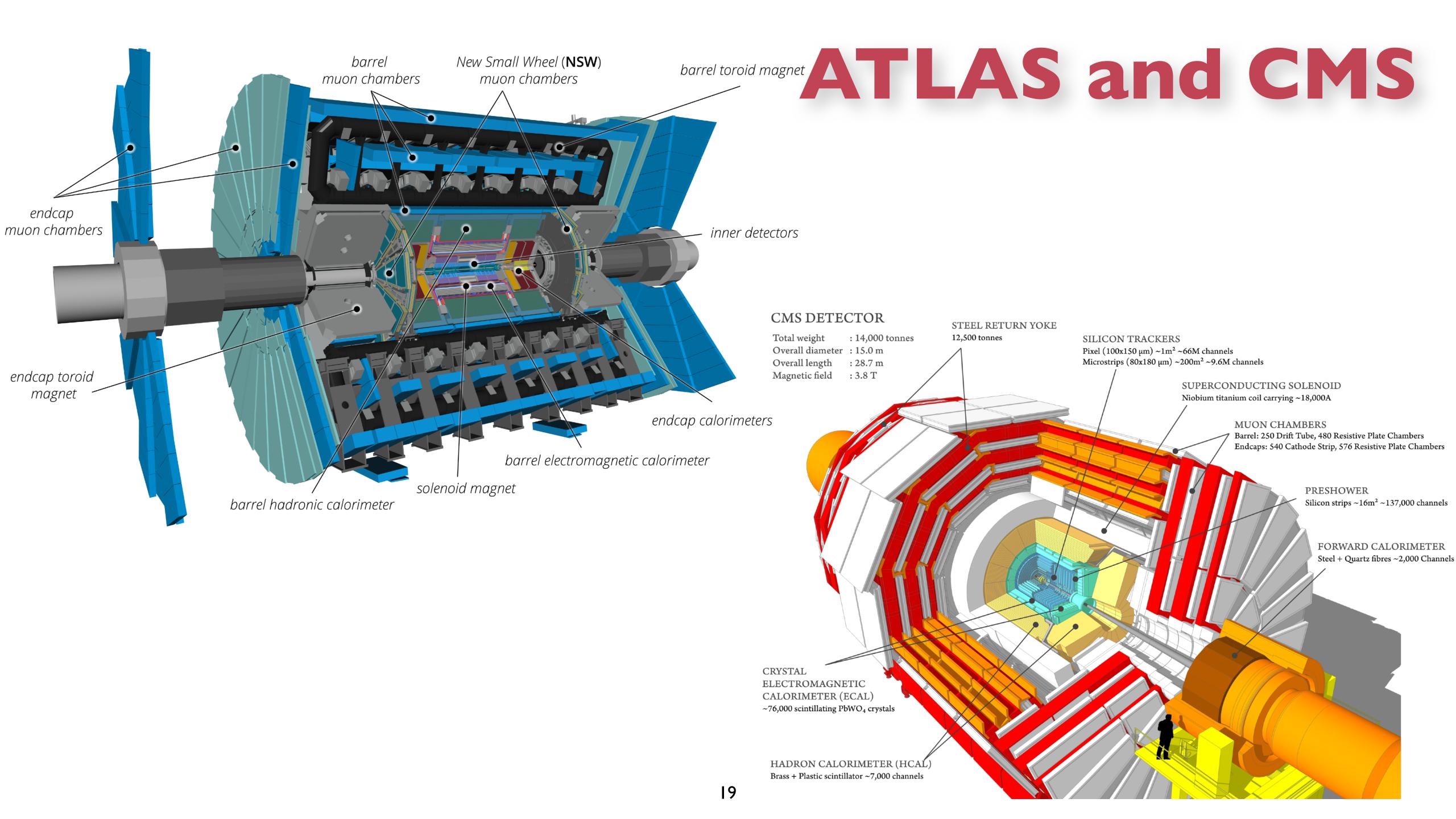
Detectors built to observe particles produced in collisions

quark/gluons seen as jets of particles in a narrow cone

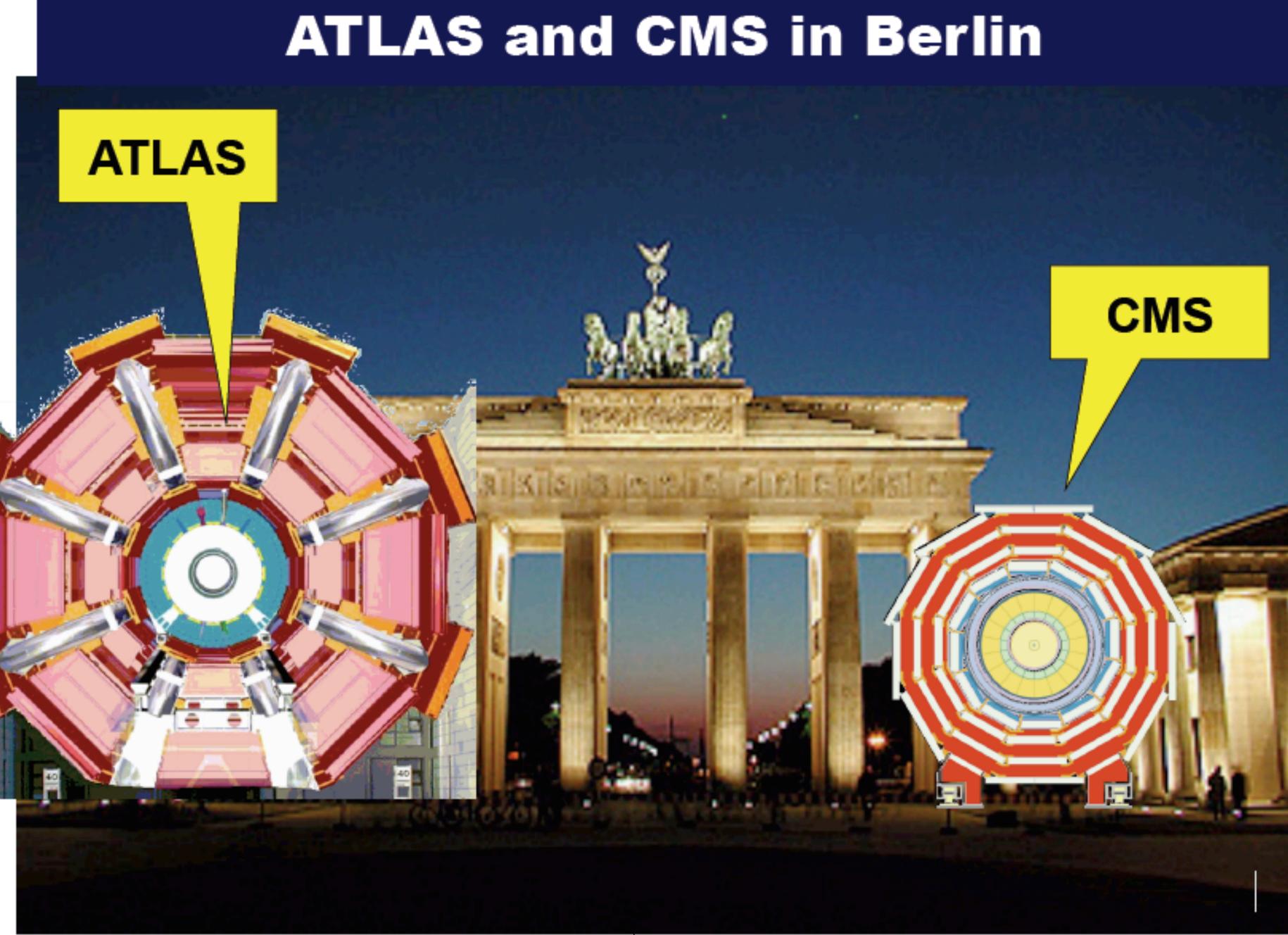
particles interact differently with detector

used to disentangle them











# ATLAS Detector Trigger

- Trigger (online event selection for permanent storage) is of paramount importance since is the first cut applied any physics analysis
- Two level trigger system

**↓** 40 MHz

#### Level-1 (L1)

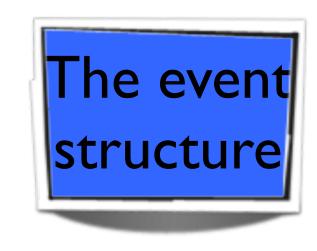
- Hardware-based trigger
- Inputs from Calorimeter and Muon systems with coarse detector granularity defining Regions of Interest (Rols)
- Latency:  $< 2.5 \mu s$

↓ 100 kHz

#### **High Level Trigger (HLT)**

- Software-based trigger
- Full detector granularity
- Latency:  $\sim$  0.5 s average

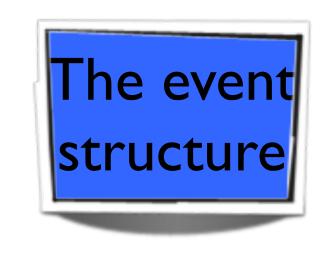
1 kHz average × 1 MB/event = 1 GB/s



# The Structure of an event

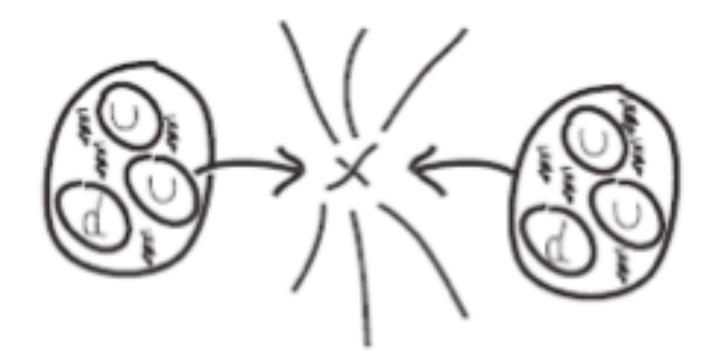
I will here introduce few concepts by showing what happens when two protons interact, i.e.

- hard scatters
- •Radiation: ISR/FSR
- Pile-up
- Parton Density Functions (PDF's)



# Few useful reminders on hadron collider kinematics

Protons (and antiprotons) are formed by quarks (uud) kept together by gluons

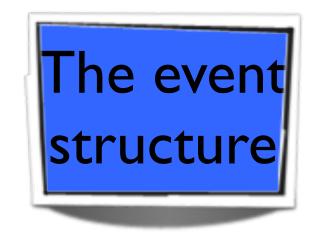


The energy of each beam is carried not by the entire proton, but by one of its constituents

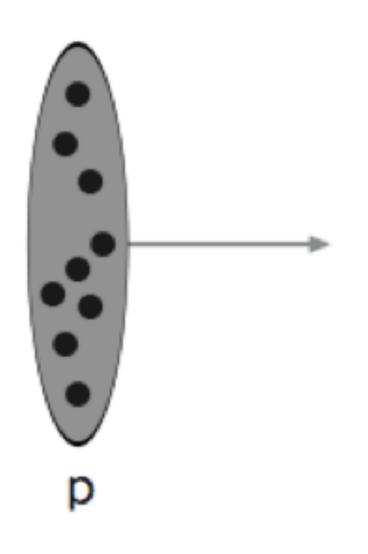
Ecollision < 2Eb

Pros: with a single energy possible to scan different processes at different energies

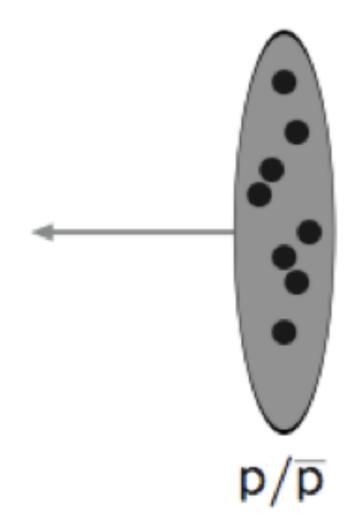
Const the energy available for the collision is lower than the accelerator energy



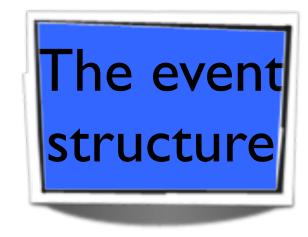
#### The Structure of an event: PDFs



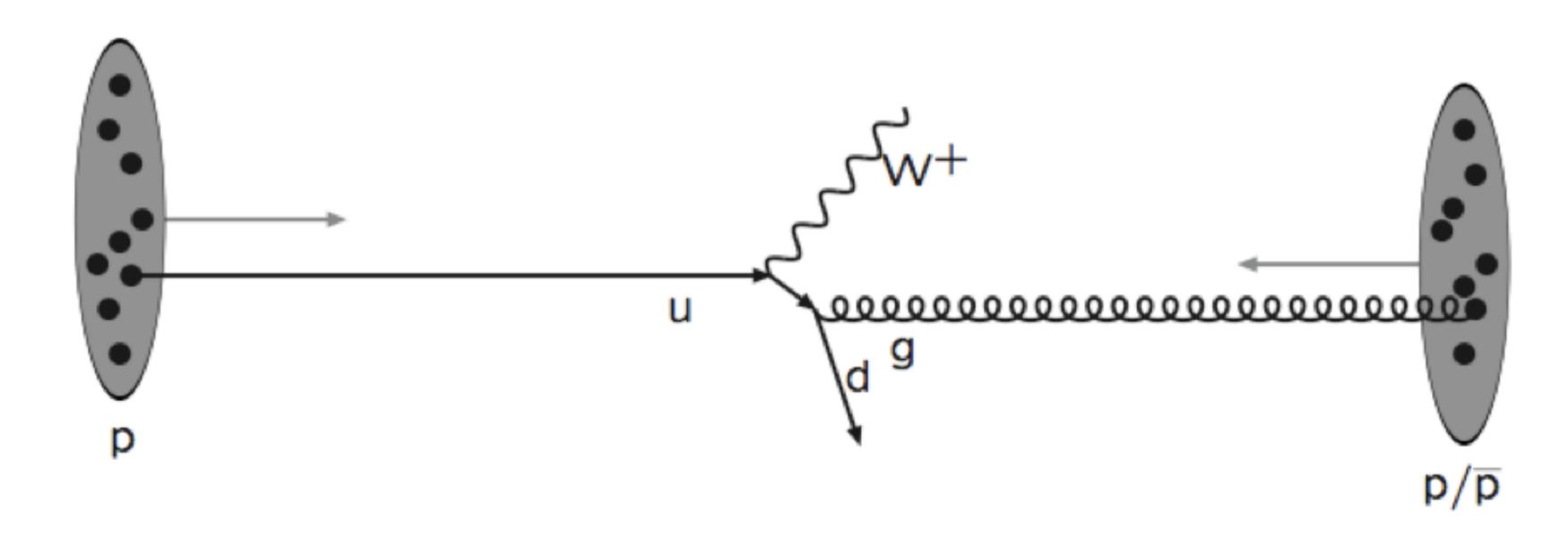
Initially two beam particles are coming in towards each other. Normally each particle is characterized by a set of parton distributions, which defines the partonic substructure in terms of flavour composition and energy sharing. This determines the energy of the interacting partons (x<sub>1</sub>, x<sub>2</sub>)



$$\begin{split} |\sigma(p(P_1)+p(P_2)\to Y) \;=\; \int_0^1 dx_1 \int_0^1 dx_2 \sum_f f_f(x_1) f_{\overline{f}}(x_2) \cdot \sigma(q_f(x_1P) + \overline{q}_f(x_2P) \to Y) \\ &\quad \text{partonic x-section:} \\ &\quad \text{phase space* matrix element} \end{split}$$

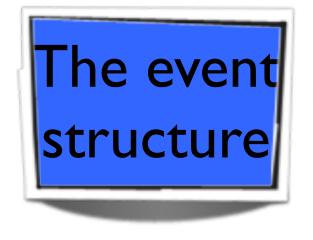


#### The Structure of an event

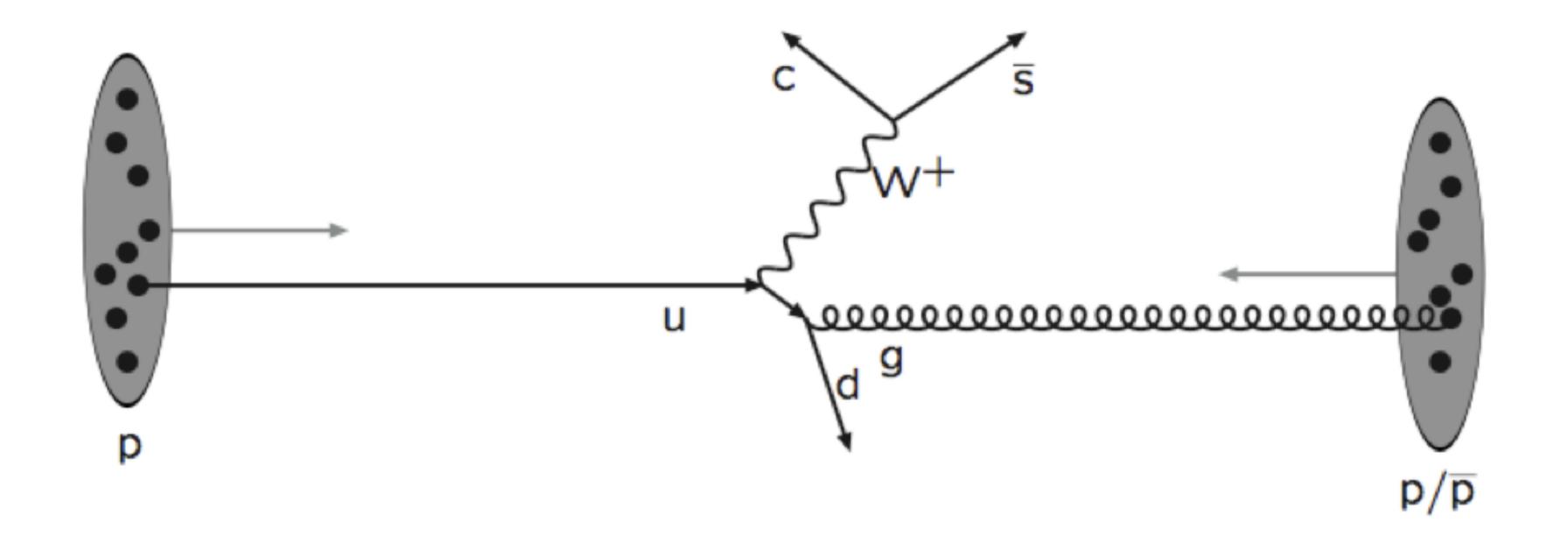


→ One incoming parton from each of the protons enters the hard process, where then a number of outgoing particles are produced. It is the nature of this process that determines the main characteristics of the event.

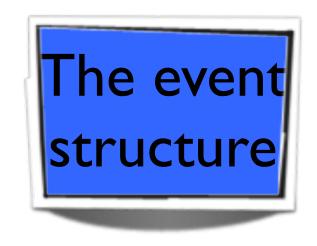
Hard subprocess: described by matrix elements



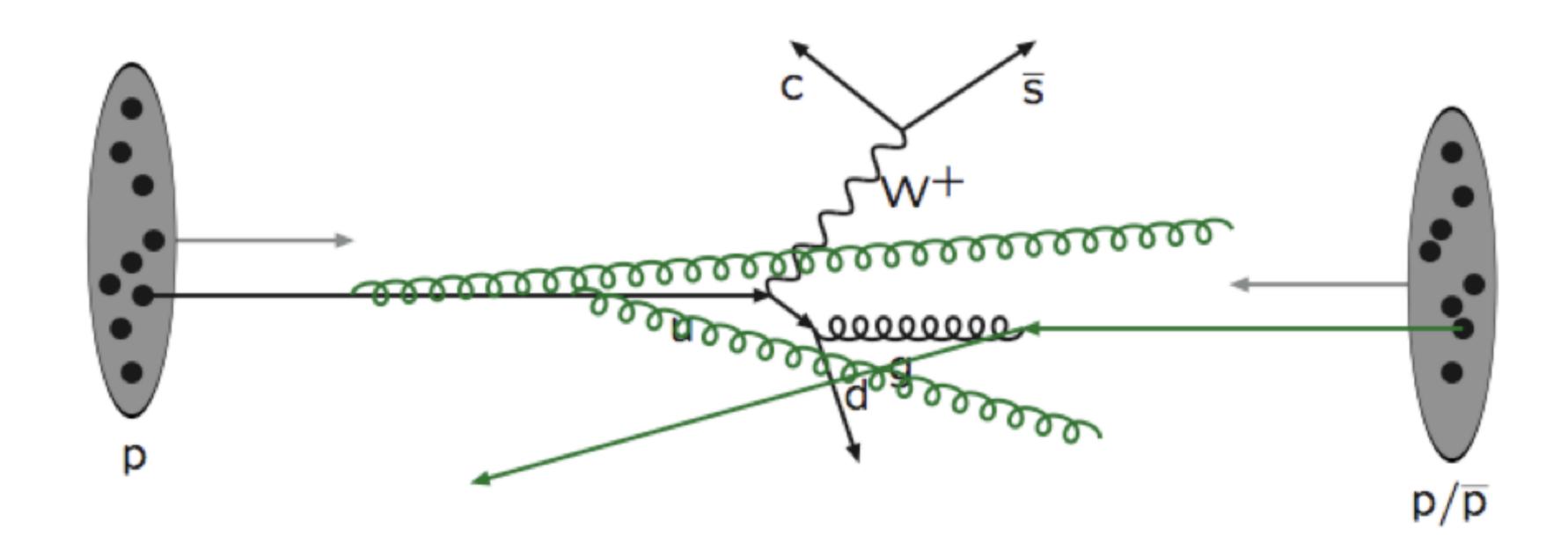
#### The Structure of an event: resonances



The hard process may produce a set of short-lived resonances, like the  $Z^0/W^{\pm}$  gauge bosons.

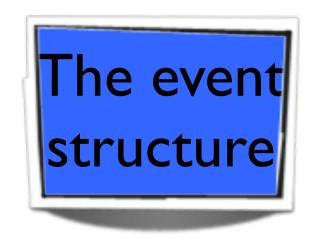


#### The Structure of an event: ISR

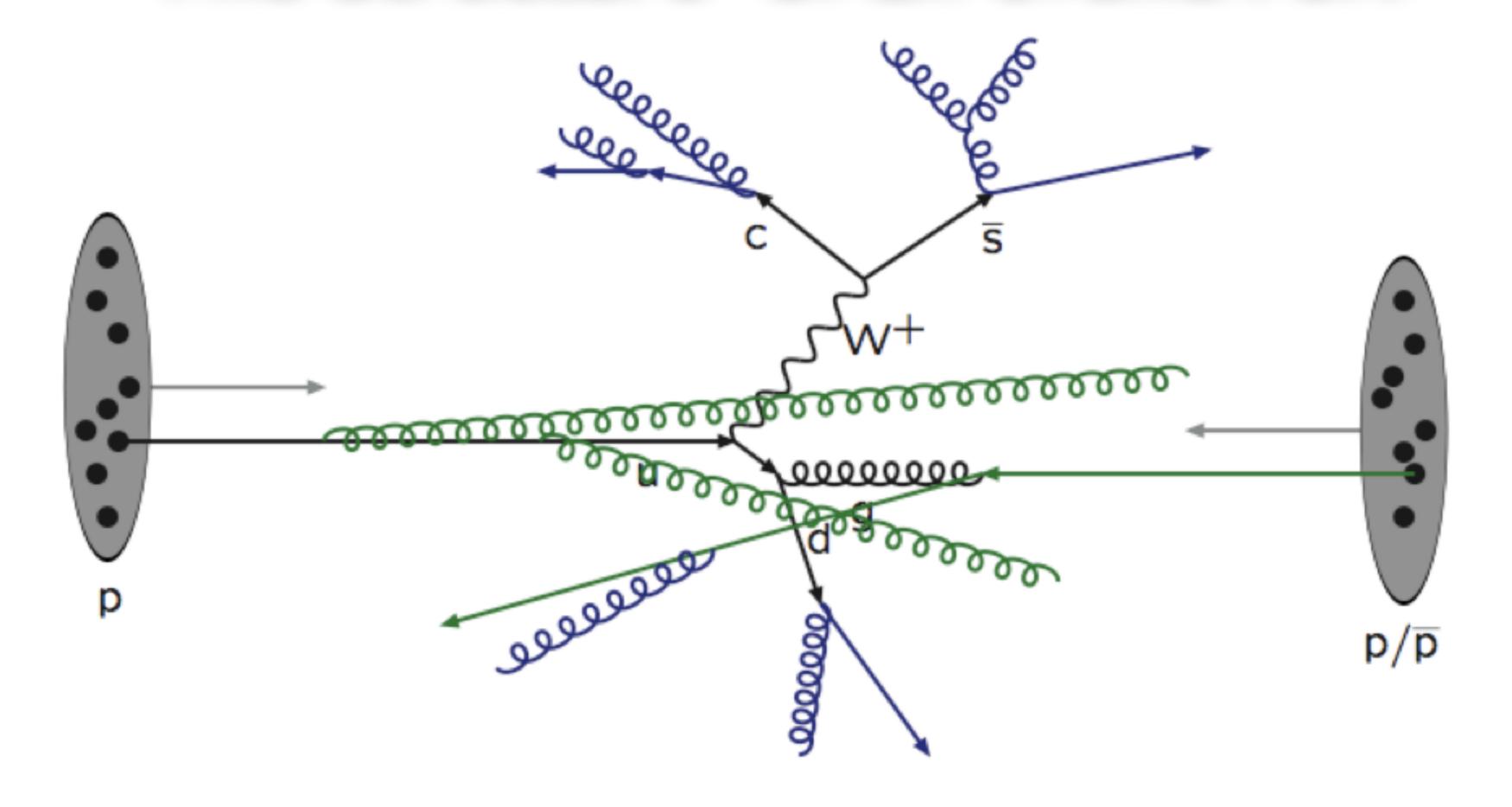


One shower initiator parton from each beam may start off a sequence of branchings, such as  $q \rightarrow qg$ , which build up an initial-state shower.

Initial-state radiation: spacelike parton showers

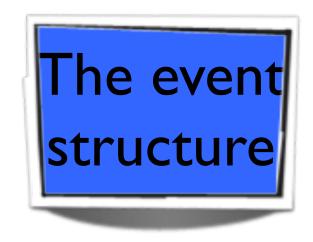


#### The Structure of an event: FSR

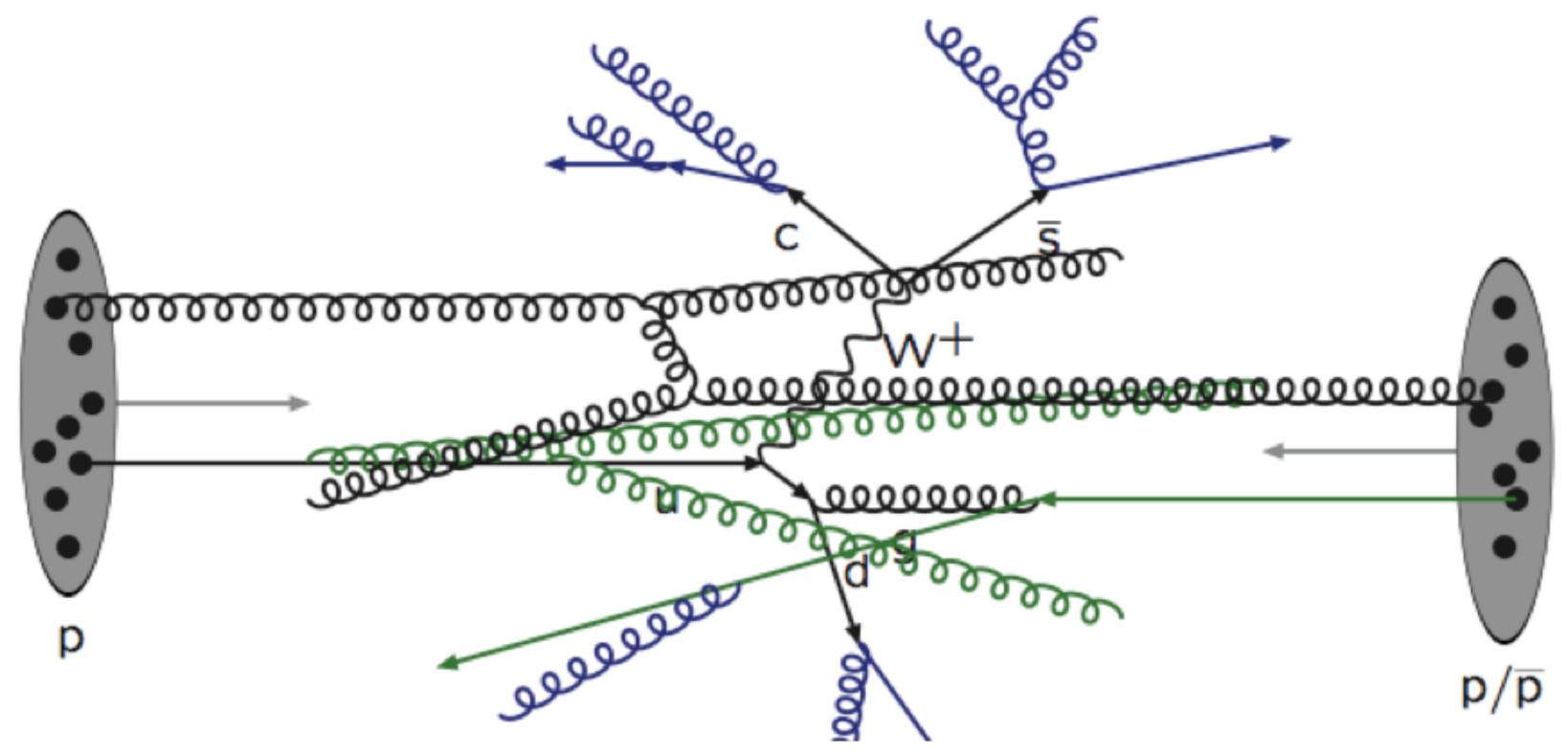


The outgoing partons may branch, just like the incoming did, to build up final-state showers.

Final-state radiation: timelike parton showers



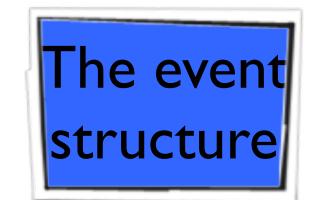
## The Structure of an event: pile-up



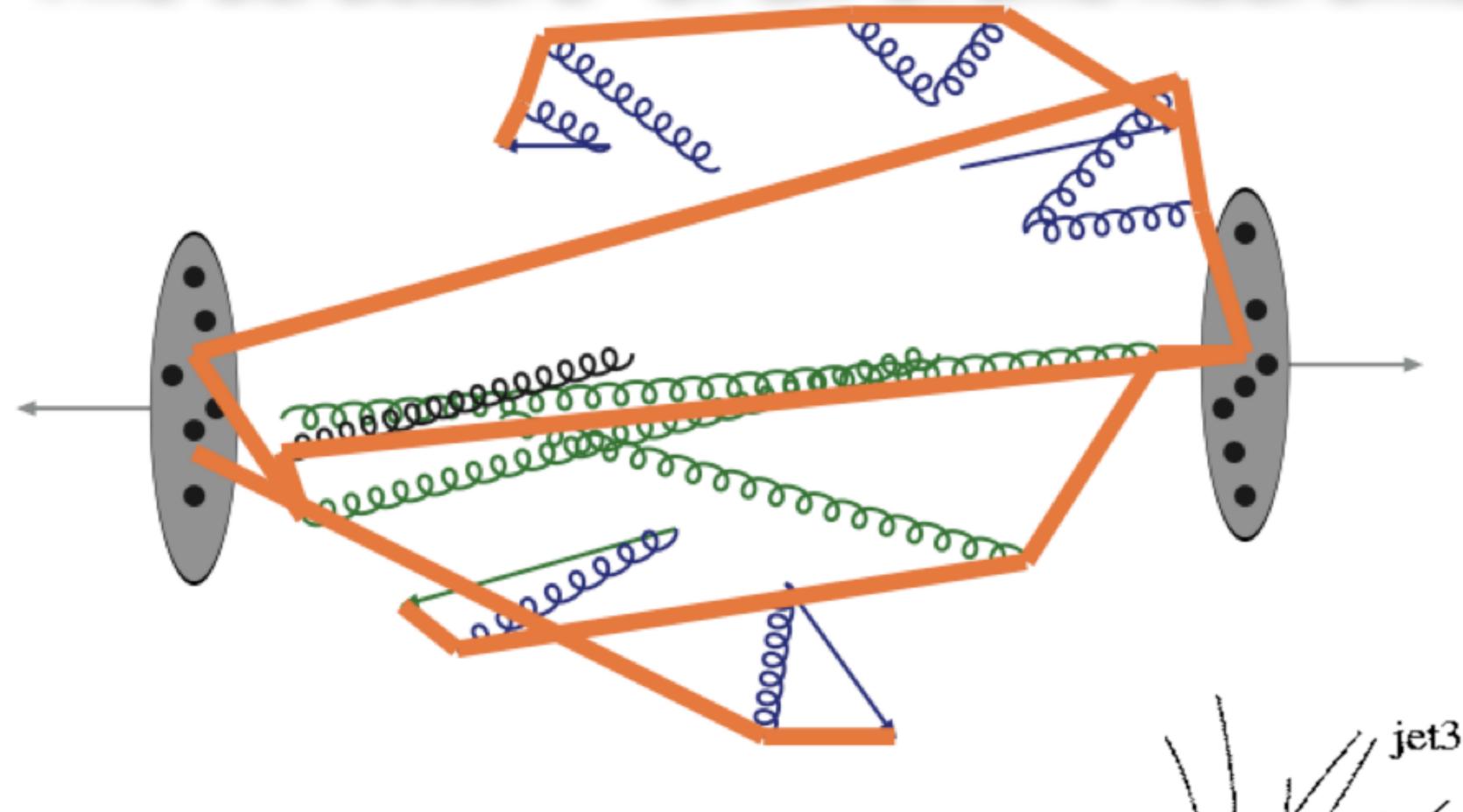
In addition to the hard process, further semihard interactions may occur between the other partons of two incoming hadrons.

There is in time pile-up which comes from the same bunch of protons from the interaction of interest, and can be resolved by setting the interaction points location by identifying vertices.

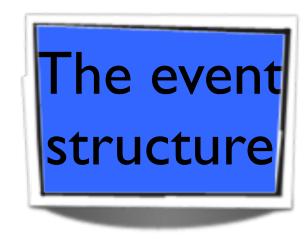
The second type is out-time pile-up, which which comes from other proton bunches when the detector has not yet recorded the signal completely due to dead time, the time needed for a certain detector to be able to record an event after a previous one



#### The event The Structure of an event: hadronisation



The result of the hadronization is that quark and gluons are not observed as free particles but as Hadrons, and actually in the detector as jets of particles in a narrow cone



# Pile-up

2010

O(2) Pile-up events

150 ns inter-bunch spacing

2011

O(10) Pile-up events

50 ns inter-bunch spacing

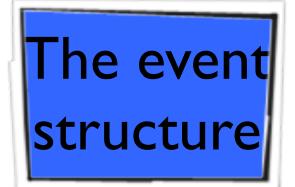
Design value (expected to be reached at L=10<sup>34</sup>!)

2012

O(20) Pile-up events

50 ns inter-bunch spacing





#### the cross-section

#### 

$$1 \text{ barn} = 10^{-28} \text{ m}^2 = 10^{-24} \text{ cm}^2$$

#### The Cross-section

Number of observed events is proportional to

- 1) Luminosity
- 2) analysis efficiency
- 3) cross section of the process

$$N_{obs} = \int Ldt \cdot \varepsilon \cdot \sigma$$

The luminosity is a parameter of the LHC and can be increased

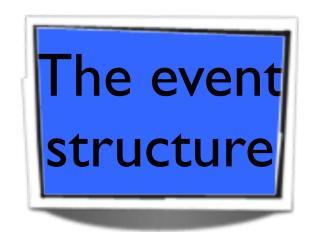
$$L = \frac{f_{rev} n_{bunch} N_p^2}{4 \pi \sigma_v \sigma_v}$$

revolving frequency: f<sub>rev</sub>=11245.5/s

#bunches:  $n_{bunch}$ =2808

#protons / bunch:  $N_p = 1.15 \times 10^{11}$ 

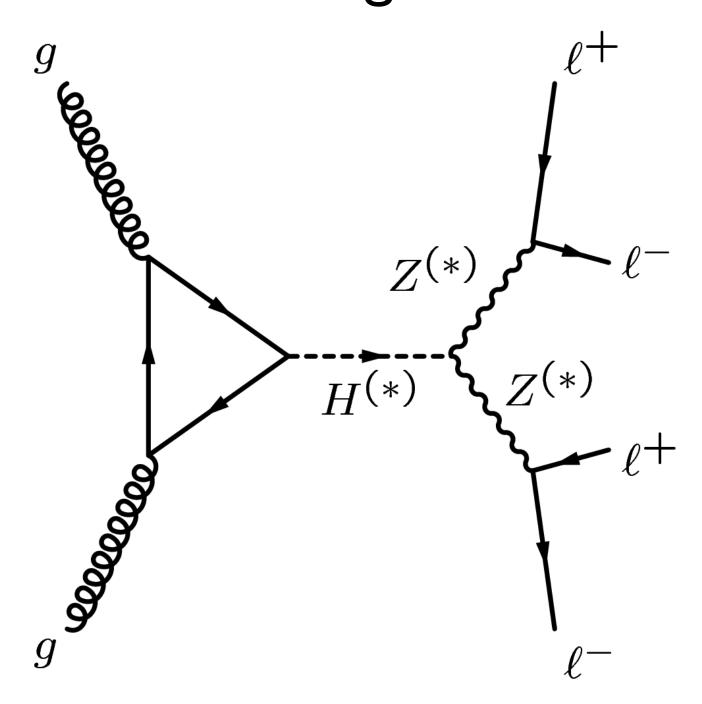
Area of beams:  $4\pi\sigma_x\sigma_v\sim40~\mu m$ 



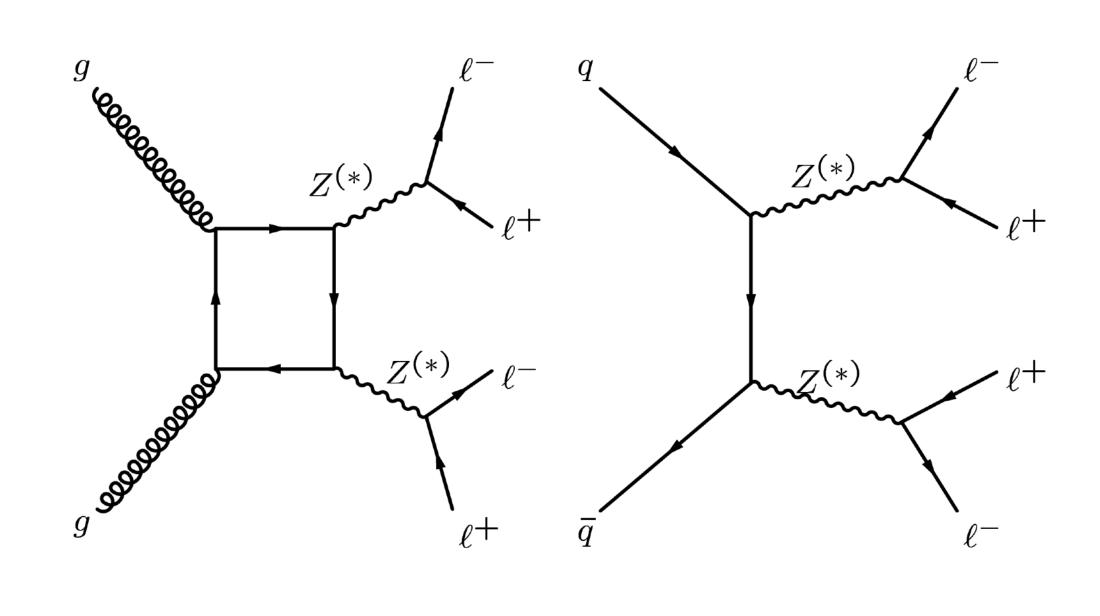
# Signal and backgrounds

Other processes (background) can mimic signal final state. Same particles in the final state!

#### HIGGS Signal

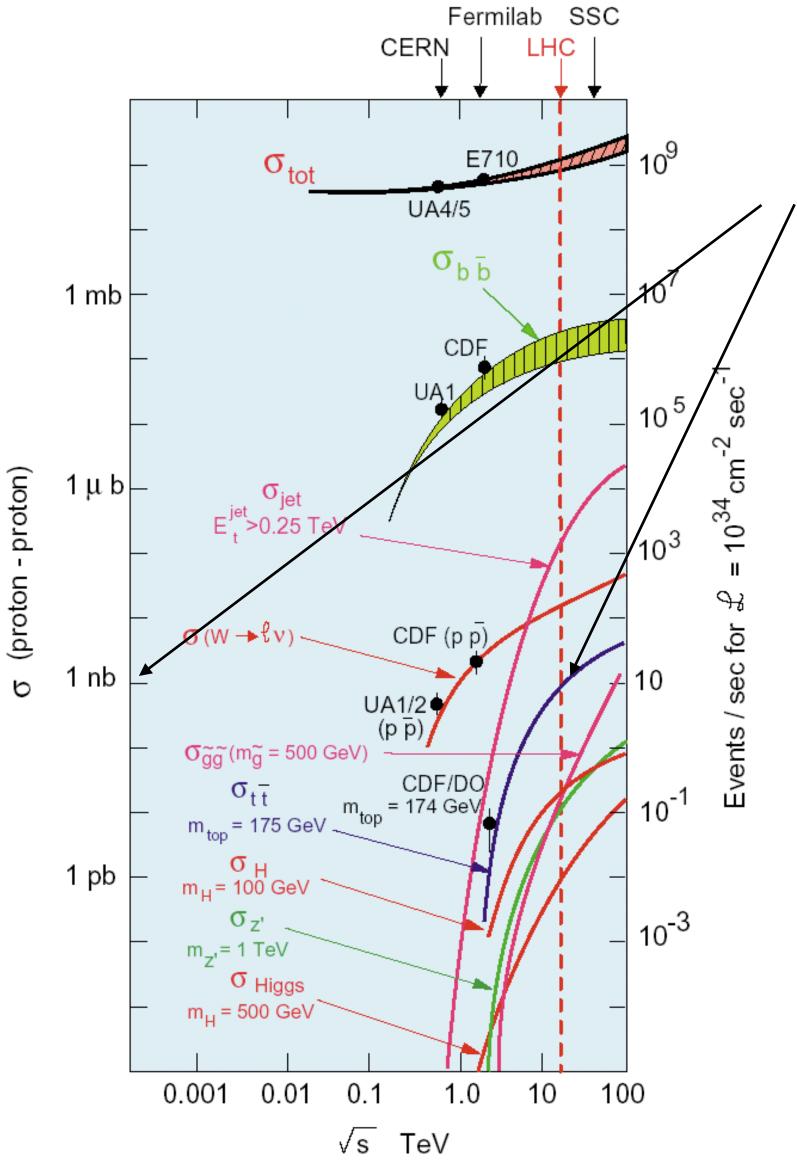


#### SM backgrounds





#### Cross-sections/ number of Events

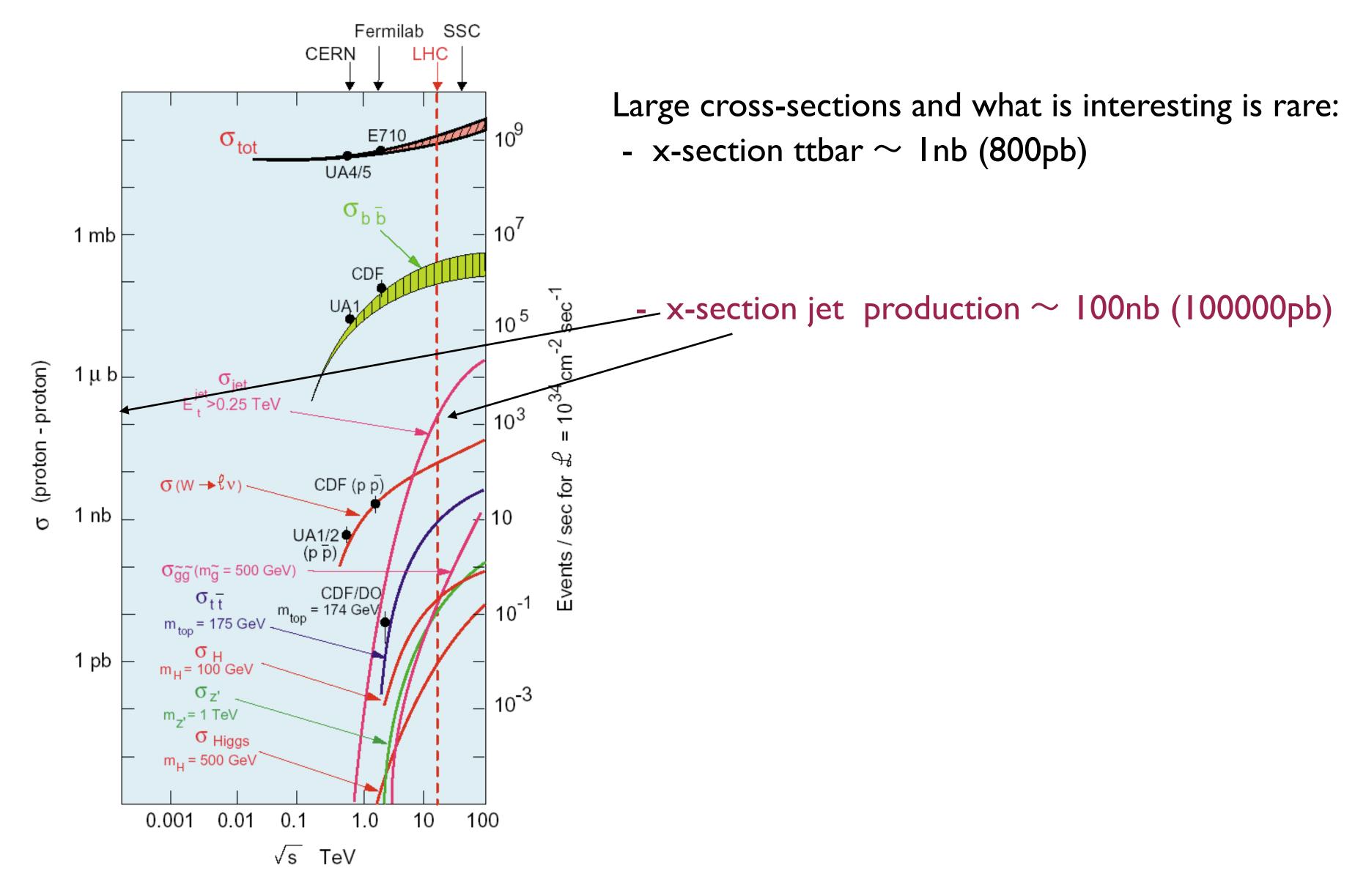


Large cross-sections and what is interesting is rare:

- x-section ttbar ~ Inb (800pb)

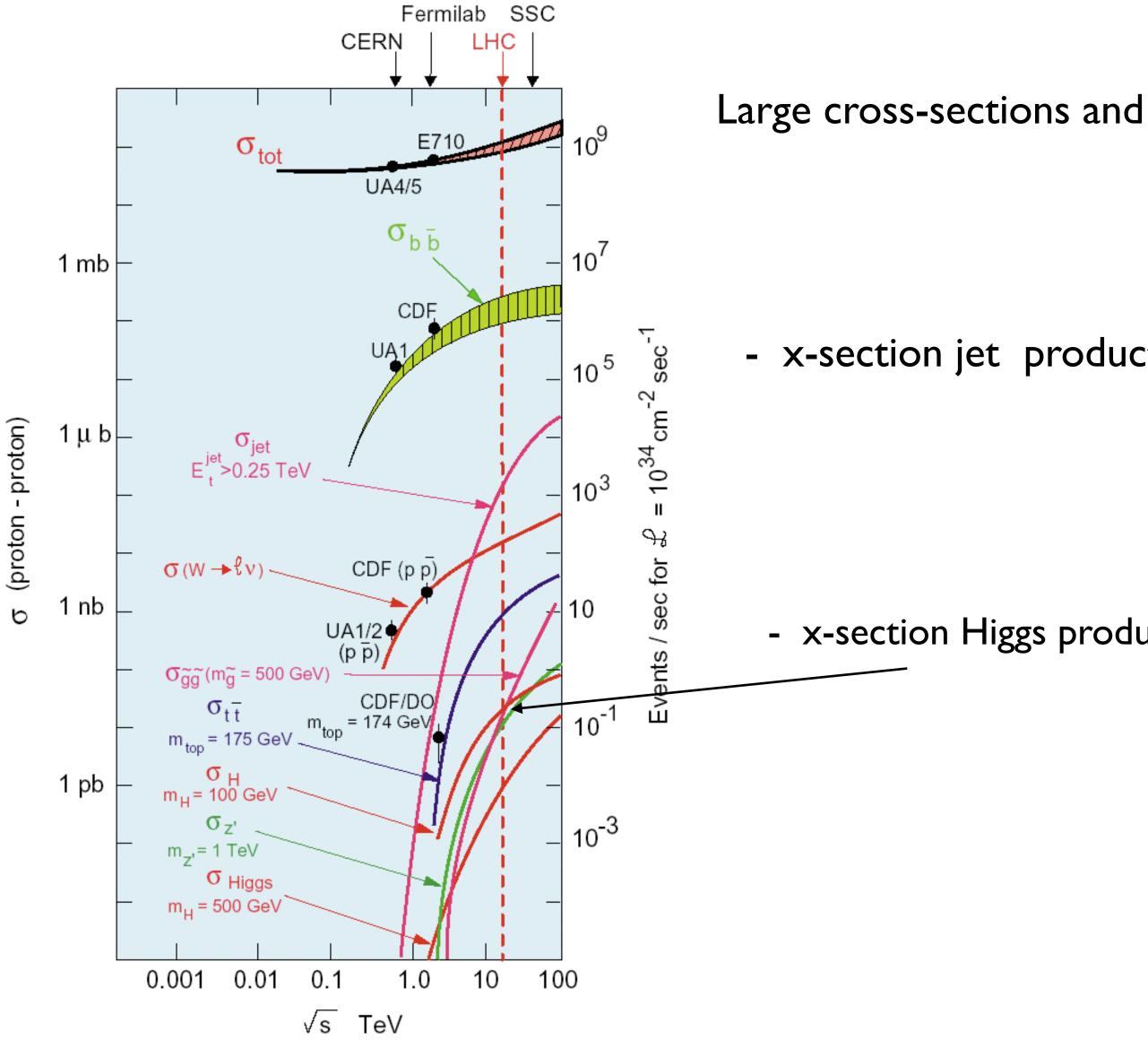


#### Cross-sections/ number of Events





#### Cross-sections/ number of Events

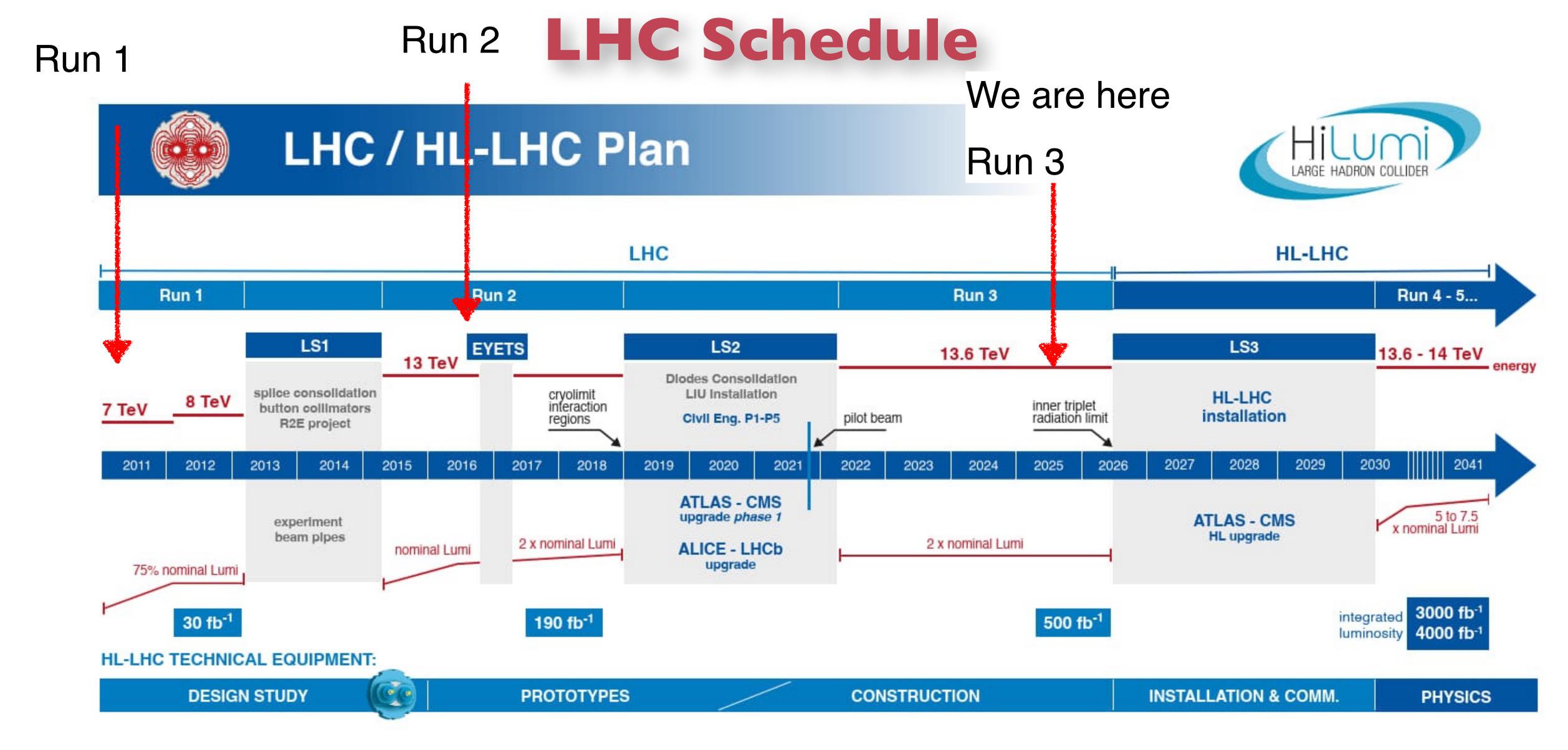


Large cross-sections and what is interesting is rare:

- x-section jet production  $\sim$  100nb (100000pb)

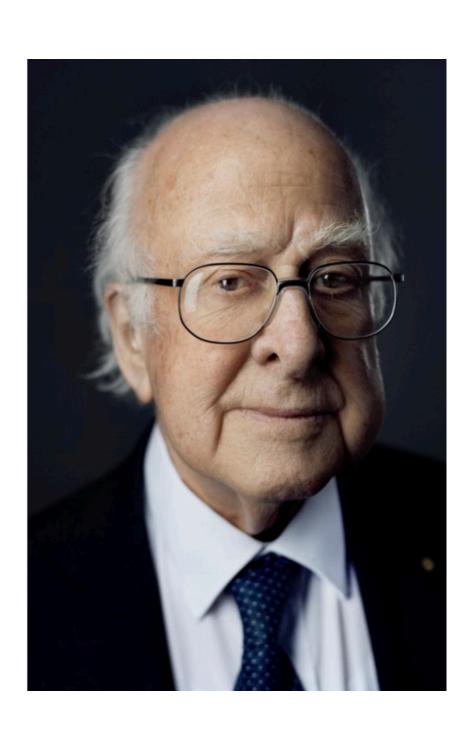
- x-section Higgs production ~ 10pb





The luminosity (L) is a parameter of the accelerator, related to how many collisions one gets and is measured in fb-1

# Single Higgs



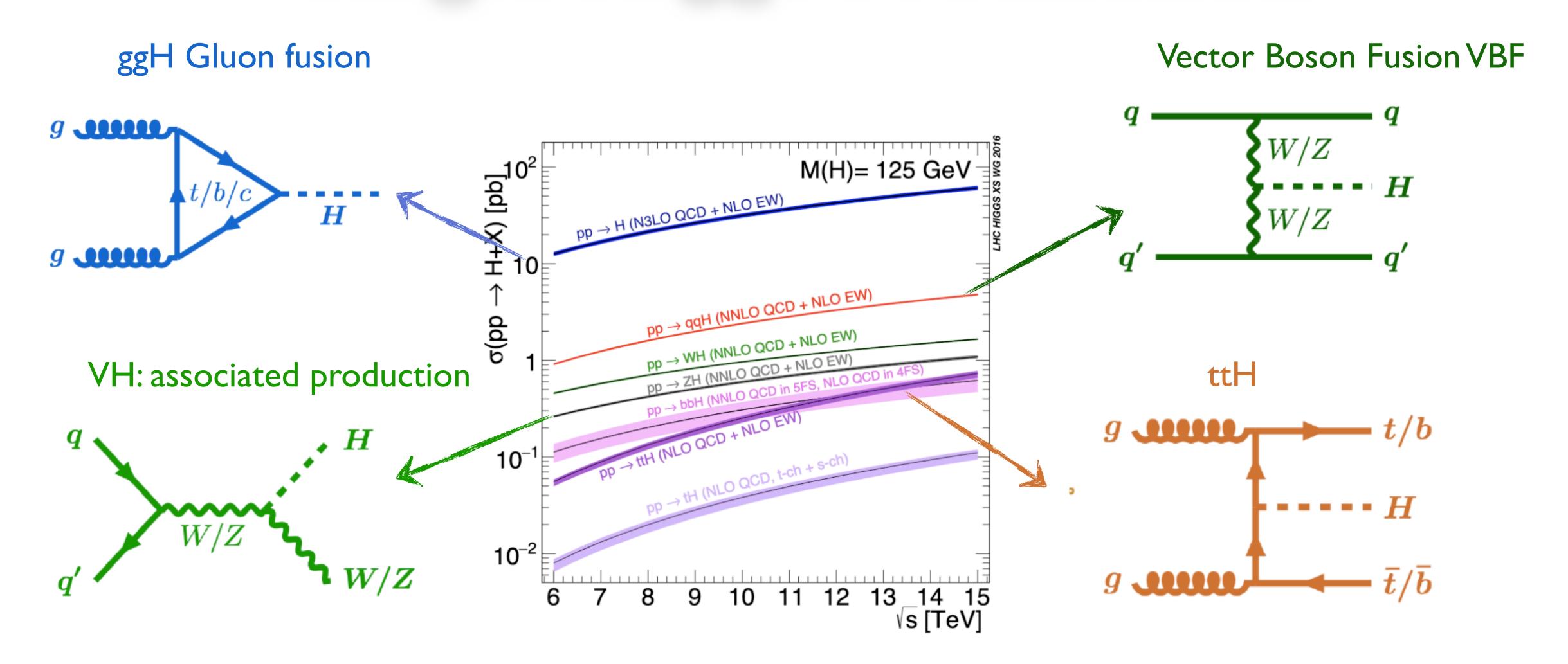
# Single Higgs



#### and François Englert & Robert Brout

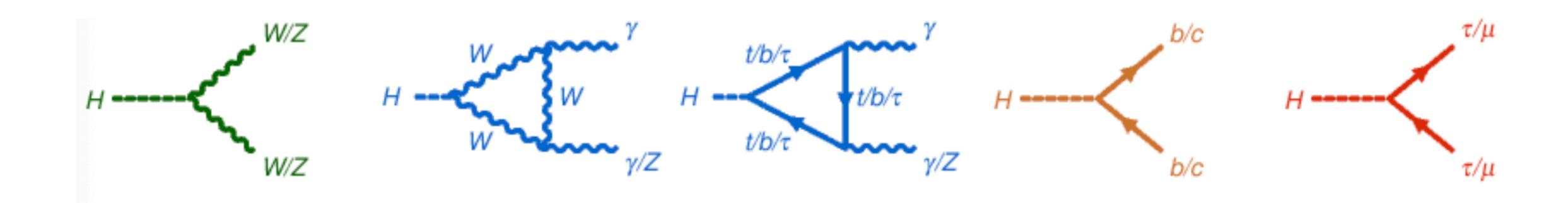


# Single Higgs Production



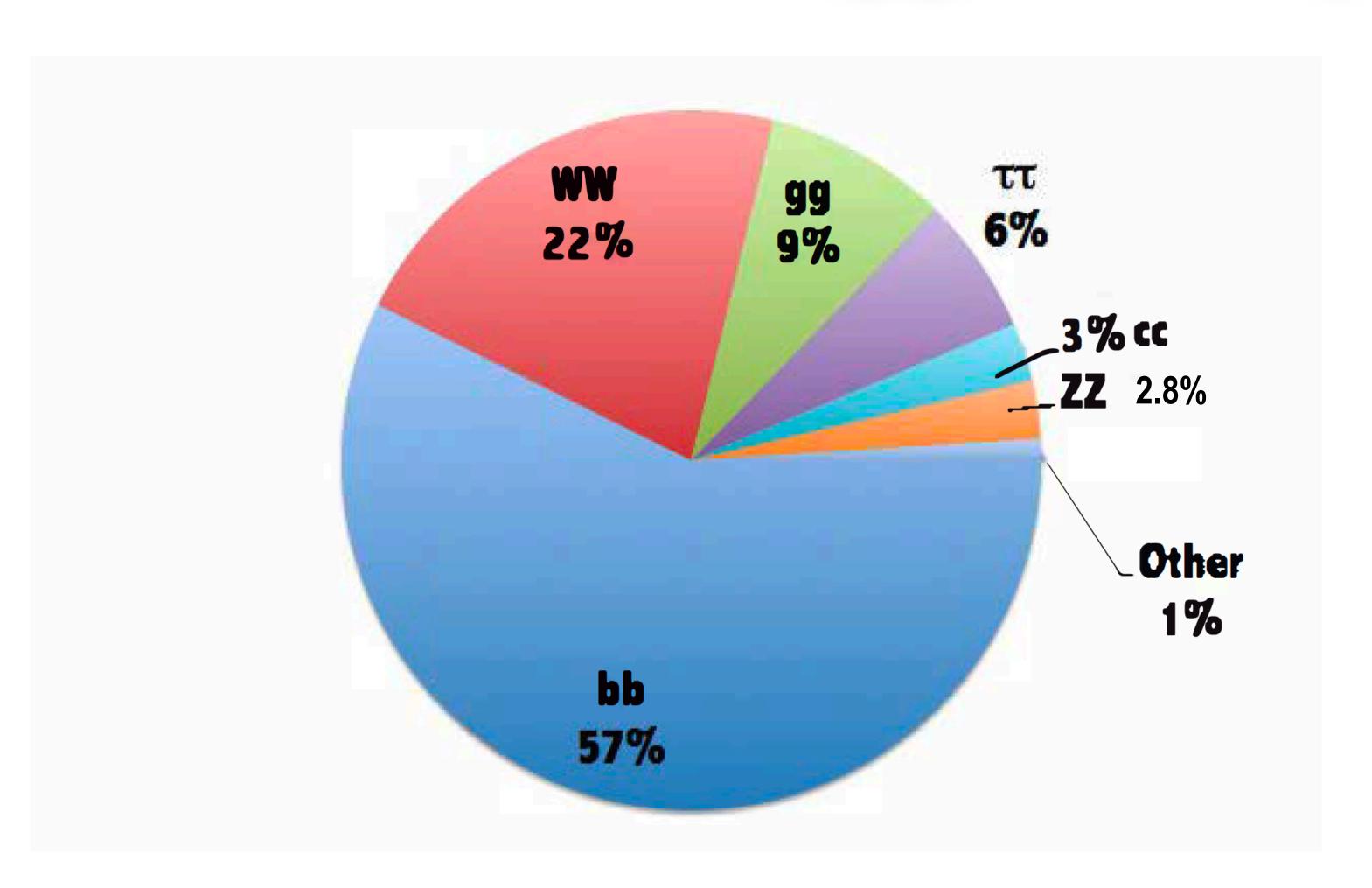
Higgs bosons are produced by a variety of mechanisms, the most important of which is gluon-gluon fusion (ggF).

# Higgs decays



- •Once produced, the Higgs boson can decay into a variety of final states.
- •By analyzing rates in each production&decay channel, we can determine Higgs couplings.

# Higgs decays



	BR(%)
bb	57
WW	22
TT	6.2
ZZ	2.8
γγ	0.23
<b>Z</b> γ	0.15
μμ	0.02

- •Once produced, the Higgs boson can decay into a variety of final states.
- •By analyzing rates in each production&decay channel, we can determine Higgs couplings.

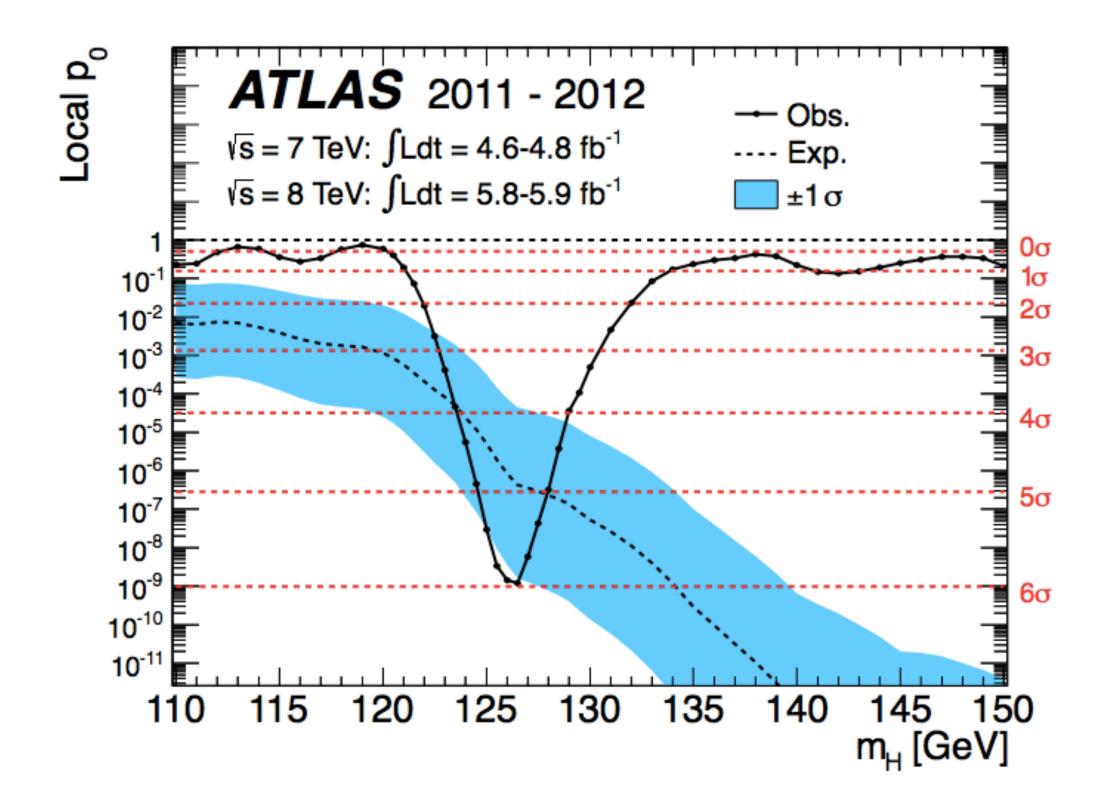
## Higgs discovery in run l

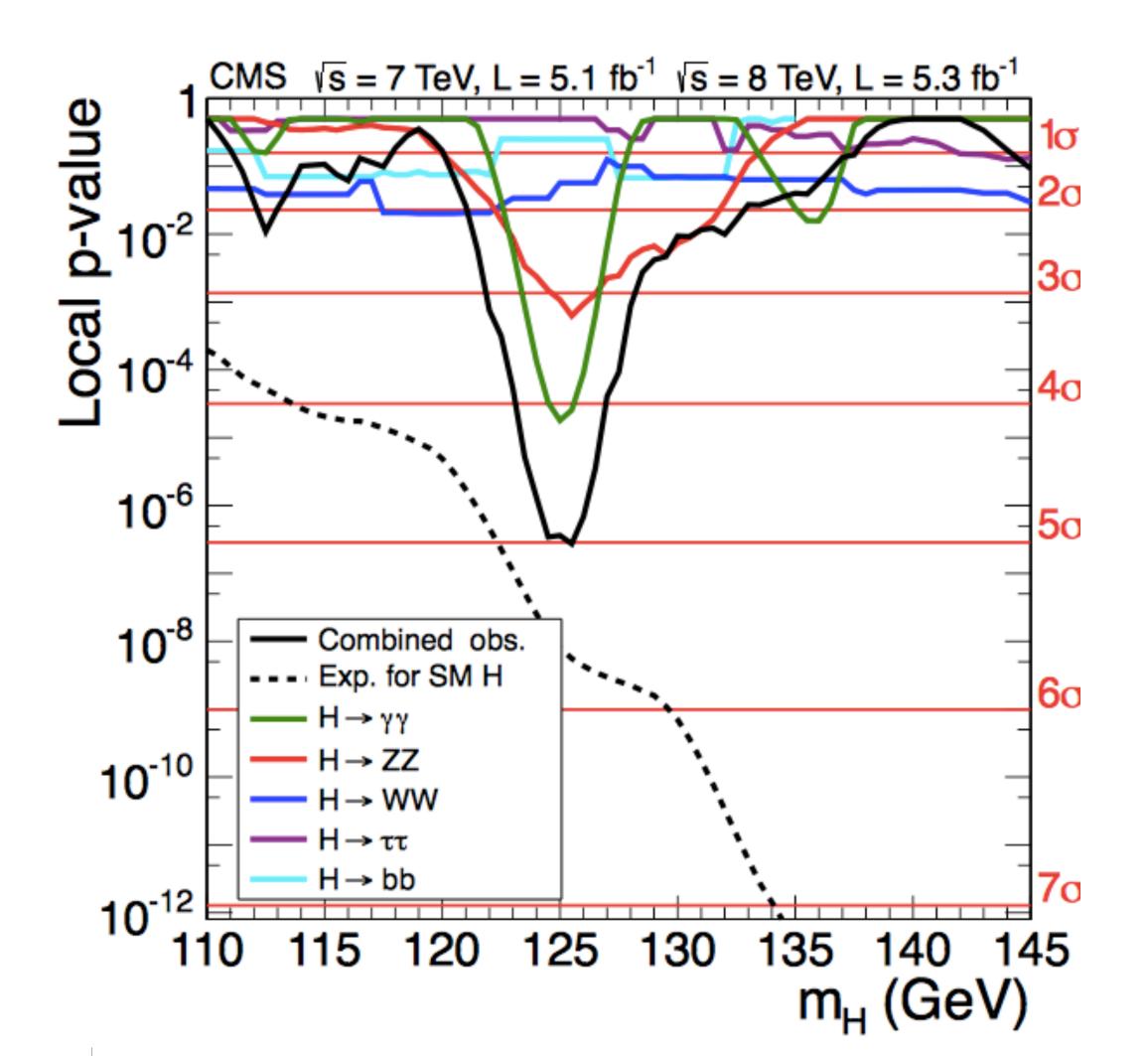
A scalar boson compatible with the SM Higgs has been discovered in run I as shown by the combination of ATLAS and CMS run I results

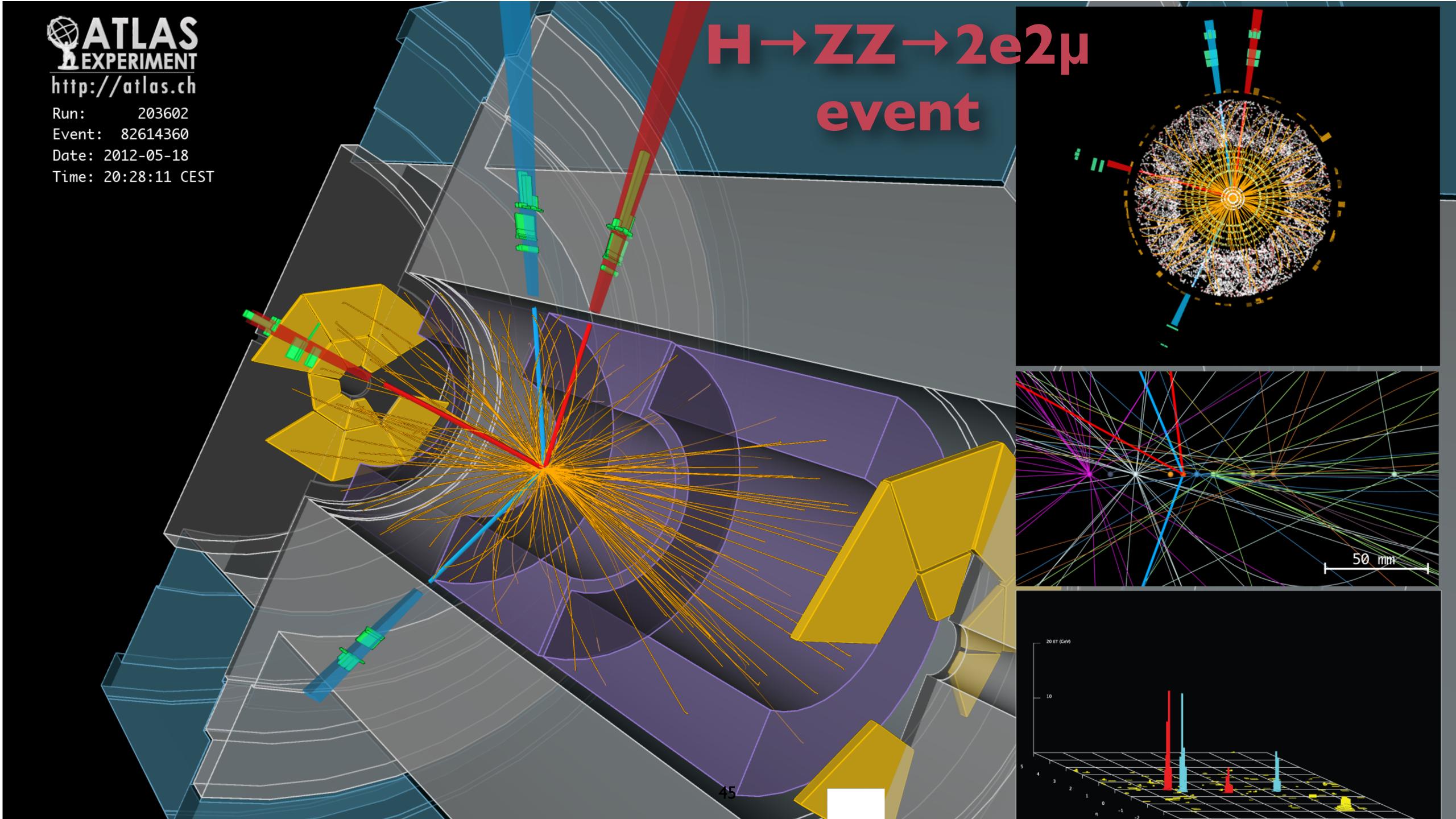
#### Greatest achievement of run I

concentrated effort on its properties:

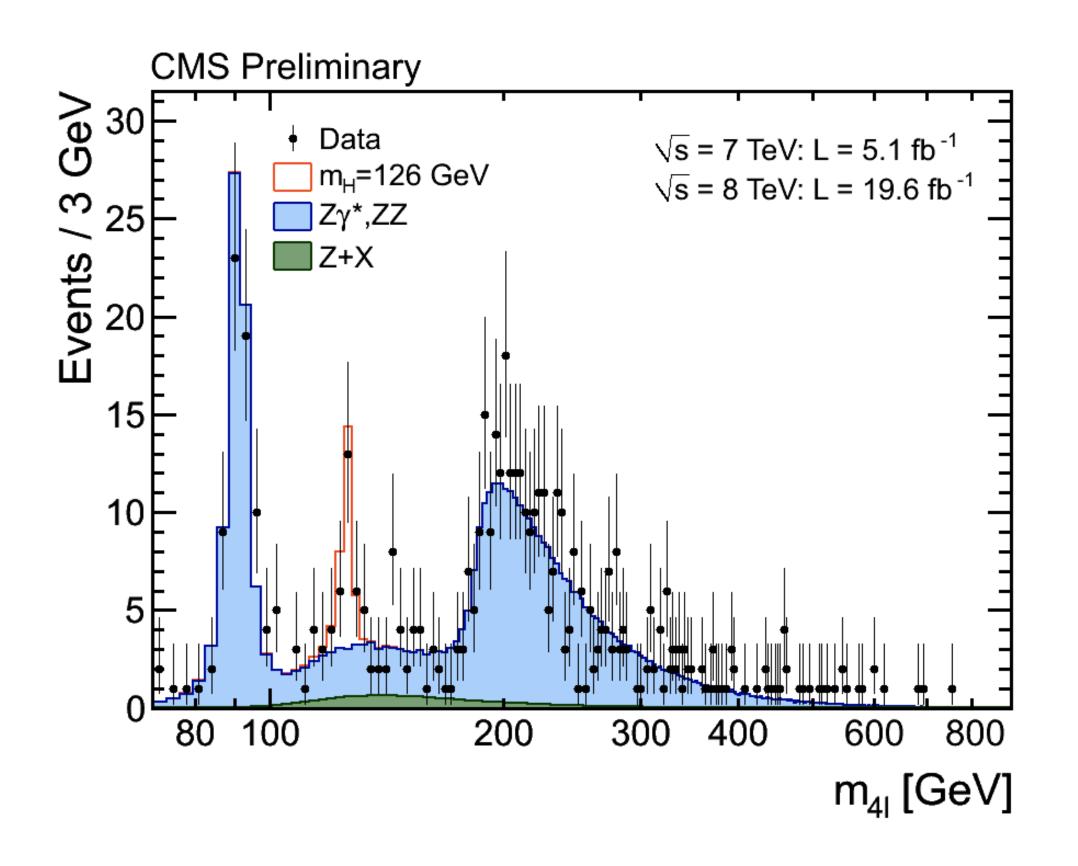
- magnitude of couplings
- mass measurements
- spin/CP

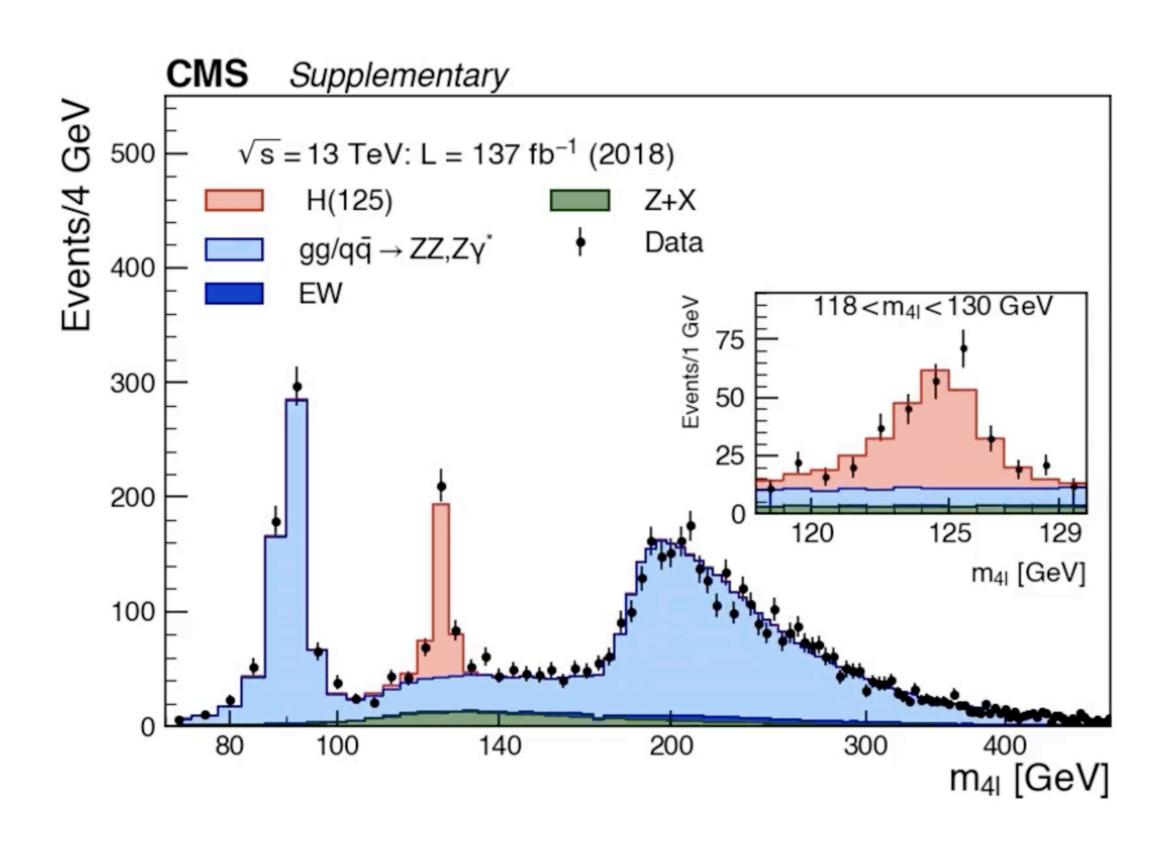






## Run I discovery Run2 precision





Run 3 trigger detector and analysis improvements!

### Object reconstruction is fundamental

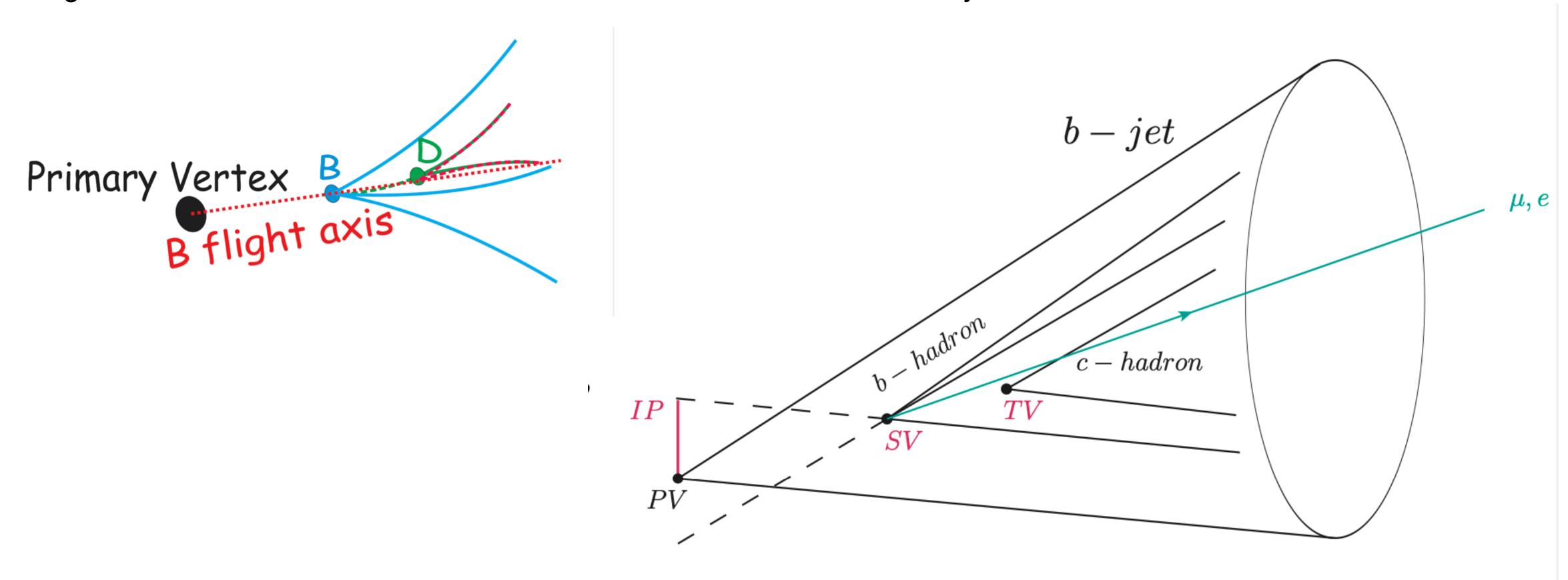
We have to reconstruct the particles in which the Higgs decays in the best possible way. This really pays off, as you will see in these lectures.

Fundamental are the reconstructions of

- leptons (muons, taus, electrons)
- photons
- b-quarks and c-quarks (flavor tagging)
- Boosted objects

## Flavor tagging

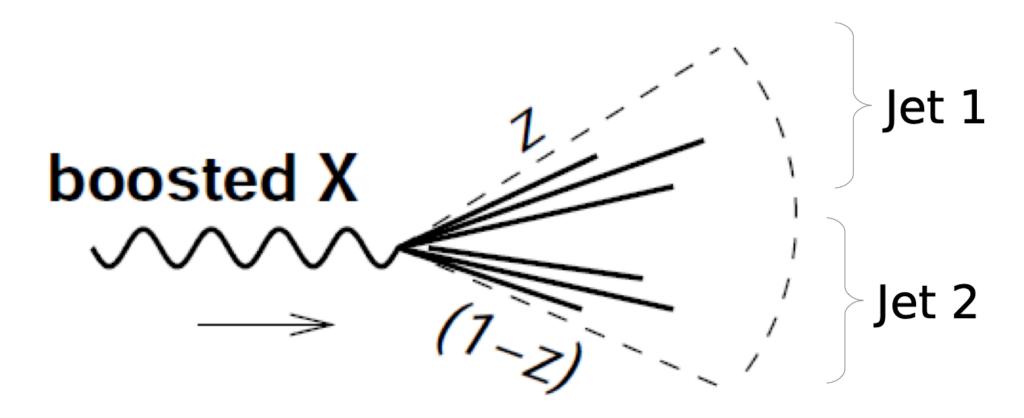
We tag b-hadrons and c-hadrons thanks to the fact that there is a secondary vertex



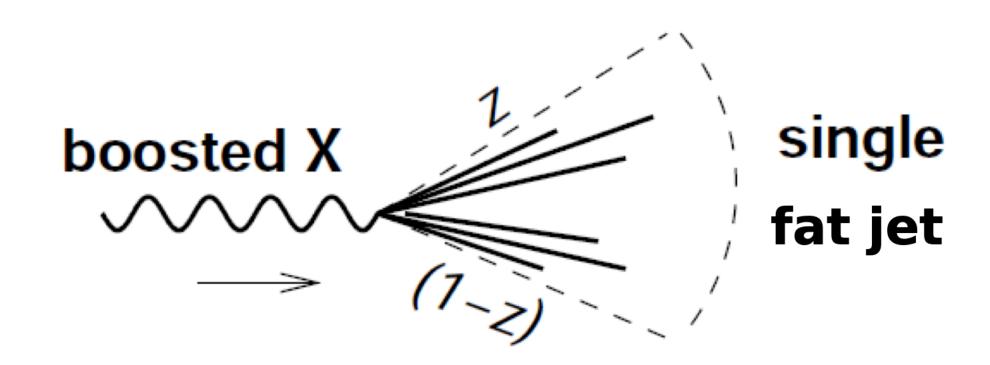
#### Boosted objects

At the LHC given the large center of mass energy and given that the SM particles have masses below 200 GeV, also the heaviest SM particles often acquire large momentum >> m → production of "boosted objects"

Normally we reconstruct jets with R=0.4, if the object is boosted the jets in which it decays cannot be resolved in small r-jets

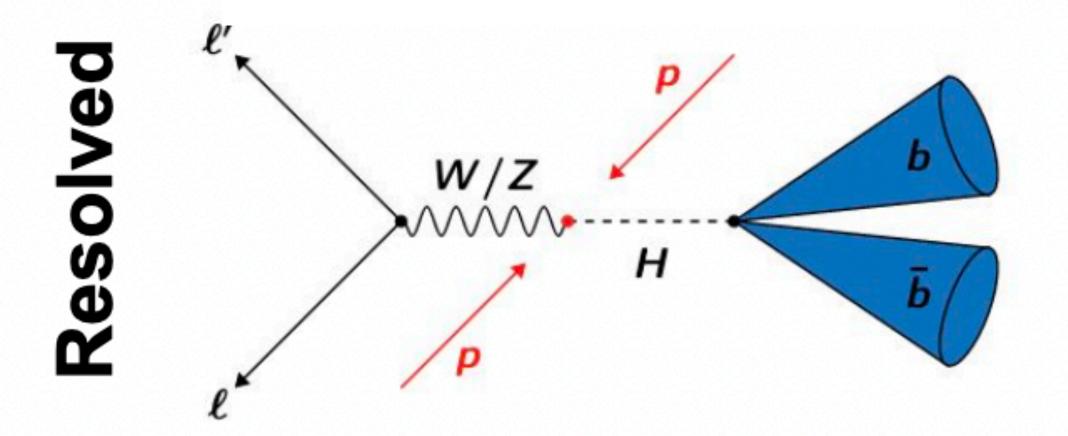


Recover sensitivity to boosted objects by developing boosted taggers, using larger R

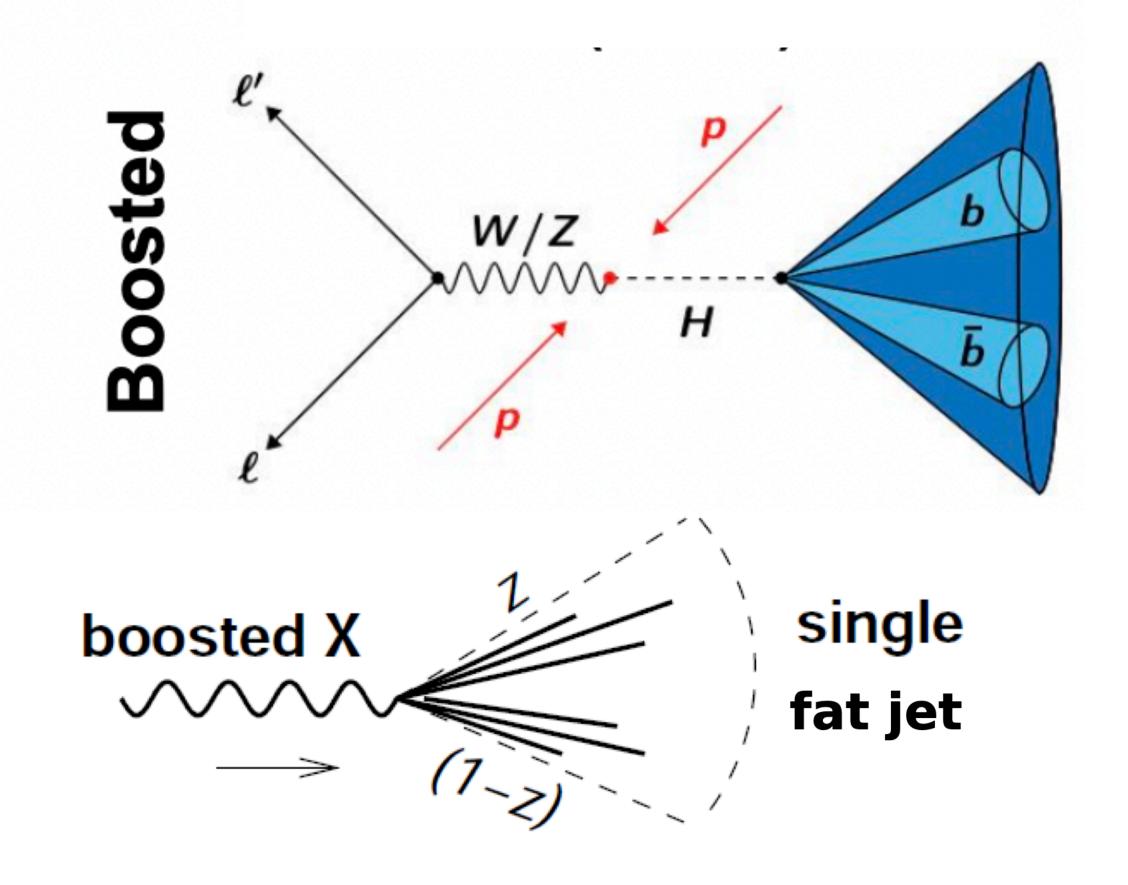


#### Boosted objects

At the LHC given the large center of mass energy and given that the SM particles have masses below 200 GeV, also the heaviest heaviest SM particles often acquire pT >> m → production of "boosted objects"



Recover sensitivity to boosted objects by developing boosted taggers, using larger R



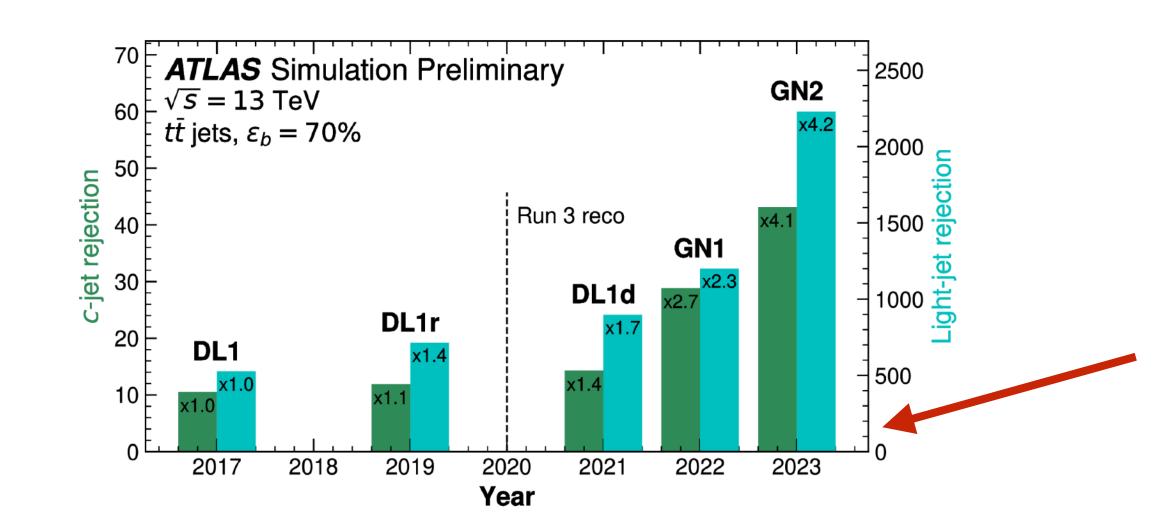


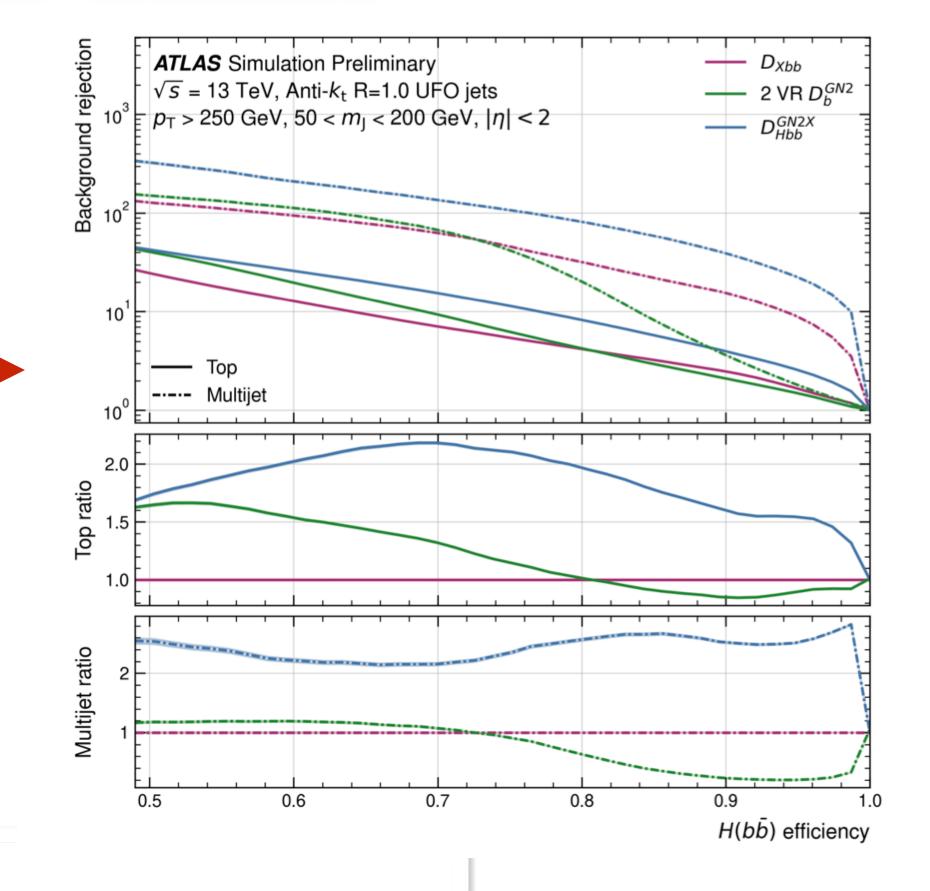
#### Flavor tagging in continuous evolution

#### Boosted H->bb/cc tagging ATL-PHYS-PUB-2023-021

• <u>Boosted b-tagging:</u> new algorithm, GN2X for largeradius jets: tagging boosted H(bb) jets and H(cc) jets.

#### small R-jet tagging Jet Flavour Tagging With GN1 and DL1d





 GN2X benefits from advances in flavour tagging of small-radius jets with Graph Neural Networks (GNNs)

# Higgs terminology

#### We talk of signal strength, defined as the $\mu$ parameter:

- $\mu$  is the ratio of the measured cross-section with respect to the SM expectation.
- $\mu$ =1 means that we measure back the SM

#### We also have introduced kappa's k-framework:

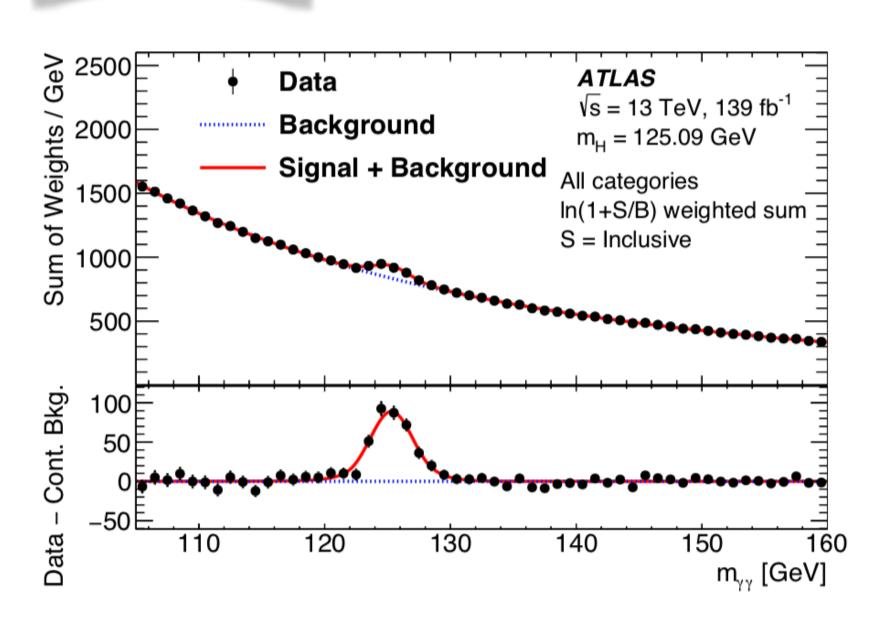
- •k is the ratio of the a Higgs coupling with respect to the coupling from SM expectation.
- •k=I means that we measure back the SM





10 % precision

$$\mu = 1.04^{+0.10}_{-0.09} = 1.04 \pm 0.06 \text{ (stat.)}^{+0.06}_{-0.05} \text{ (theory syst.)}^{+0.05}_{-0.04} \text{ (exp. syst.)}.$$



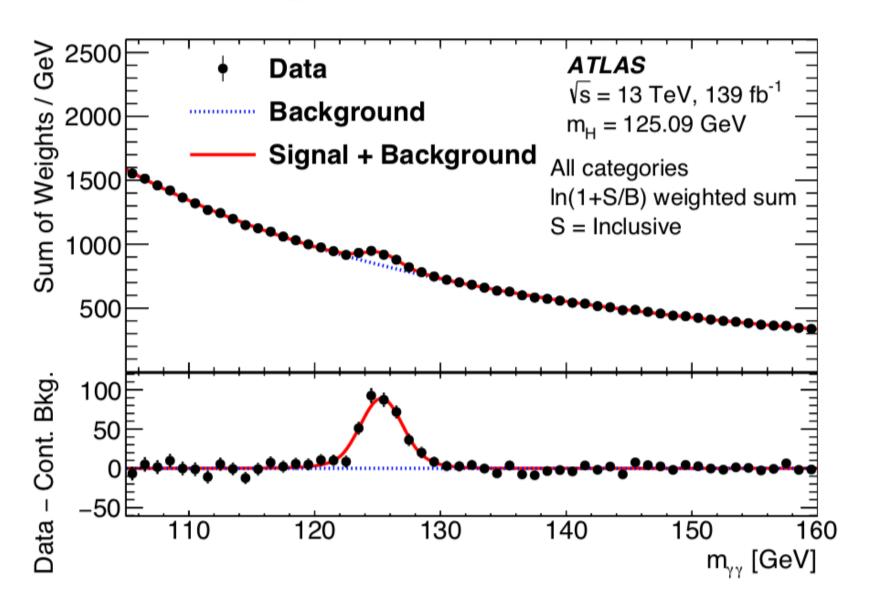
Please note that theoretical error is at level of other errors!

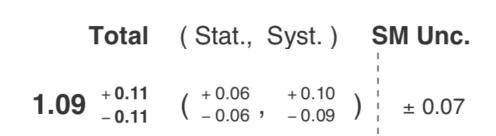


#### $H \rightarrow \gamma \gamma$

10 % precision

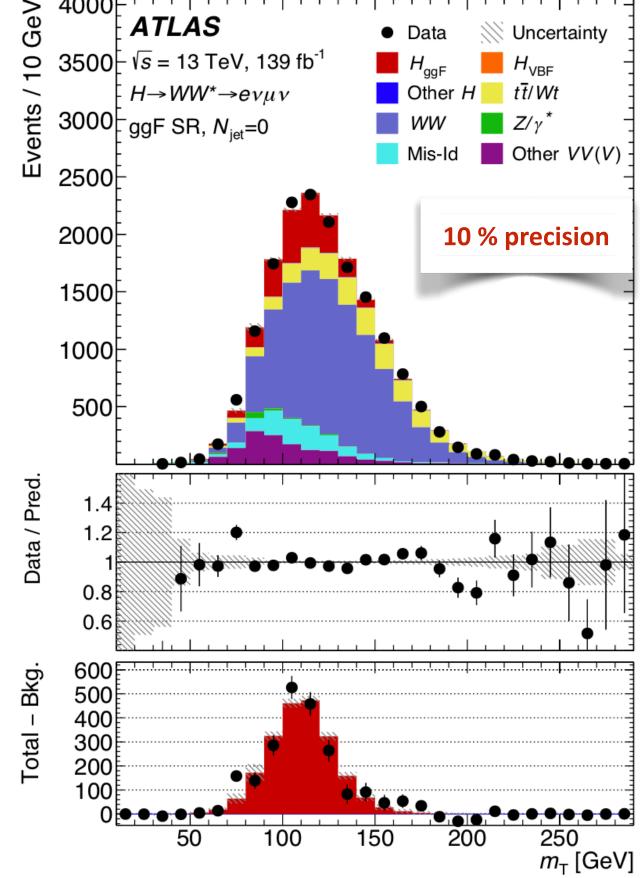
$$\mu = 1.04^{+0.10}_{-0.09} = 1.04 \pm 0.06 \text{ (stat.)}^{+0.06}_{-0.05} \text{ (theory syst.)} ^{+0.05}_{-0.04} \text{ (exp. syst.)}.$$











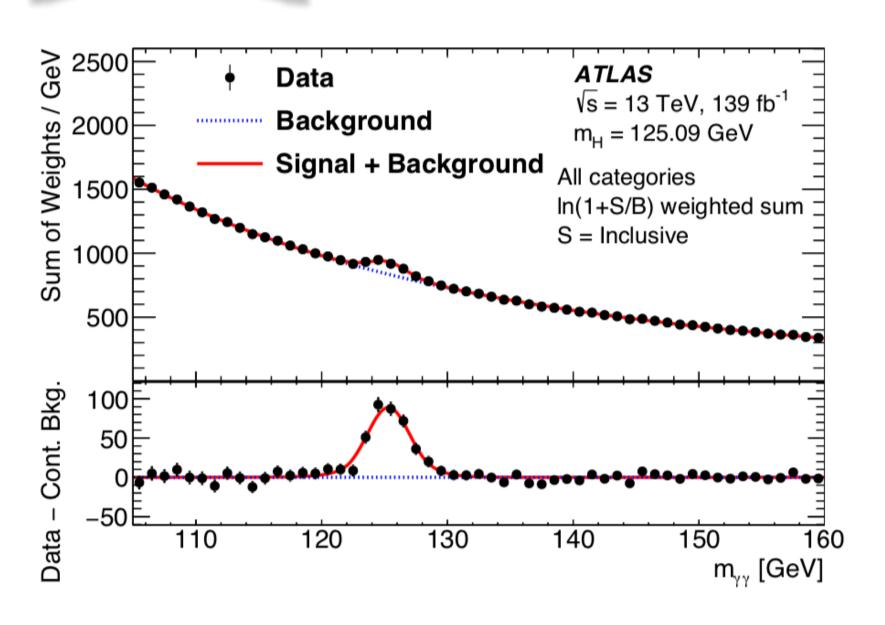
Please note that theoretical error is at level of other errors!

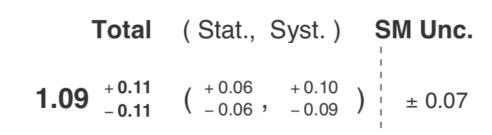


## $H \rightarrow \gamma \gamma$

10 % precision

$$\mu = 1.04^{+0.10}_{-0.09} = 1.04 \pm 0.06 \text{ (stat.)}^{+0.06}_{-0.05} \text{ (theory syst.)}^{+0.05}_{-0.04} \text{ (exp. syst.)}.$$

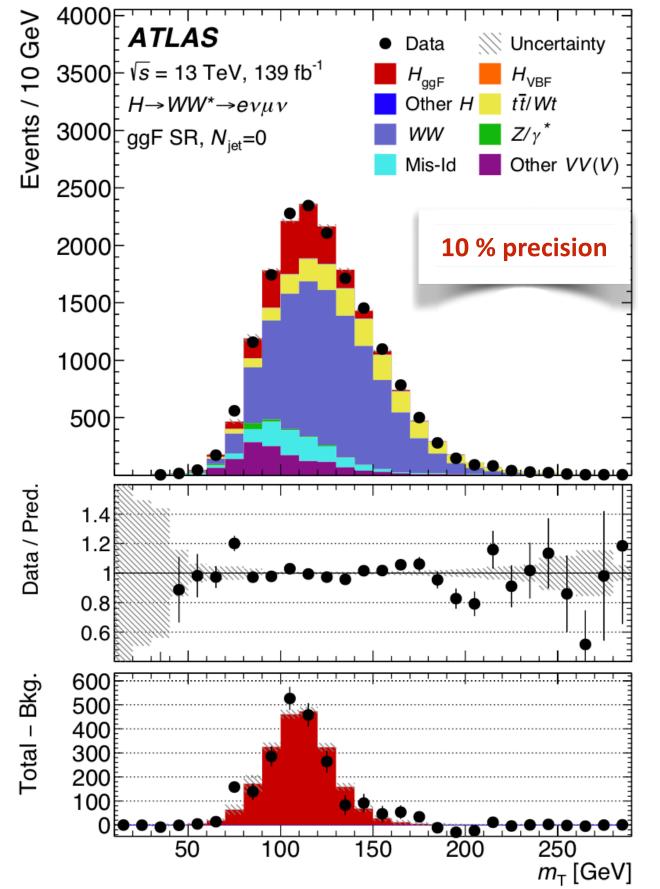


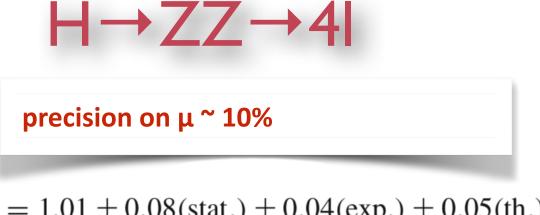




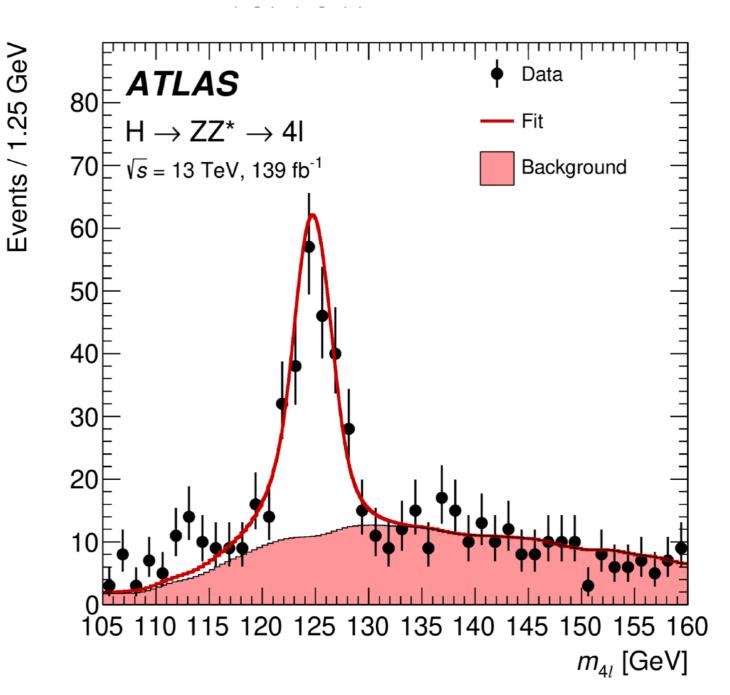
Please note that theoretical error is at level of other errors!



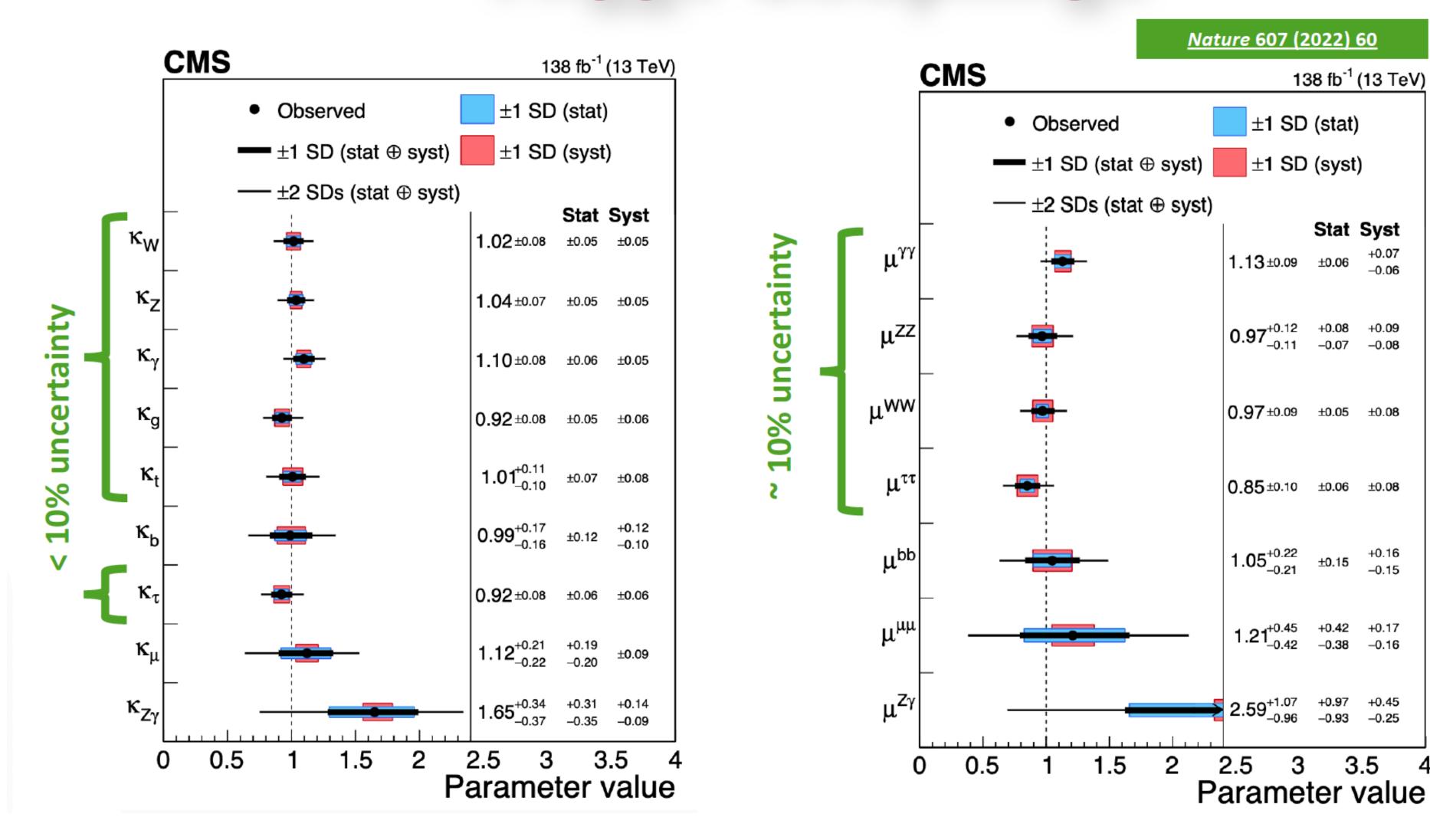




 $\mu = 1.01 \pm 0.08(\text{stat.}) \pm 0.04(\text{exp.}) \pm 0.05(\text{th.})$ 



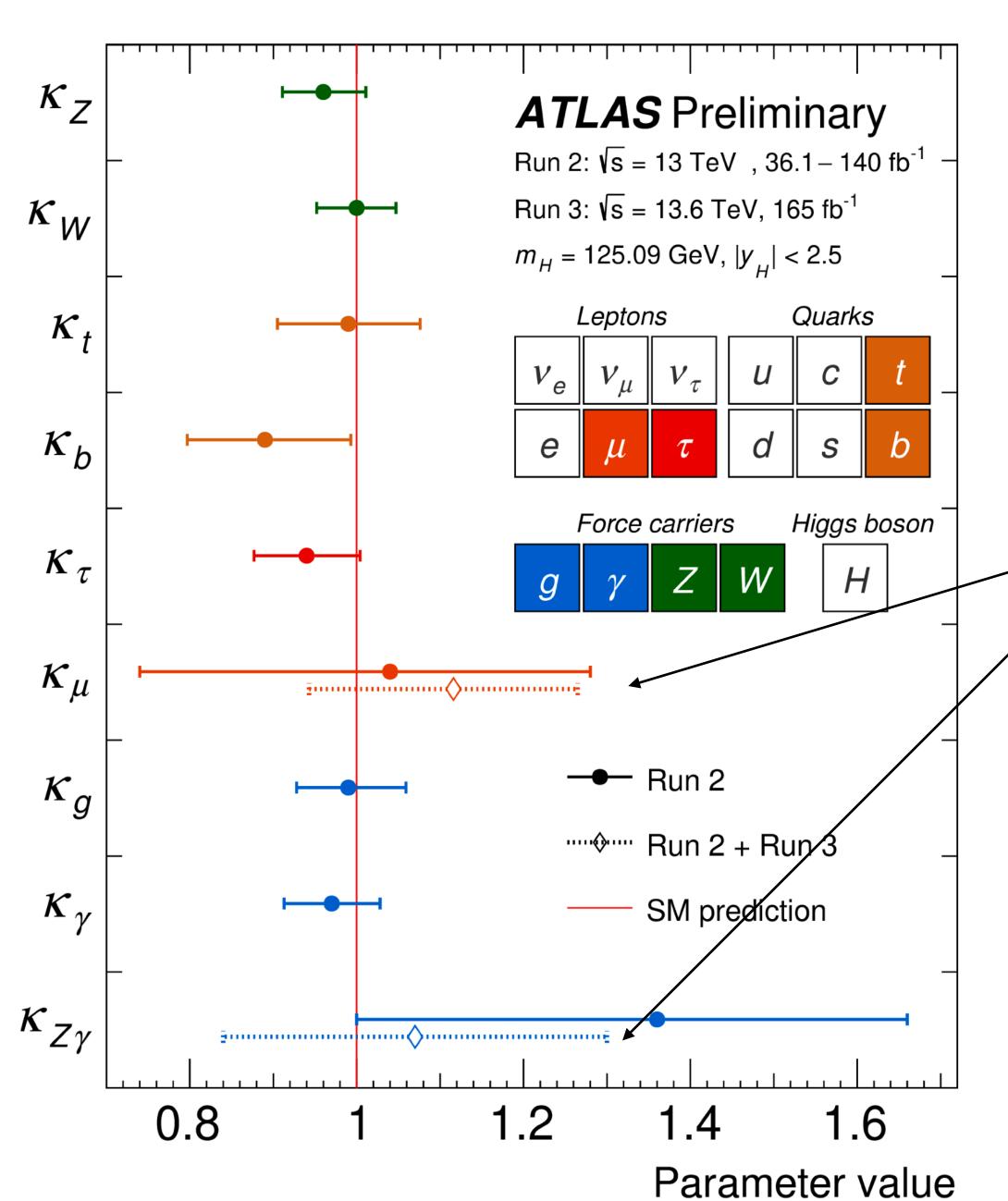
## Higgs Couplings



Run 2 Great precision on Higgs Couplings 30 times discovery statistics

Couplings to gauge bosons  $H \to \gamma \gamma$ ,  $H \to ZZ$ ,  $H \to WW < 10\%$  and 3rd generation fermions  $H \to \tau \tau$ , tt precision 10-12%;  $H \to bb$  observed ( 7 sigma)

# Higgs Couplings

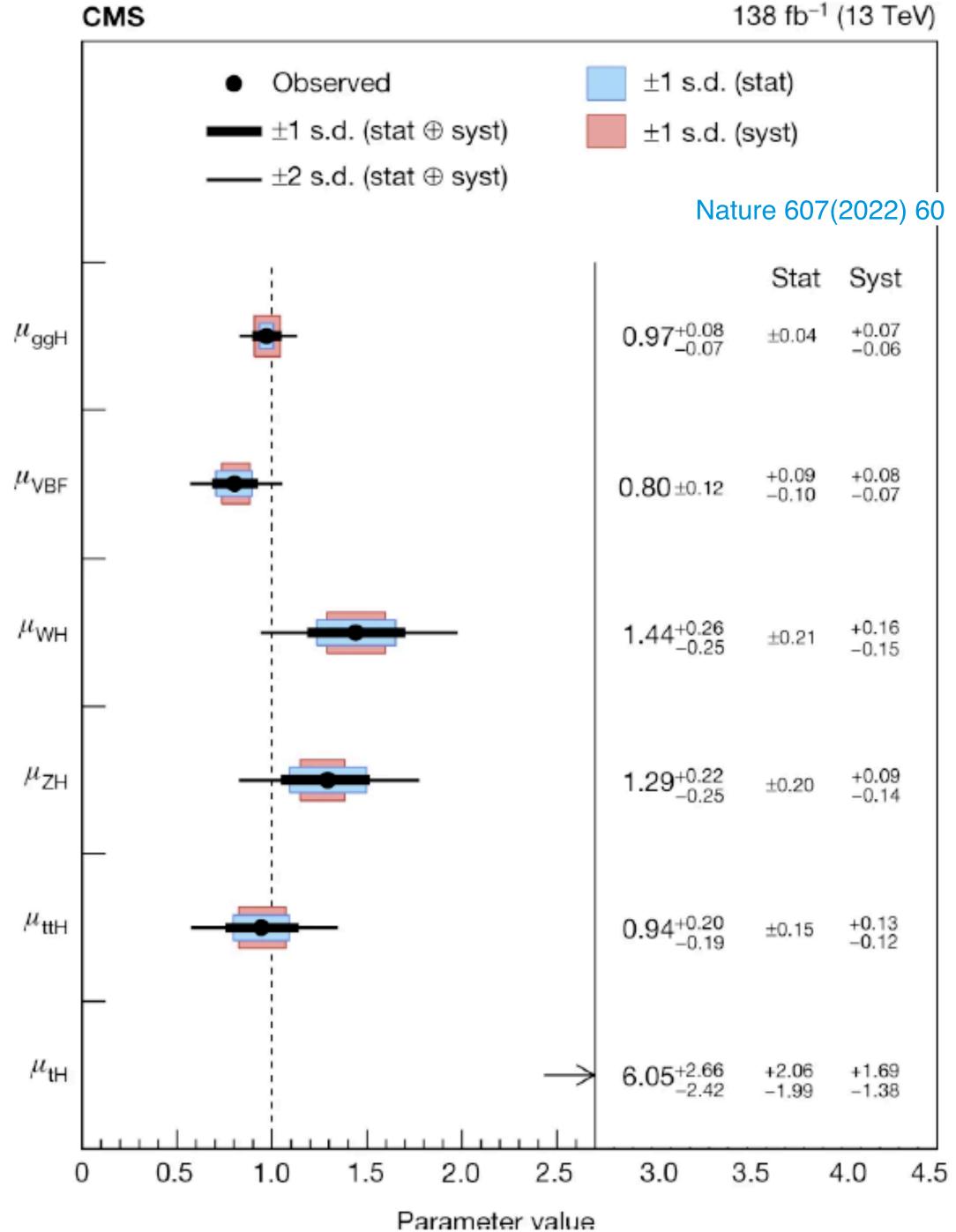


First run 3 measurements entering the coupling plots!
30% / 38% improvement in κμ / κΖγ

#### Higgs Production

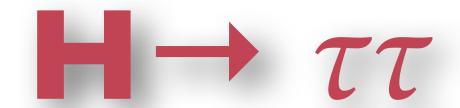
All the main production modes have been observed with a 5  $\sigma$  significance

- Main production modes (ggH, VBF) probed at ~
   10% level
- ttH ~ 21% level, VH~18%



# We are continuously significantly improving wrt to projections. Few very recent examples!

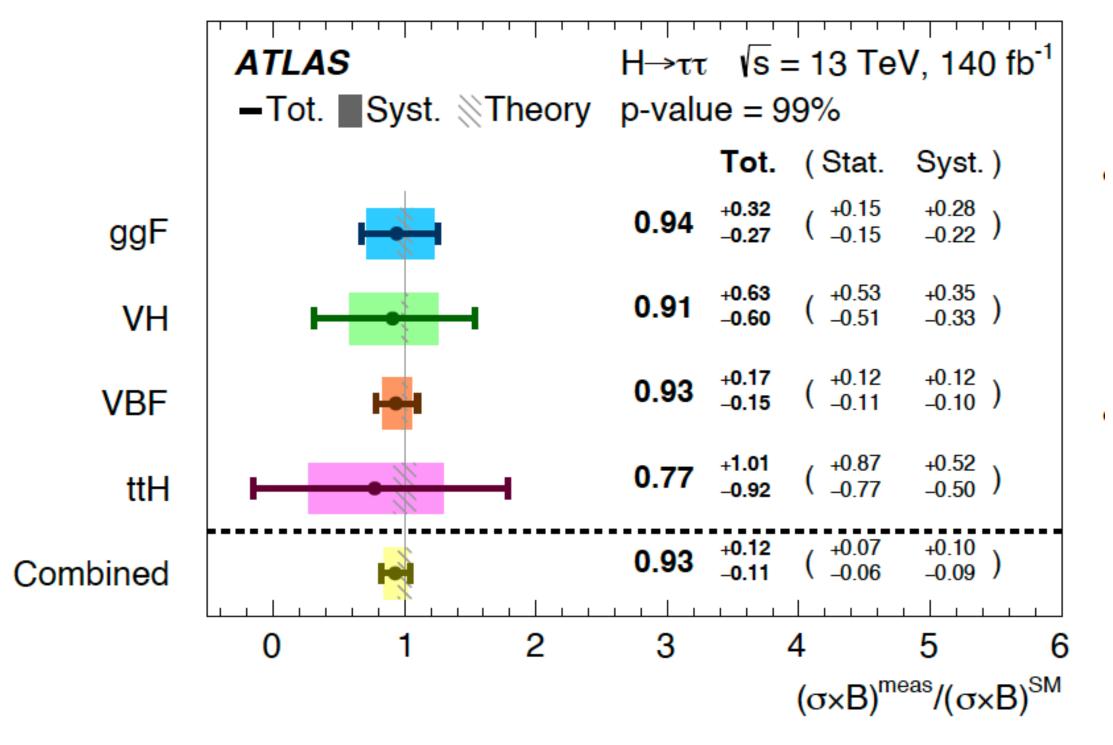
- $\bullet$  H $\rightarrow$   $\tau\tau$
- VH, H→cc and H→bb combination
- Couplings to 2nd generation
- rare processes
- Run 3 measurements



#### arXiv 2407.16320

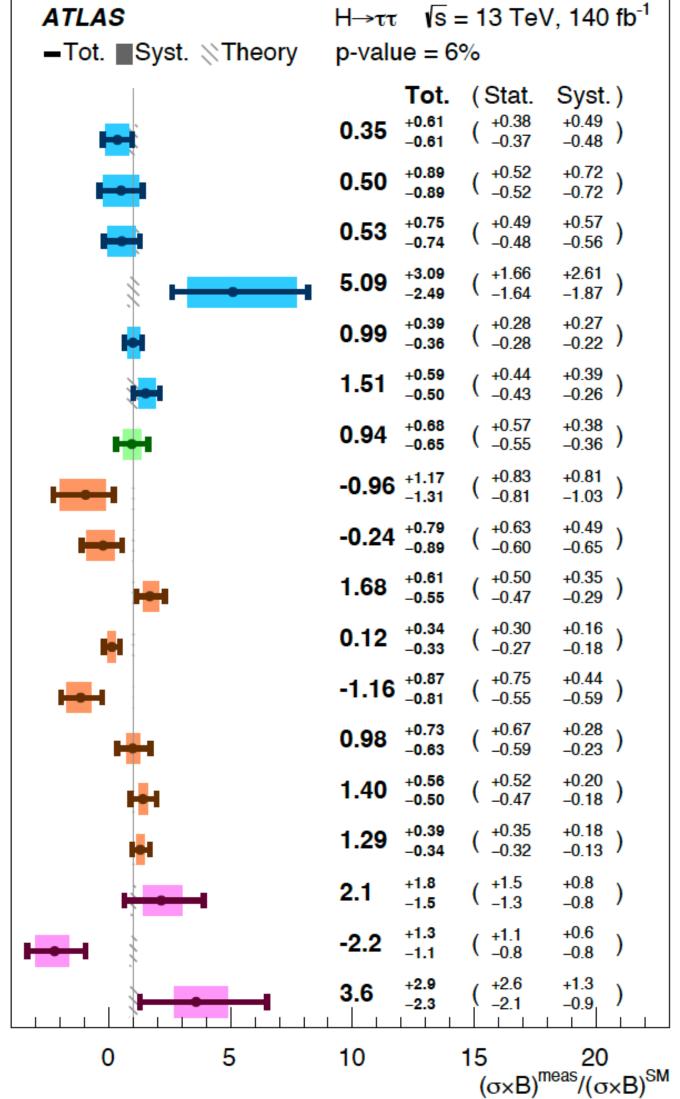
# Re-analysis of run 2 data $H \rightarrow \tau \tau$ has the strongest BR to leptons: Neural Network improves pTH resolution by 50%

#### all production modes



gg→H, 1-jet, 120 ≤ p<sub>T</sub> < 200 GeV gg→H, ≥ 1-jet,  $60 \le p_{\tau}^{H} < 120 \text{ GeV}$ gg→H, ≥ 2-jet, m<sub>ii</sub> < 350, 120 ≤ p<sub>T</sub><sup>H</sup> < 200 GeV gg→H, ≥ 2-jet, m<sub>ii</sub> ≥ 350 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV gg→H, 200 ≤ p<sub>T</sub> < 300 GeV gg→H, p<sub>T</sub> ≥ 300 GeV qq'→Hqq', ≥ 2**-**jet, 60 ≤ m<sub>ii</sub> < 120 GeV qq'→Hqq', ≥ 2-jet, 350 ≤ m<sub>ii</sub> < 700 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV qq'→Hqq', ≥ 2-jet, 700 ≤ m<sub>ii</sub> < 1000 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV qq'→Hqq', ≥ 2-jet, 1000 ≤ m<sub>ii</sub> < 1500 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV qq'→Hqq', ≥ 2-jet, m<sub>ii</sub> ≥ 1500 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV qq'→Hqq', ≥ 2-jet, 350 ≤ m<sub>||</sub> < 700 GeV, p<sub>T</sub> ≥ 200 GeV qq' $\rightarrow$ Hqq',  $\geq$  2-jet, 700  $\leq$  m<sub>ii</sub> < 1000 GeV, p<sub> $\tau$ </sub><sup>H</sup>  $\geq$  200 GeV qq'→Hqq', ≥ 2-jet, 1000 ≤ m<sub>ii</sub> < 1500 GeV, p<sub>T</sub><sup>H</sup> ≥ 200 GeV qq'→Hqq', ≥ 2-jet, m<sub>ii</sub> ≥ 1500 GeV, p<sub>T</sub><sup>H</sup> ≥ 200 GeV ttH,  $p_{\tau}^{H}$  < 200 GeV  $ttH, 200 \le p_{_{T}}^{H} < 300 \text{ GeV}$ 

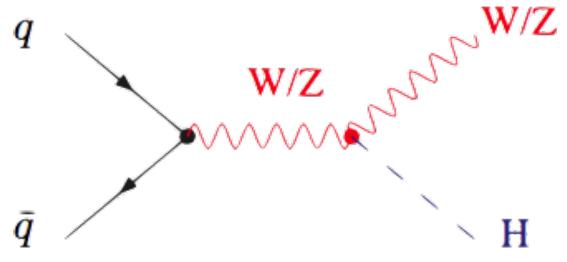
ttH,  $p_{\tau}^{H} \ge 300 \text{ GeV}$ 



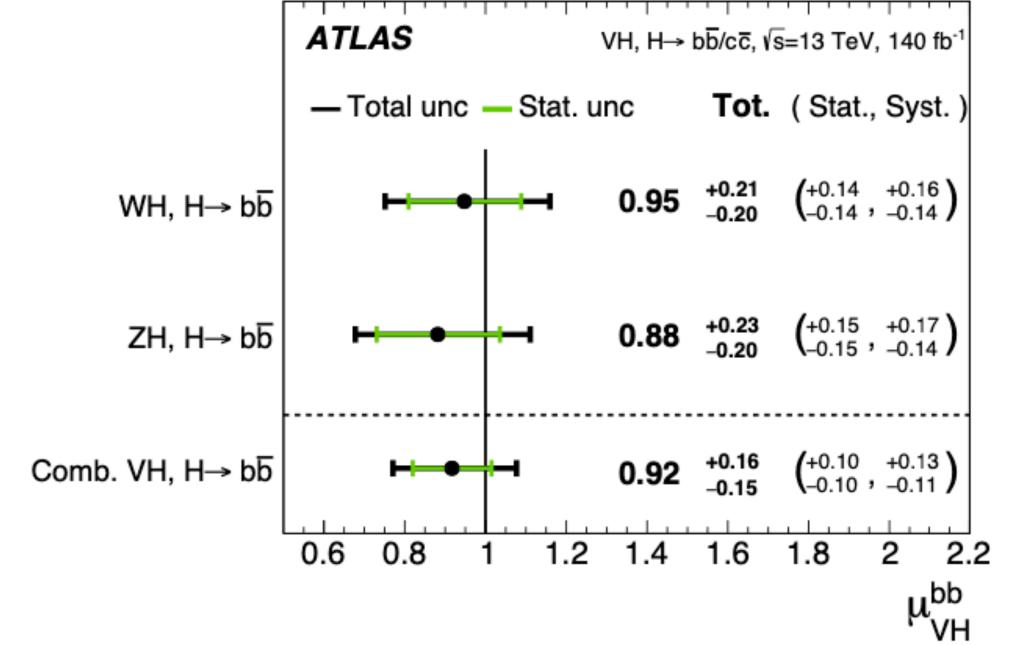
60

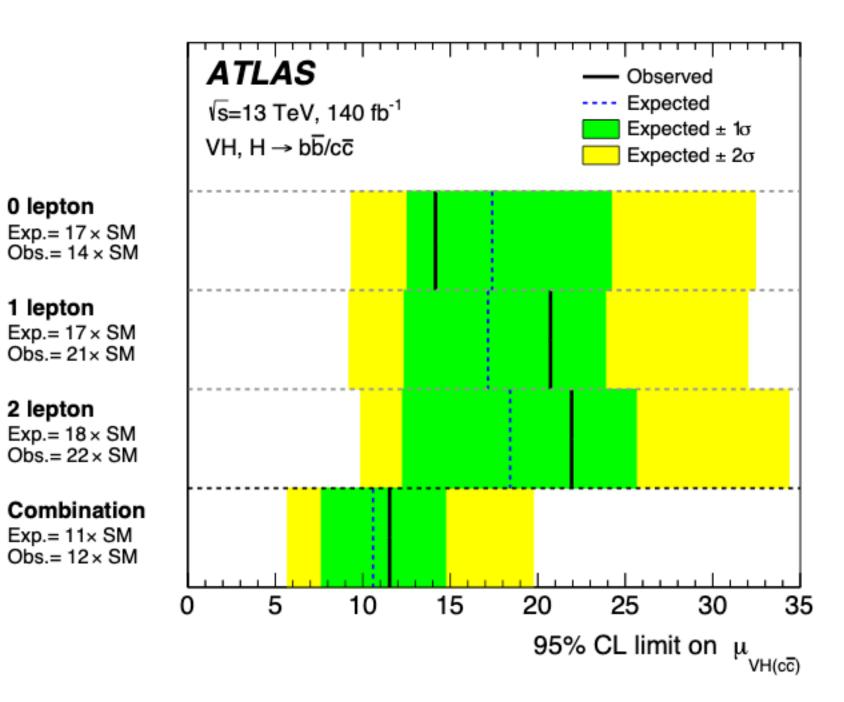
Most precise VBF measurement VBF at lowest pT First VBF measurements for pTH>200 GeV (30-50% accuracy)

#### VH, H→cc and H→bb combination

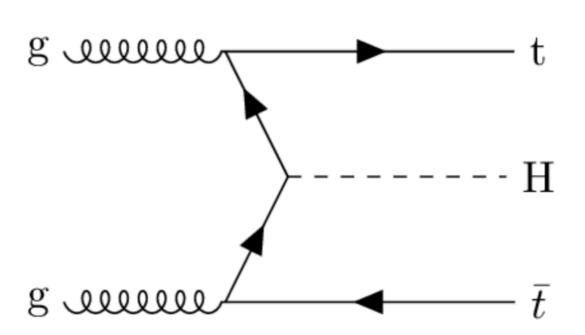


Run2 re-analysis: better reconstruction, better flavor tagging with b and c identification, better calibration, machine learning





H->cc limit on signal strength improved by a factor of 3 wrt first full Run 2 result !!



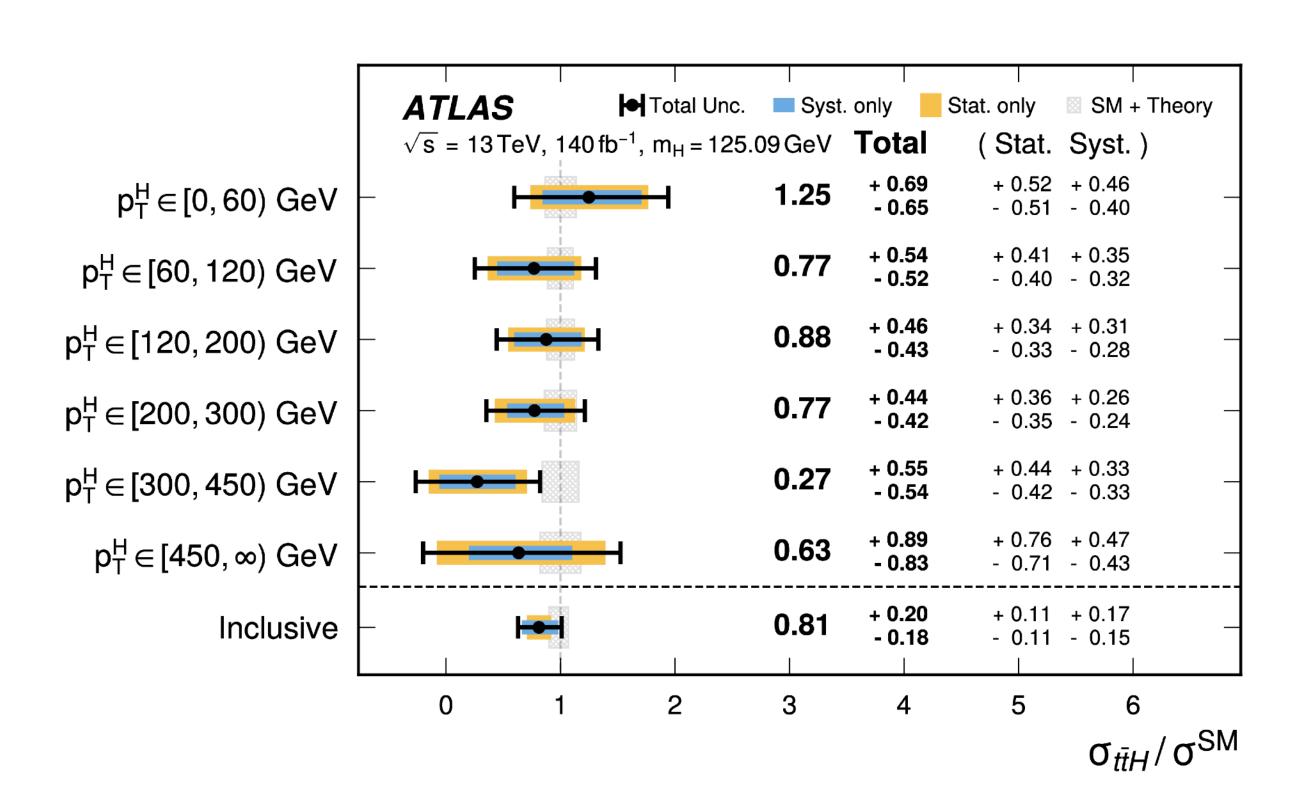


arXiv:2407.10904

4.6 sigma observation in this channel alone.

expected sensitivity improved from 2,7 to 5.4 sigma with same data

Various improvements NN-based classifier using transformers with attention mechanism, improvements in MC modeling



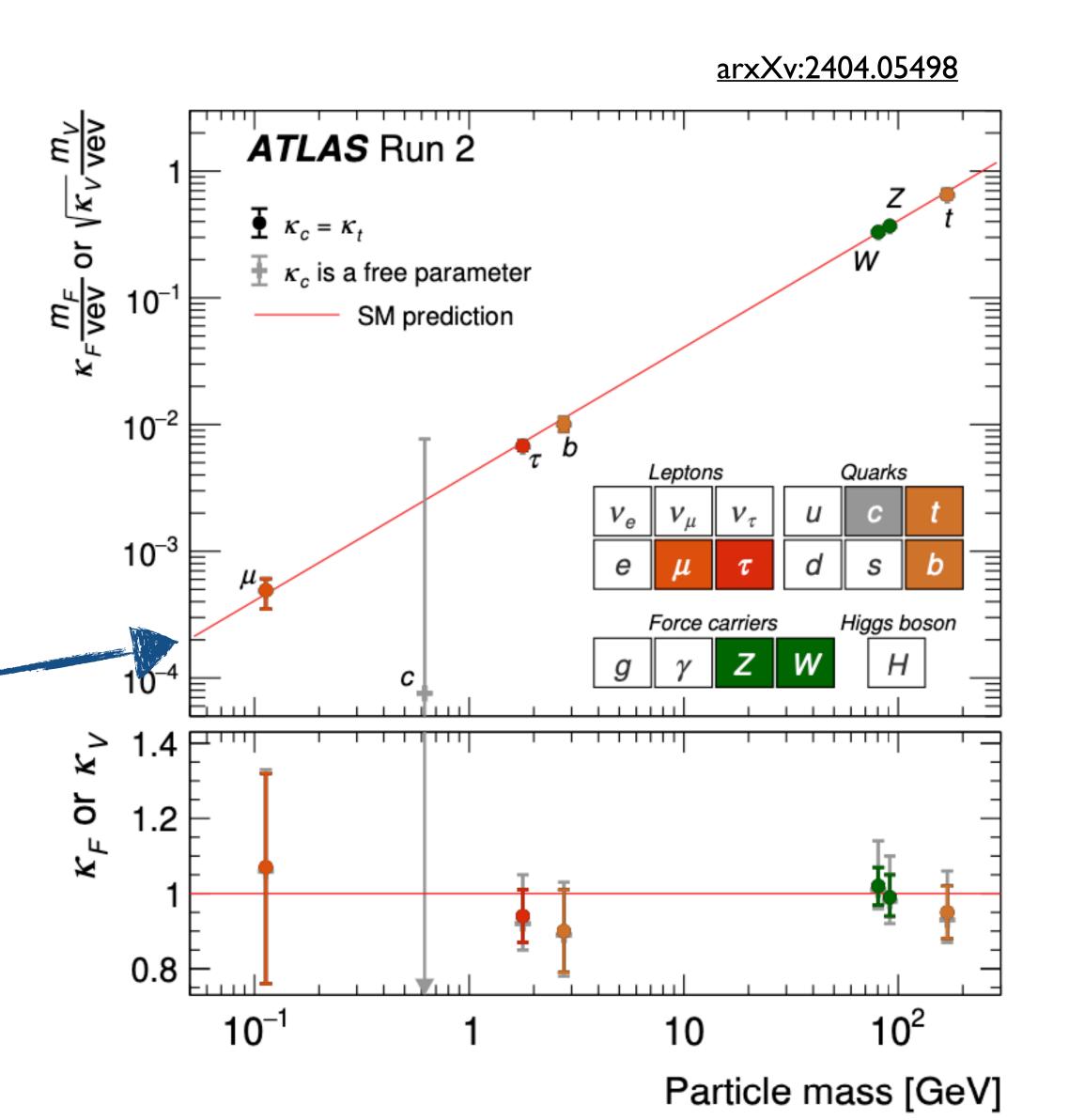
## Higgs couples to mass, 2<sup>nd</sup> generation?

Evidence that the Higgs couples to mass

Third generation and gauge boson couplings follow the SM predictions

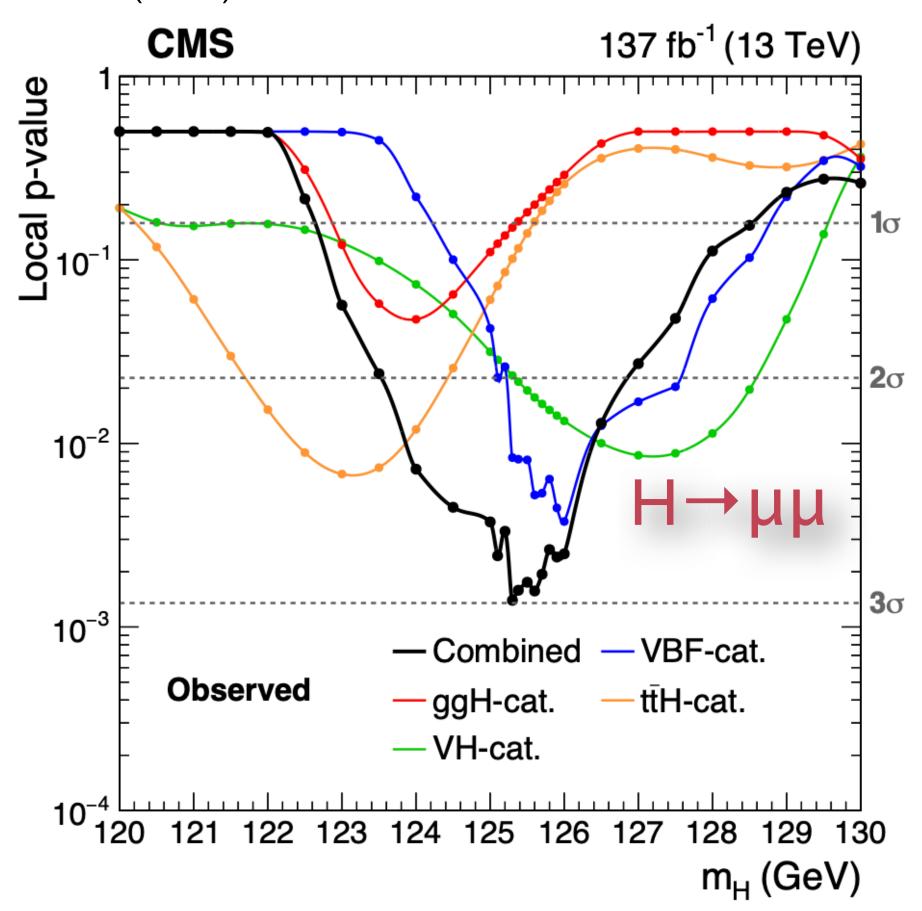
Next goal is to observe couplings to second generation (muons, c-quarks)

muons in reach soon (both experiments have already 3 sigma evidence)

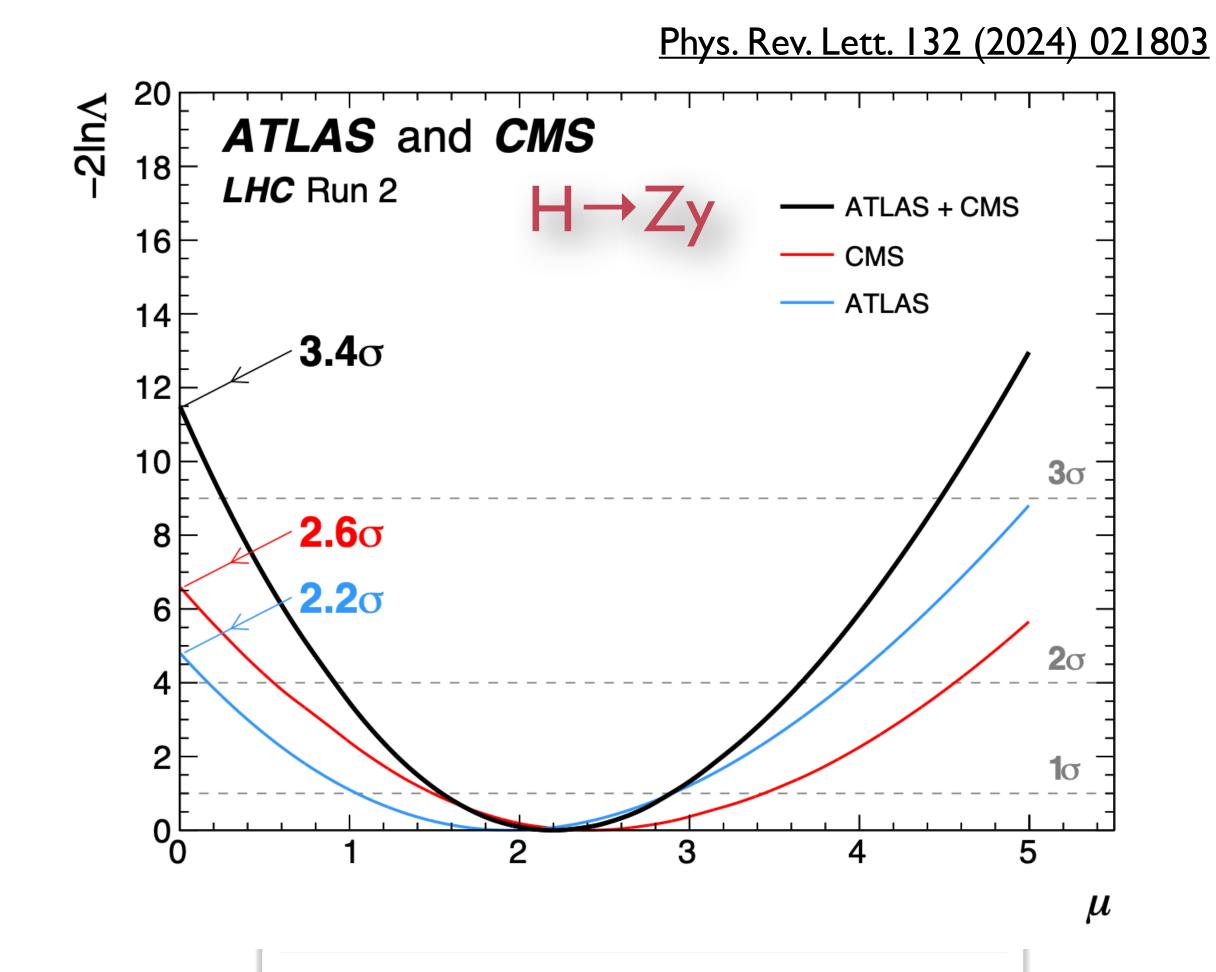


## Rare decays

JHEP 01 (2021) 148



First evidence for the SM Higgs decay to fermions of the 2nd generation ggH, VBF, VH and TT

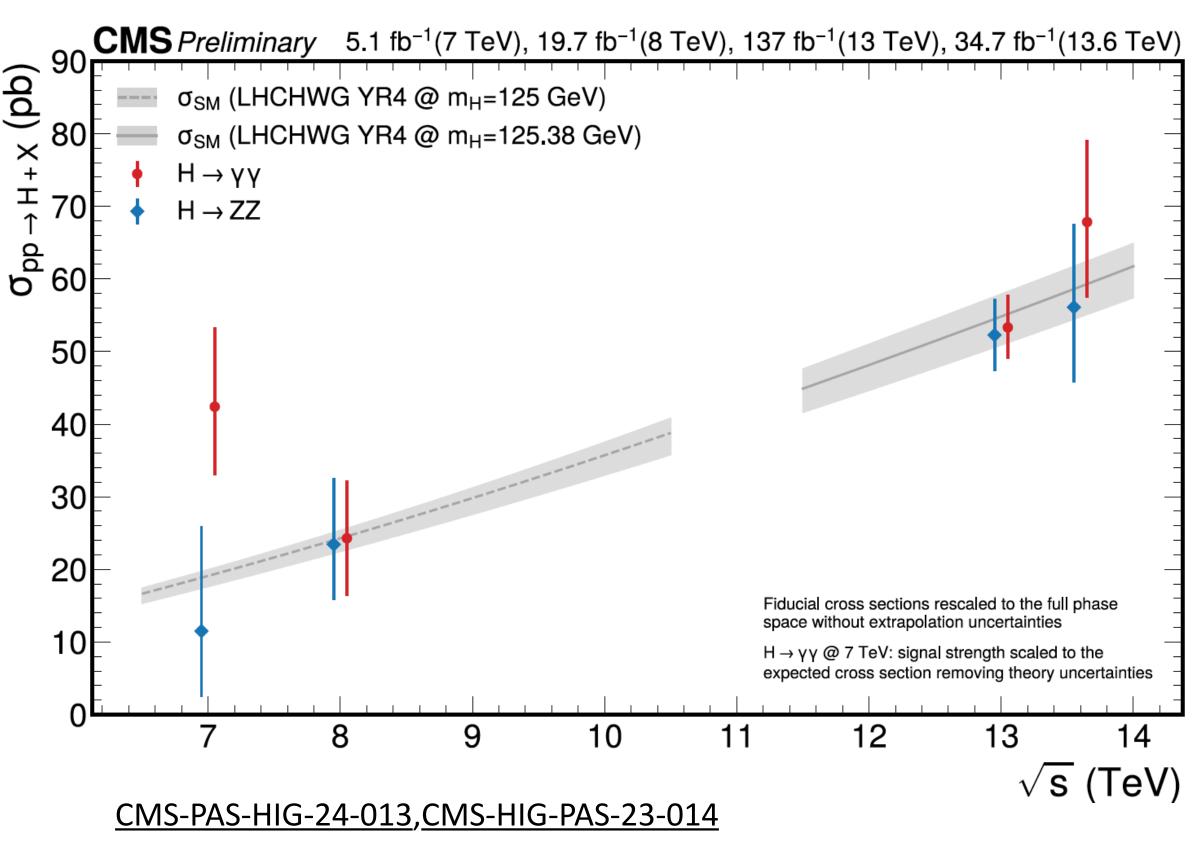


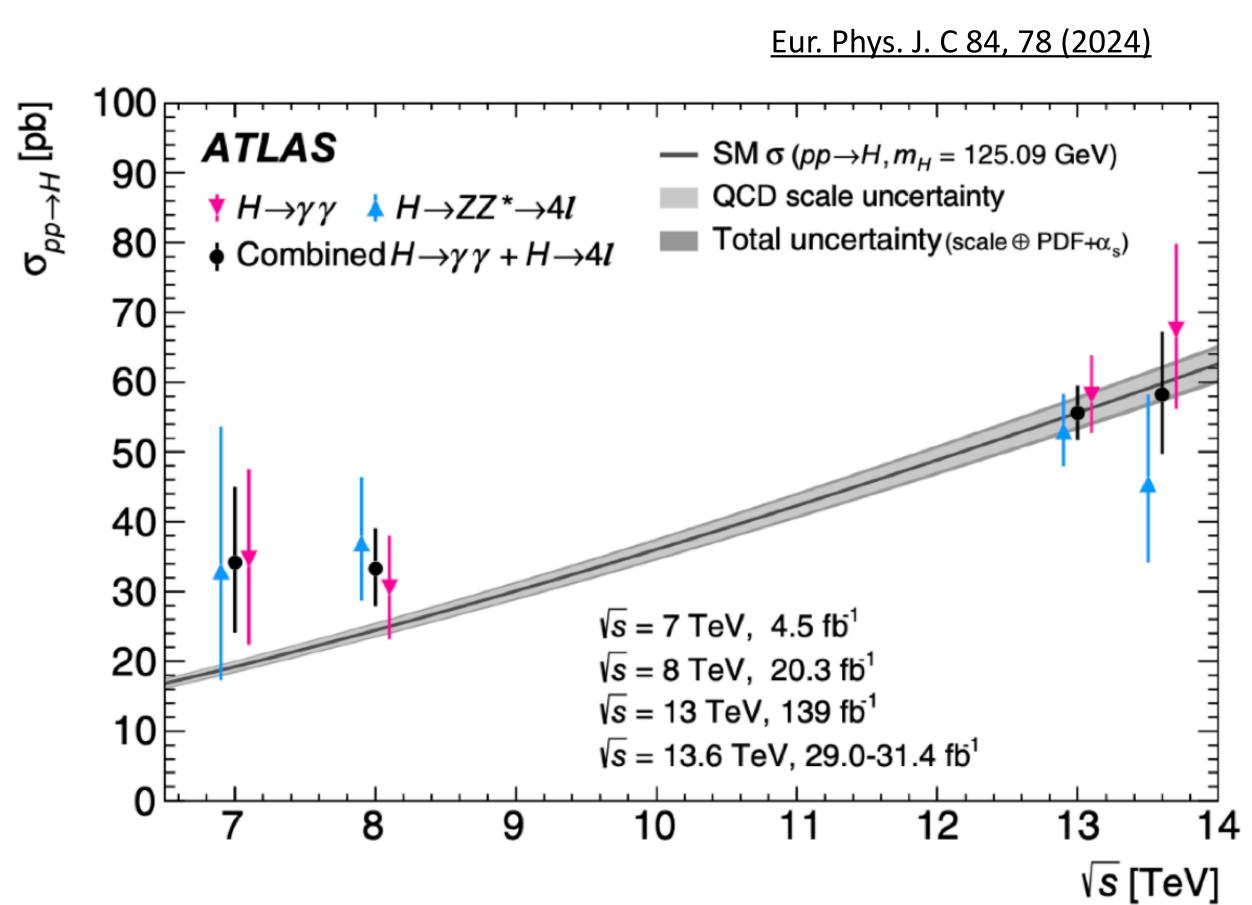
combination ATLAS+CMS

evidence at  $3.4 \sigma$ 

#### Run 3 measurements!

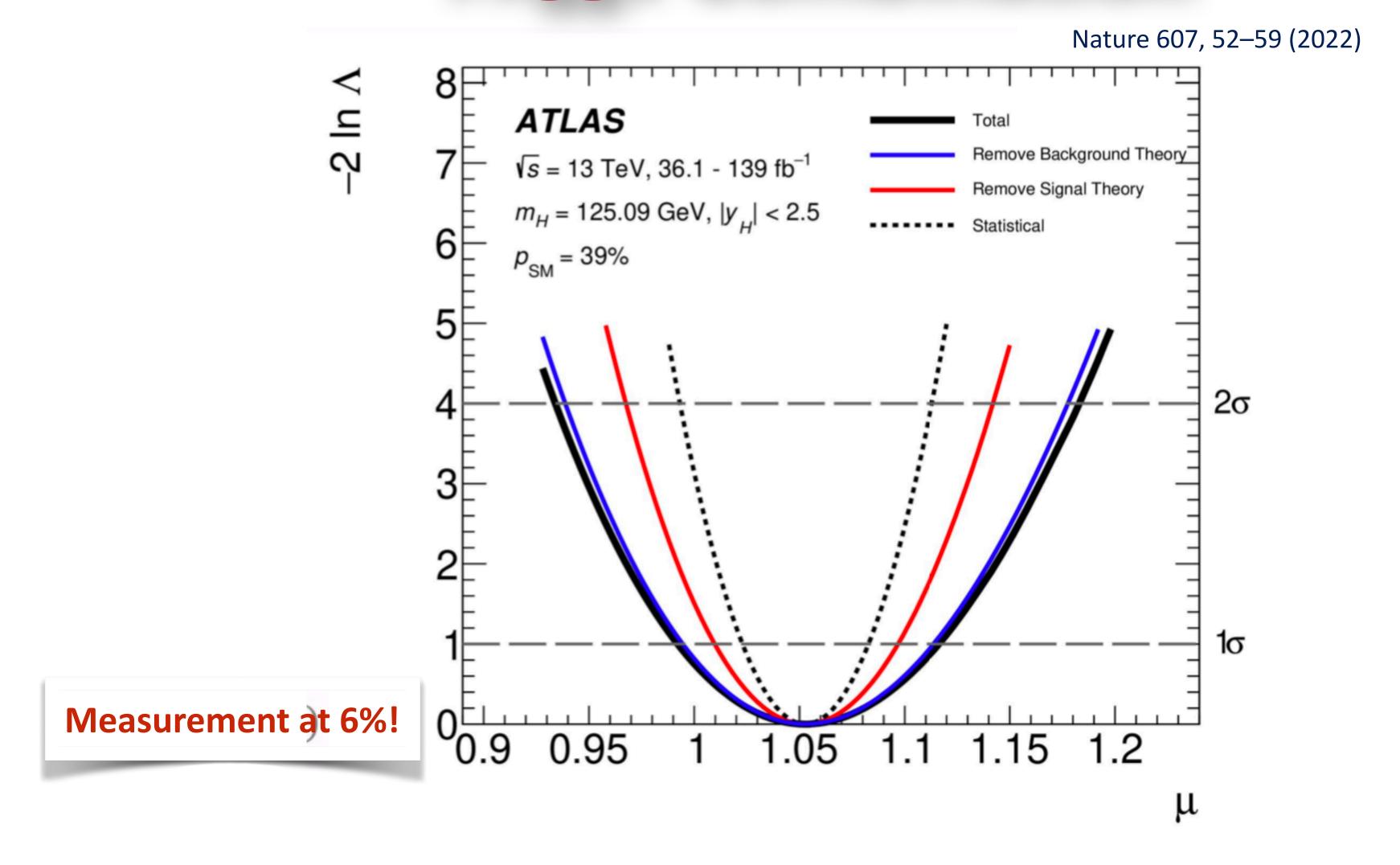
#### Fiducial cross-sections (rescaled to inclusive)





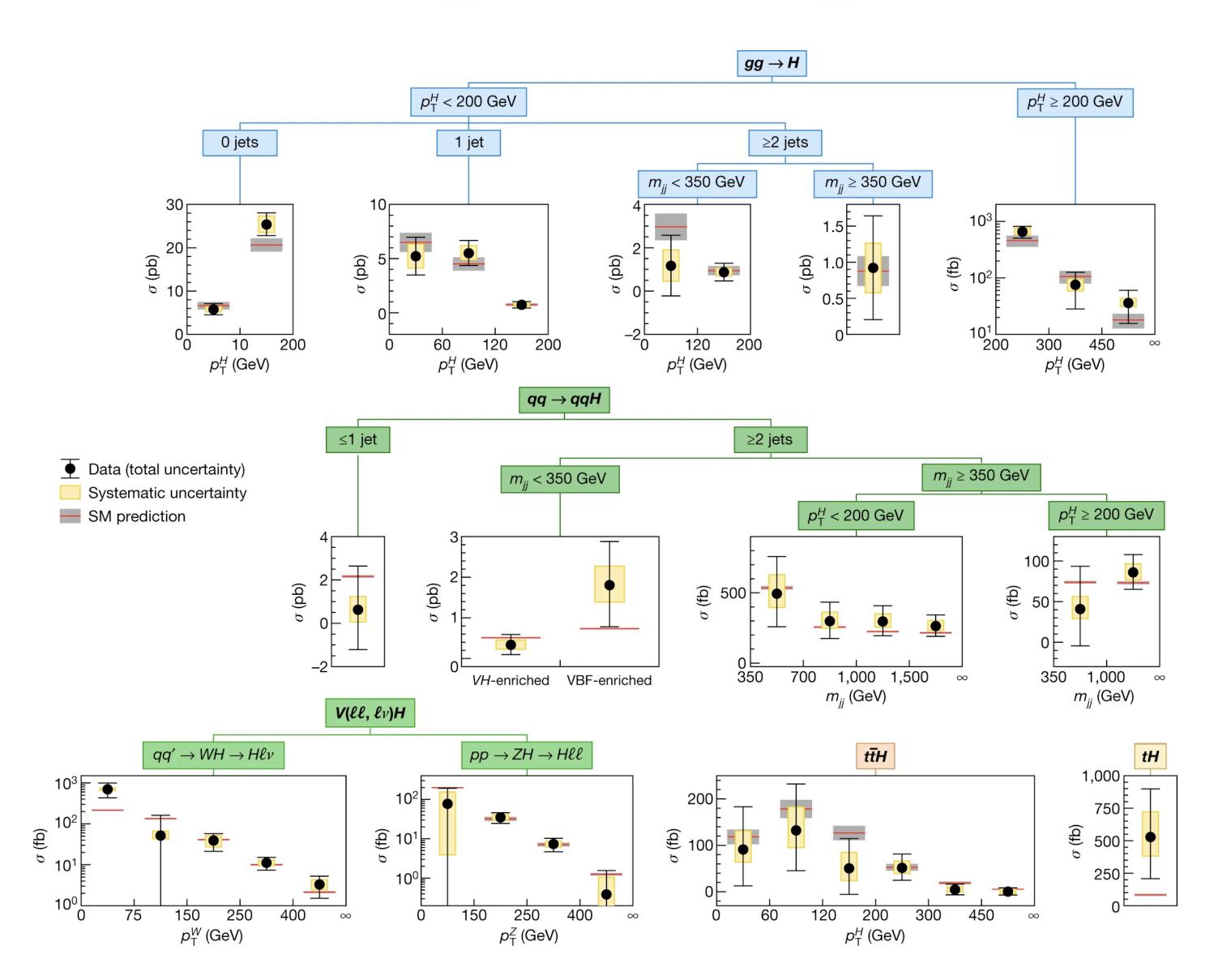
# Putting it all together

#### Higgs Combination



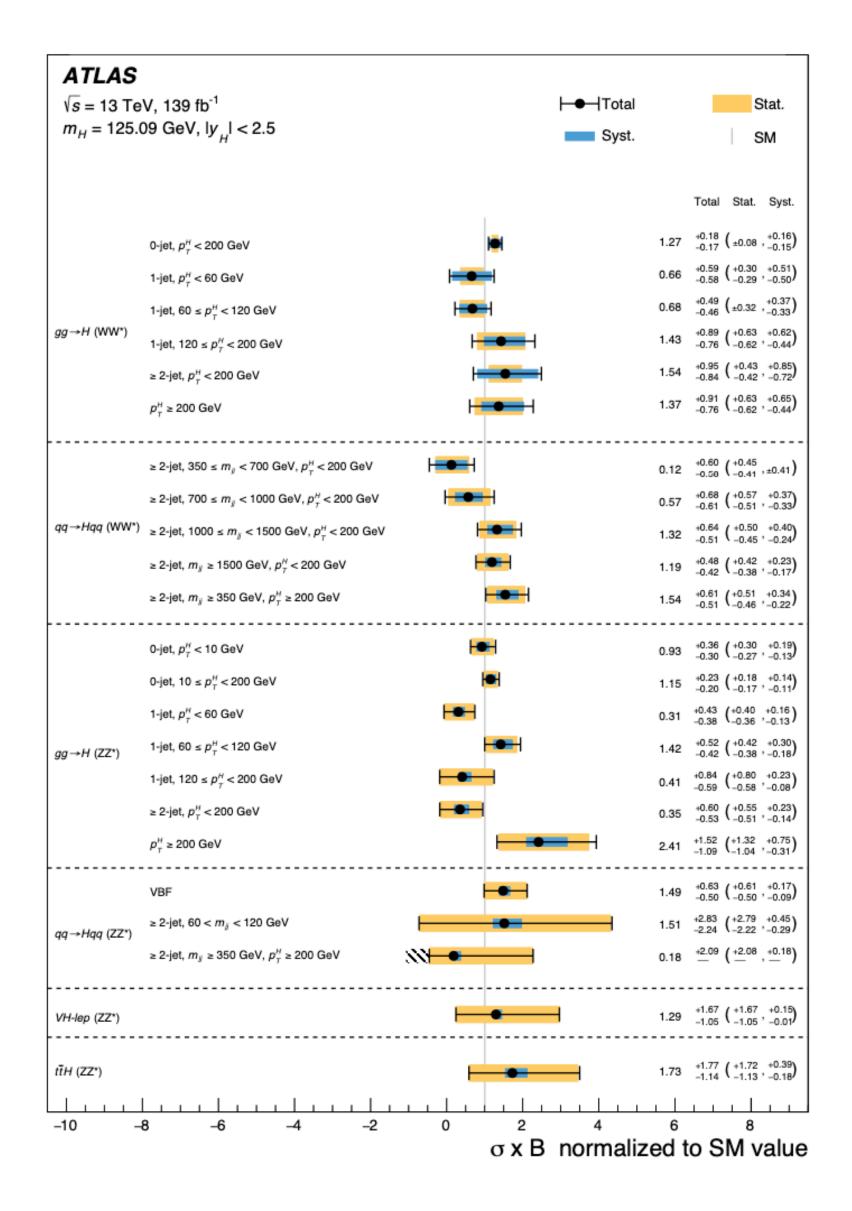
$$\mu = 1.05 \pm 0.06 = 1.05 \pm 0.03$$
 (stat.)  $\pm 0.03$  (exp.)  $\pm 0.04$  (sig. th.)  $\pm 0.02$  (bkg. th.).

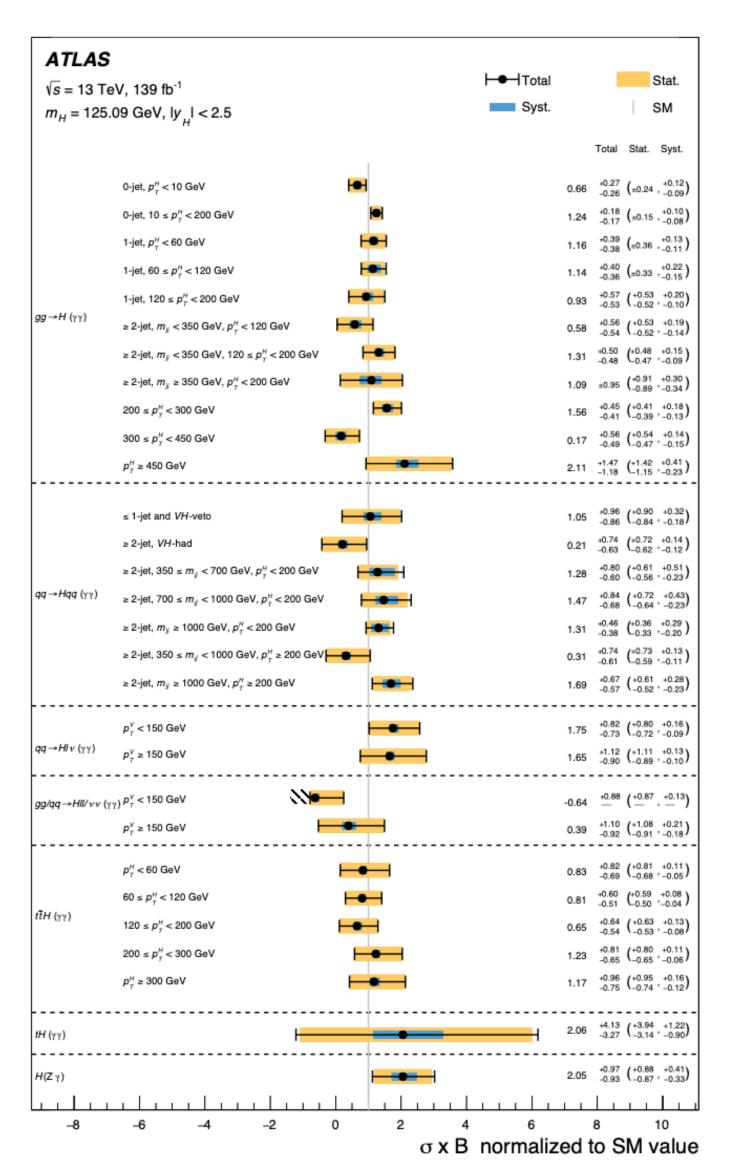
#### Simplified Template X-Sections full Run 2

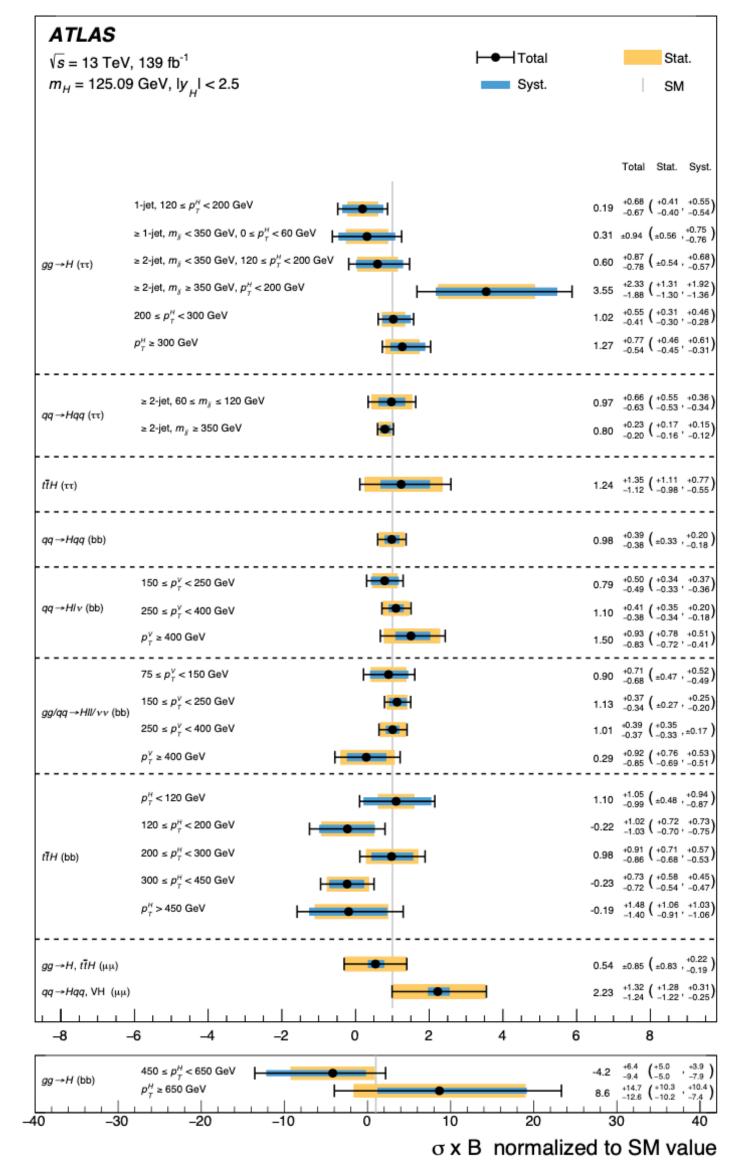


- Measurements of cross sections σ in mutually exclusive regions split based on Higgs kinematics (+ W or Z bosons and associated jets)
- Reduce large theory uncertainties and minimize model-dependence when extrapolating to signal regions
- Maximize sensitivity to possible BSM

#### STXS







#### Effective Field Theories

How to be sensitive to New Physics beyond the Standard Model in a Model Independent way?

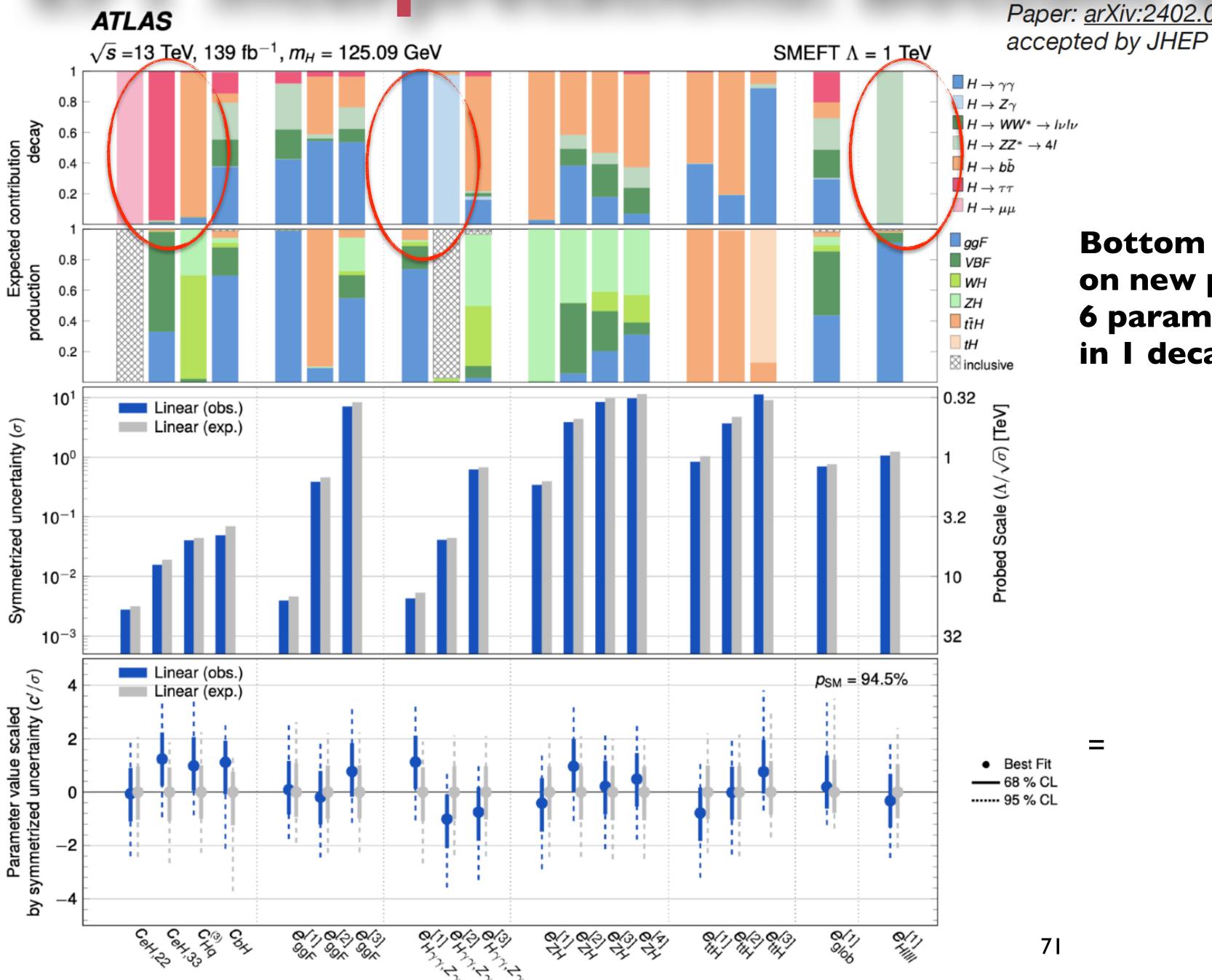
Effective Field Theory parametrizes Beyond Standard Model (BSM) effects at high energies ( $\Lambda \gg v$ , above electroweak scale) at low energies,  $E \ll \Lambda$ , in terms of higher-dimensional operators in an effective Lagrangian:

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{i}^{N_{d=6}} \frac{c_i}{\Lambda^2} O_i^{(6)} + \sum_{j}^{N_{d=8}} \frac{b_j}{\Lambda^4} O_j^{(8)} + \dots,$$



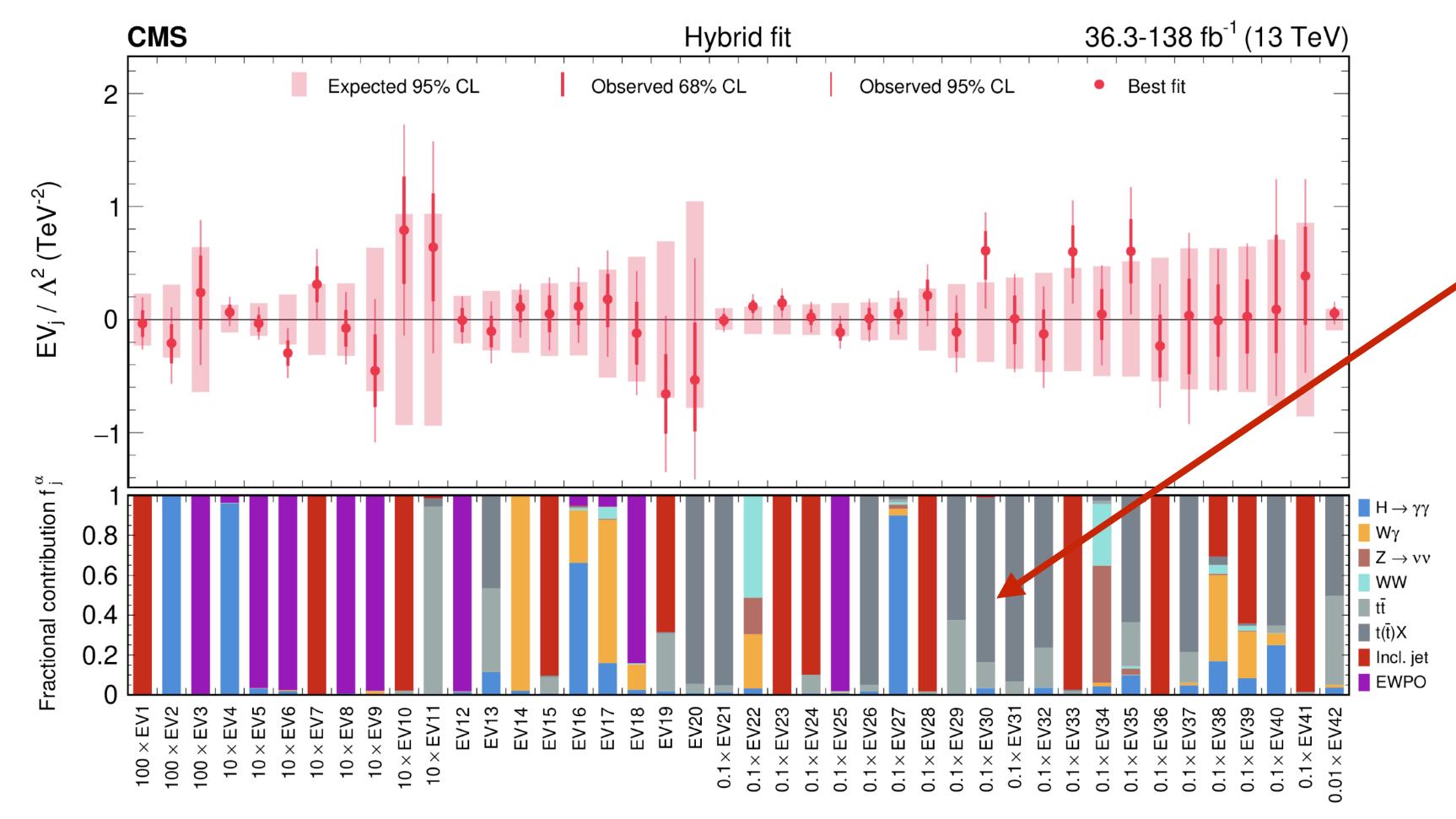
dimensionless Wilson coefficients of higher dimension operators

# EFT interpretations: Bottom up approach



Bottom up approach means that we are on new physics model.
6 parameters nearly exclusively measured in I decay mode

## Higgs+EW+top+EWPO EFTs



This is just NOT just Higgs!
We use EW+top+Higgs +EWPrecision
Observables from LEP

We also adding in the game the Charge and Parity of the Higgs.

The Higgs sector could contain a partial violation of the Charge and Parity that could explain the matter antimatter asymmetry.

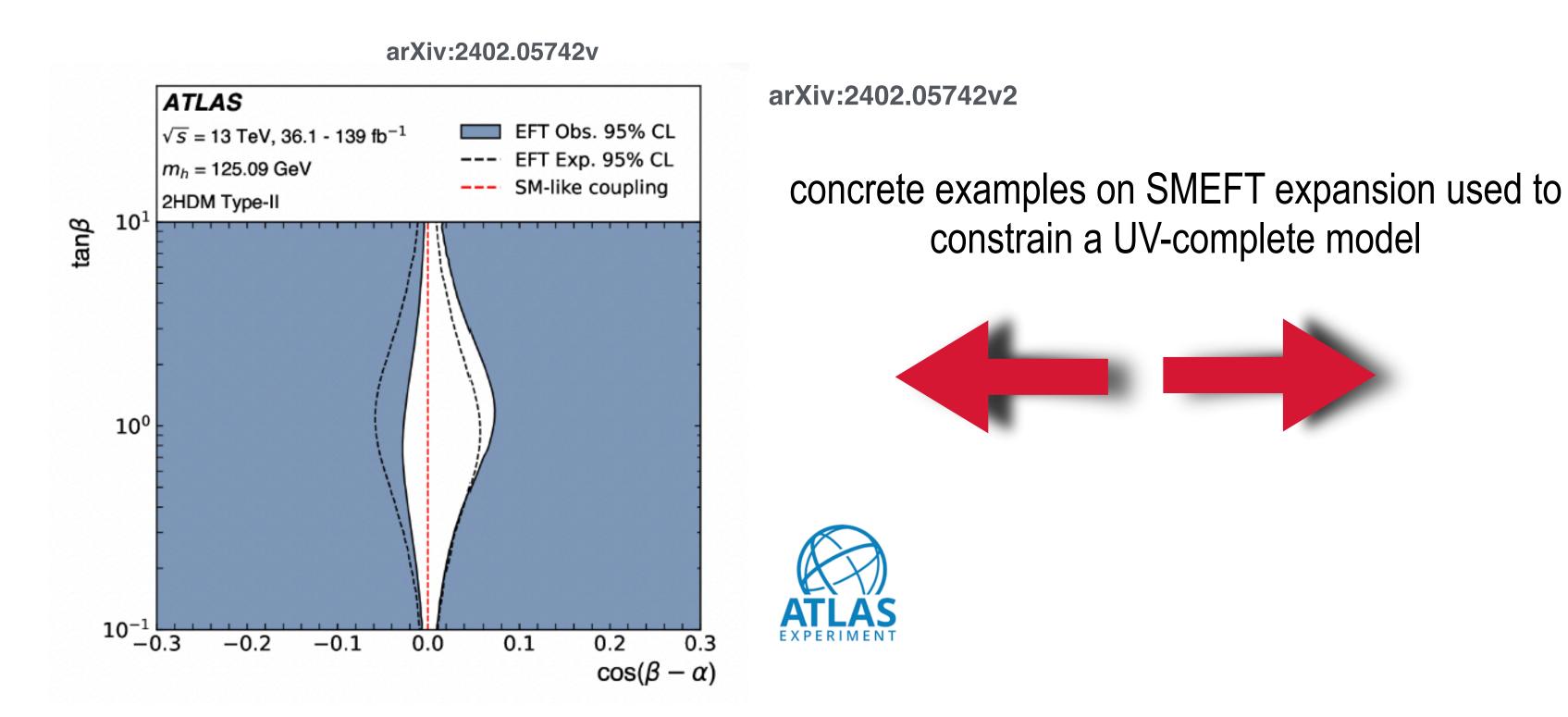
#### EFT: Top down approach

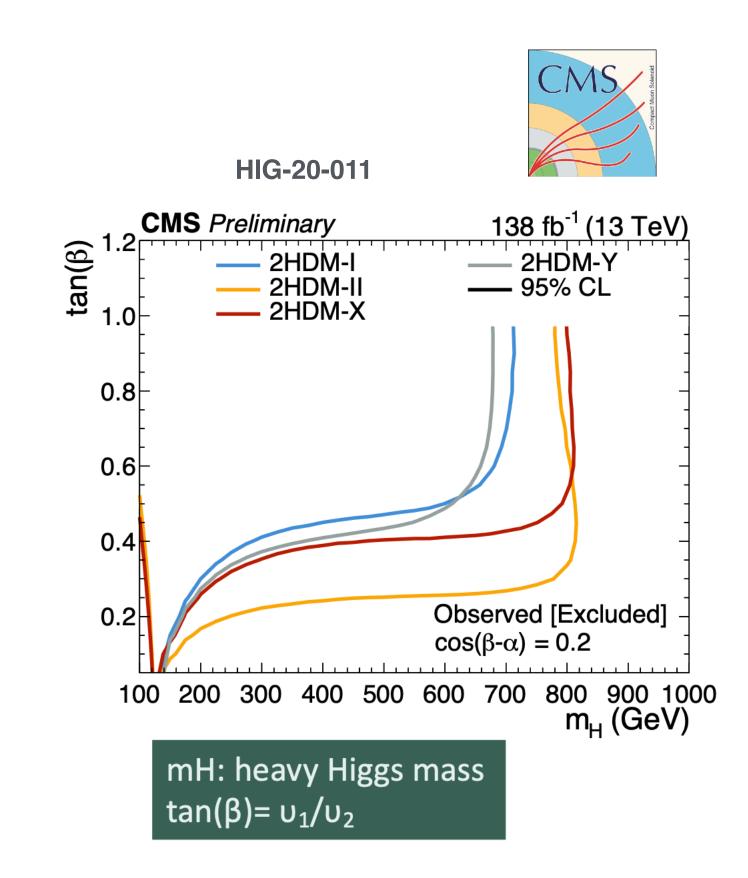
What about guiding the measurements such that they can be more sensitive to Wilson Coefficients that would present larger variations within specific models?

- Comes at the cost of losing generality, but can lead help us not to miss important features in the data.
- Requires matching models with EFTs.

The leading IR/UV dictionary (tree-level, dimension 6 SMEFT) is already there since few years ago.

#### Example: The 2HDM Lagrangian can be written as a SMEFT expansion!





arXiv:2112.10787v1

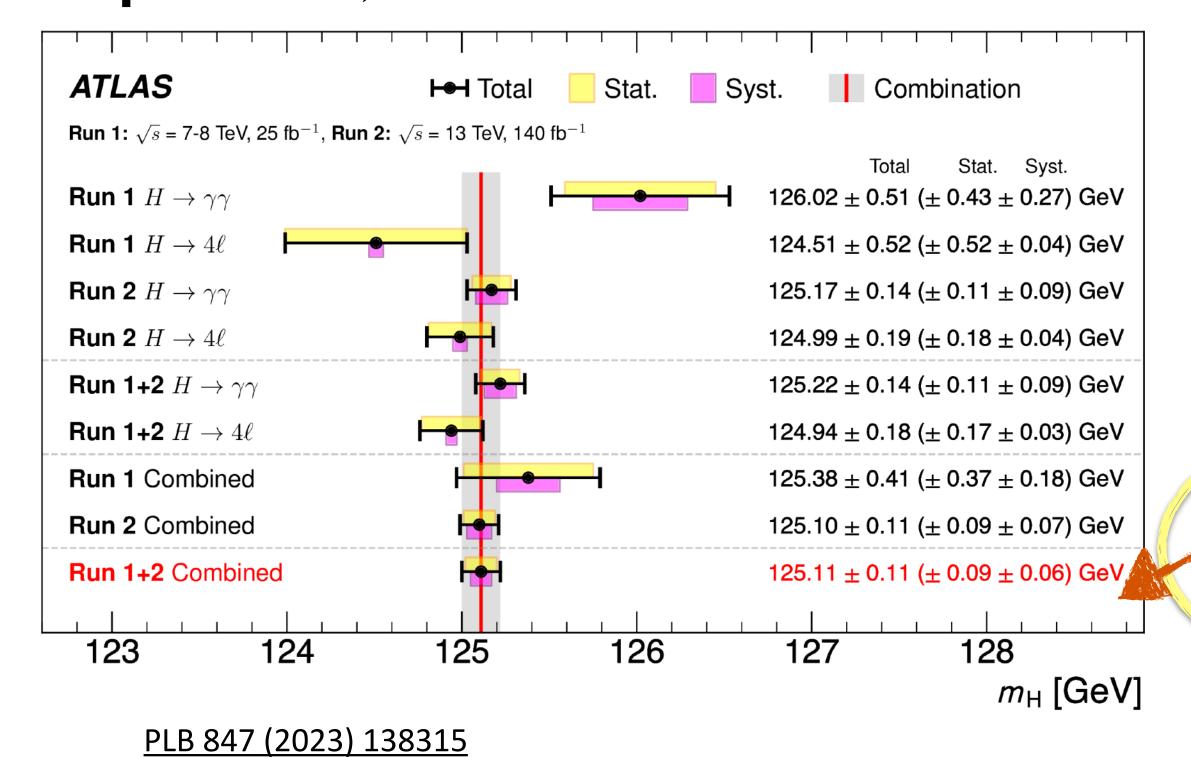
### Lecture 2

### Higgs mass

 $m_H = 125.04 \pm 0.12$ 



Higgs mass precision is now at the 0.1% level for **both experiments**, was 0.2% after Run 1.



PRL 131 (2023) 251802

arXiv:2409.13663 **CMS Run 2**: 138 fb<sup>-1</sup> (13 TeV) Stat. Only **Run 1**: 5.1 fb<sup>-1</sup> (7 TeV) + 19.7 fb<sup>-1</sup> (8 TeV) Total (Stat. Only) 124.90 <sup>+0.15</sup><sub>-0.15</sub> (<sup>+0.14</sup><sub>-0.14</sub>) GeV **4**μ 124.70 <sup>+0.53</sup><sub>-0.51</sub> (<sup>+0.49</sup><sub>-0.47</sub>) GeV 4e 125.50 <sup>+0.27</sup><sub>-0.26</sub> (<sup>+0.25</sup><sub>-0.24</sub>) GeV **2e2**μ 125.20 <sup>+0.29</sup><sub>-0.27</sub> (<sup>+0.27</sup><sub>-0.26</sub>) GeV **2**μ**2e** 125.04 +0.12 (+0.11) GeV Run 2 125.60 <sup>+0.46</sup><sub>-0.45</sub> (<sup>+0.43</sup><sub>-0.41</sub>) GeV Run 1 125.08 <sup>+0.12</sup><sub>-0.12</sub> (<sup>+0.10</sup><sub>-0.10</sub>) GeV Run 1 + Run 2 122 124 126 128 130

m<sub>H</sub> (GeV)

ATLAS most precise m<sub>H</sub> measurement at 0.09%

 $m_H = 125.11 \pm 0.11 \text{ GeV}$ 

#### Profits of various performance improvements:

•~4x improvements in photon energy calibration!

**→**Reduces  $H \rightarrow \gamma \gamma$  systematics by factor 3:

**320 MeV** → **80 MeV** 

#### Higgs width measurement H\* → ZZ → 4ℓ

Measurement of the Higgs boson width is indirect and as such it is 3 orders of magnitude smaller than the detector mass resolution! (LHC was not supposed to measure it!)

$$\sigma_{onshell}^{pp o H o ZZ} \sim rac{g_{Hgg}^2 g_{HZZ}^2}{m_H \Gamma_H}$$

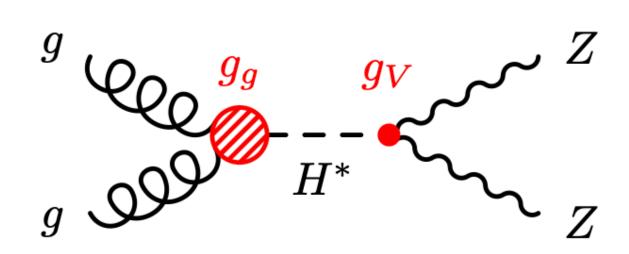
$$\sigma_{offshell}^{pp o H o ZZ} \sim rac{g_{Hgg}^2 g_{HZZ}^2}{M_{ZZ} - m_H}$$

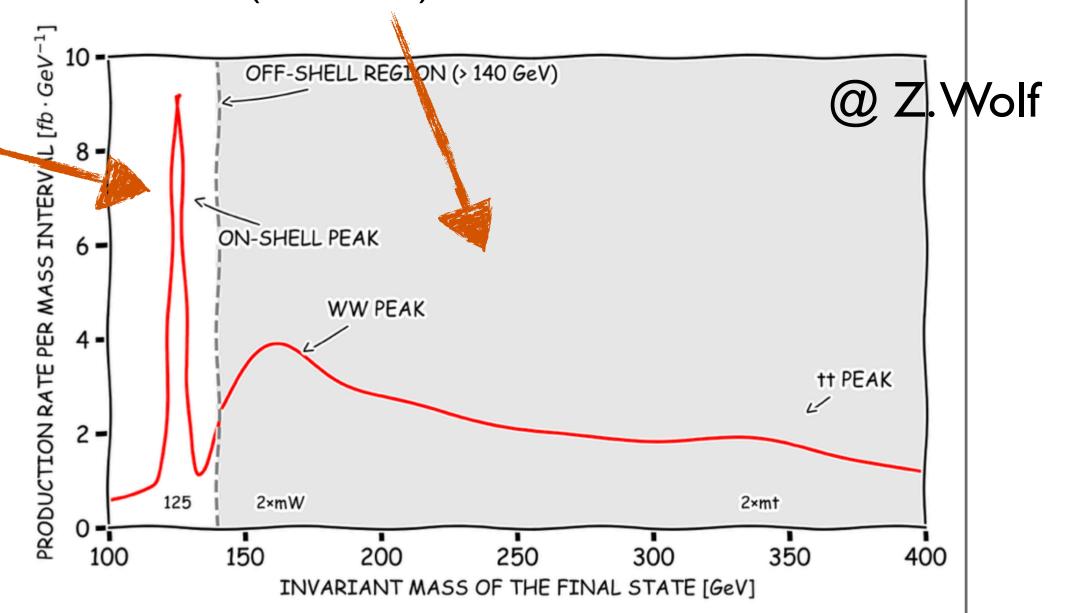


$$rac{\sigma_{offshell}}{\sigma_{onshell}} \propto \Gamma_{H}$$

It is measured from the region where mzz> mH (offhsell) vs

 $M_{ZZ}\sim M_H$  (onshell)

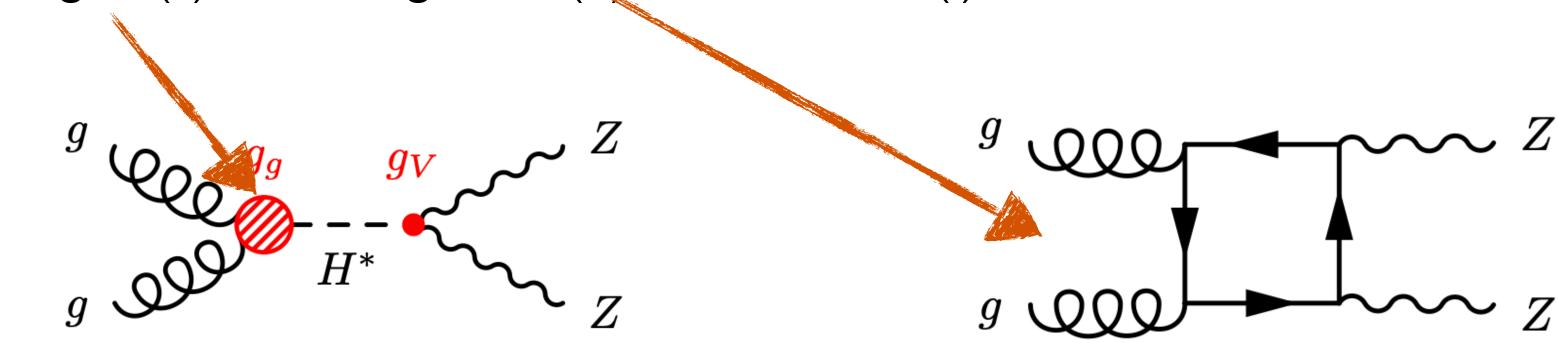




#### Higgs with measurement H\* → ZZ → 4ℓ

$$\frac{\sigma_{offshell}}{\sigma_{onshell}} \propto \Gamma_{H}$$

New Idea: Neural networks used before in off shell analysis, but not ideal due to Signal (S) and Background (B) interference (I).



Now using Neural simulation-based inference.

ATLAS Innovative method using neural simulation based-inference (NSBI) significance improved from 0.5 to 1.3 $\sigma$ .

NSBI handles high-dimensional parameter estimation without the need to bin data into low-dimensional summary histograms.

Statistical inference is the process of using data from a sample to make conclusions (inferences) about a larger population or about the underlying process that generated the data.

#### Higgs width

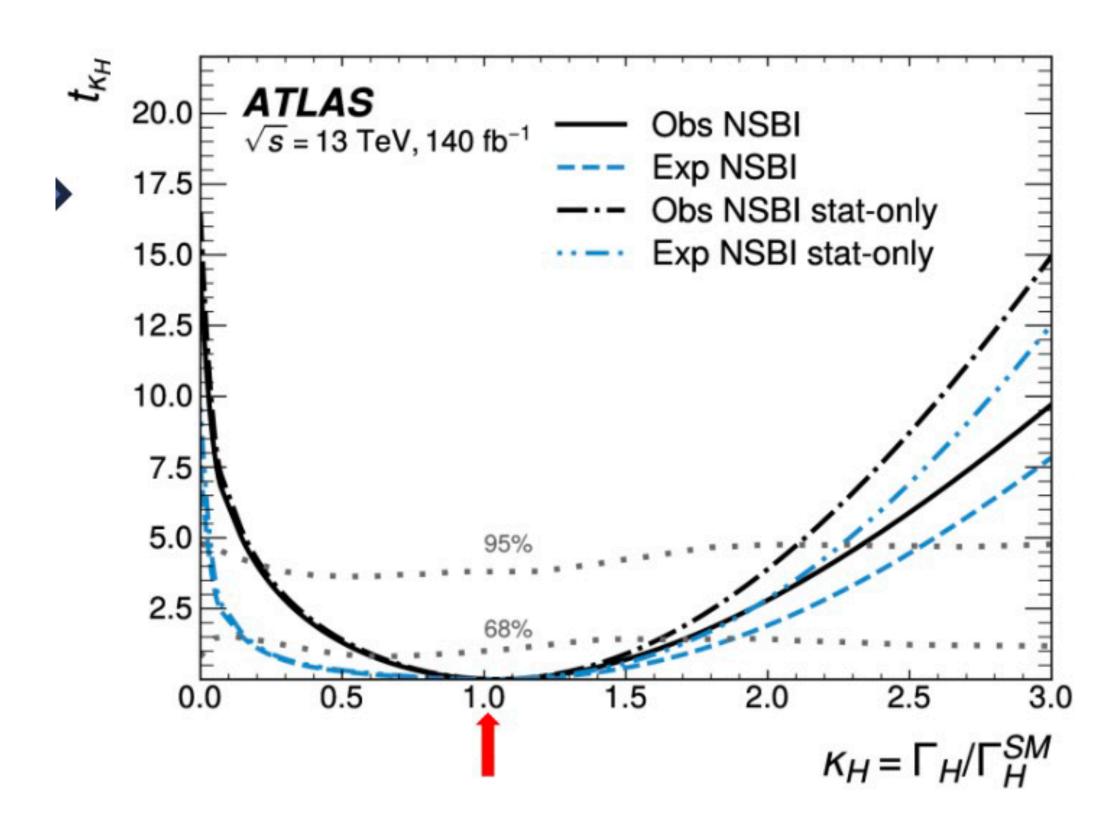
Log Likelihood tells you how compatible are the data with a certain hypothesis

In final state with Z-> 4I / 2I2v

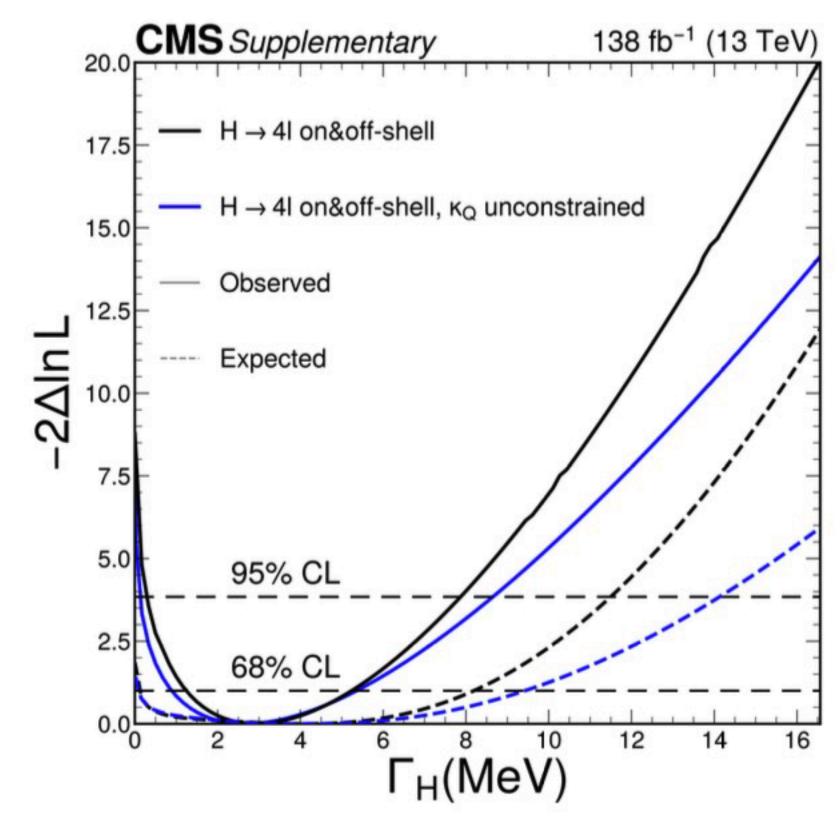
ATLAS 
$$+2.7$$
  $\Gamma_H = 4.3$   $-1.9$  MeV

CMS  

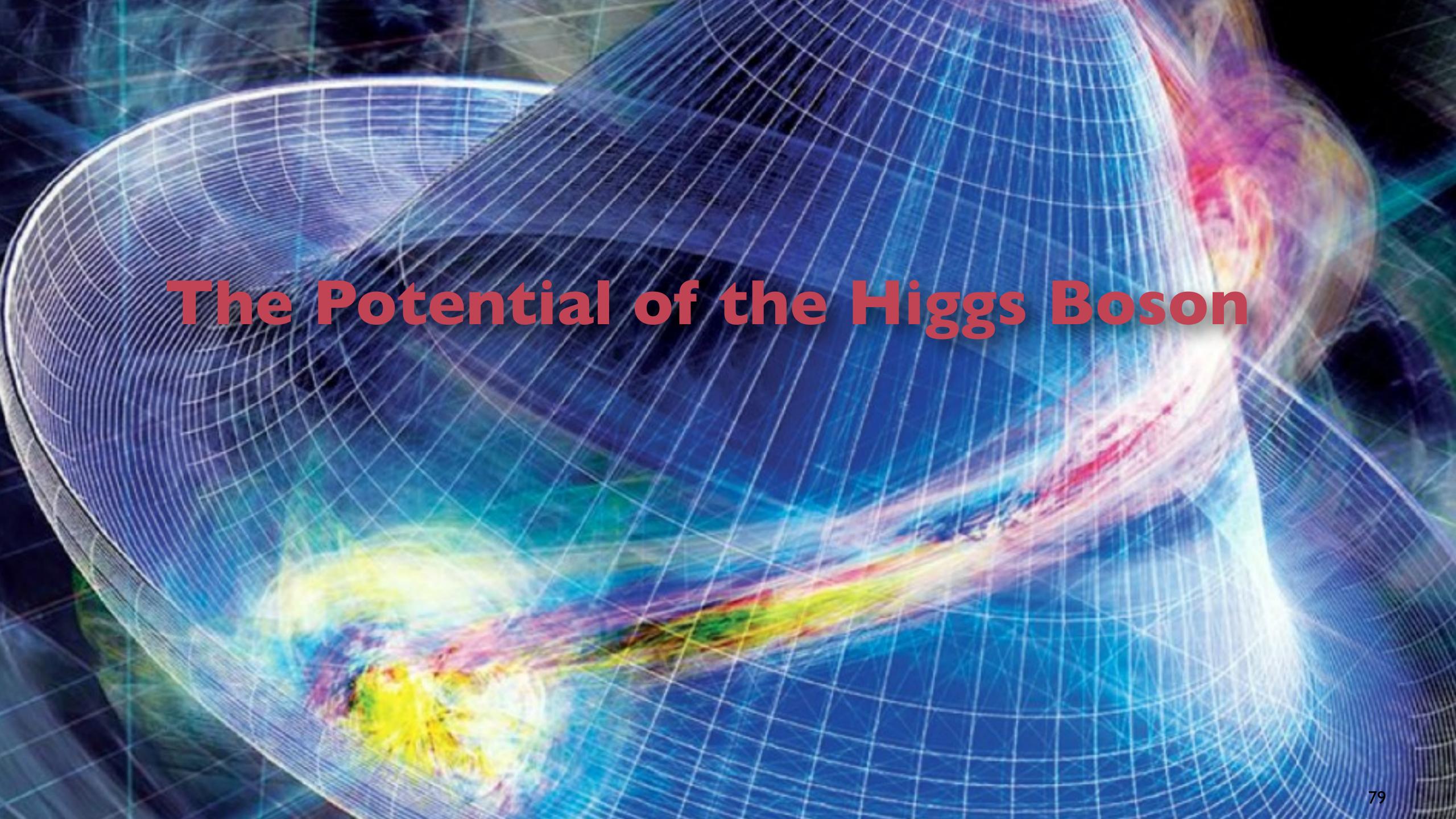
$$\Gamma_{H} = 3.0$$
 +2.0  
 $-1.5$  MeV



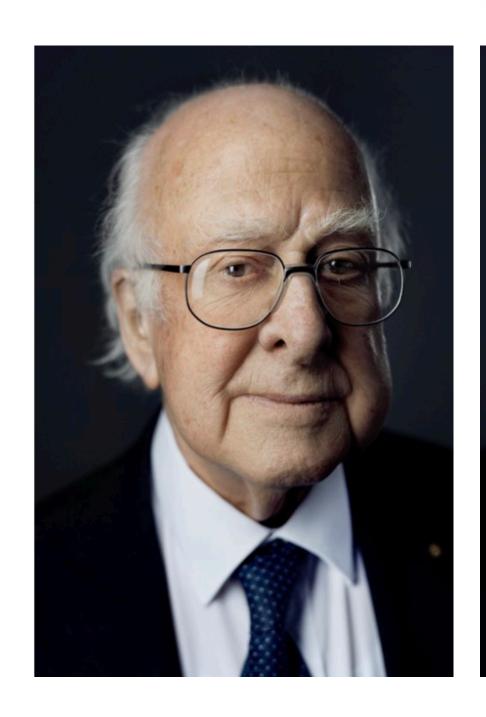
ATLAS Innovation NSBI



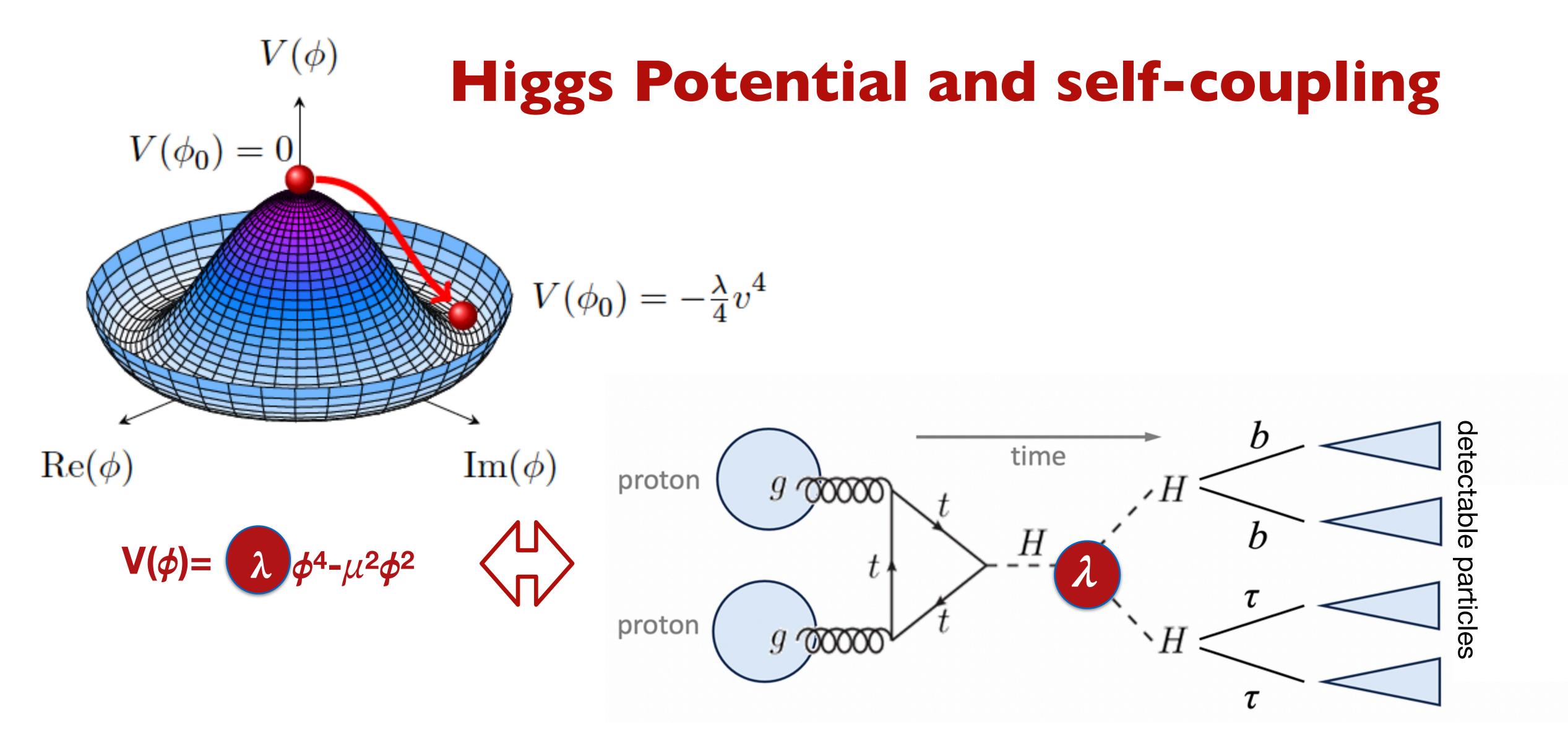
CMS Innovation include new physics in loops



# 2 Higgses

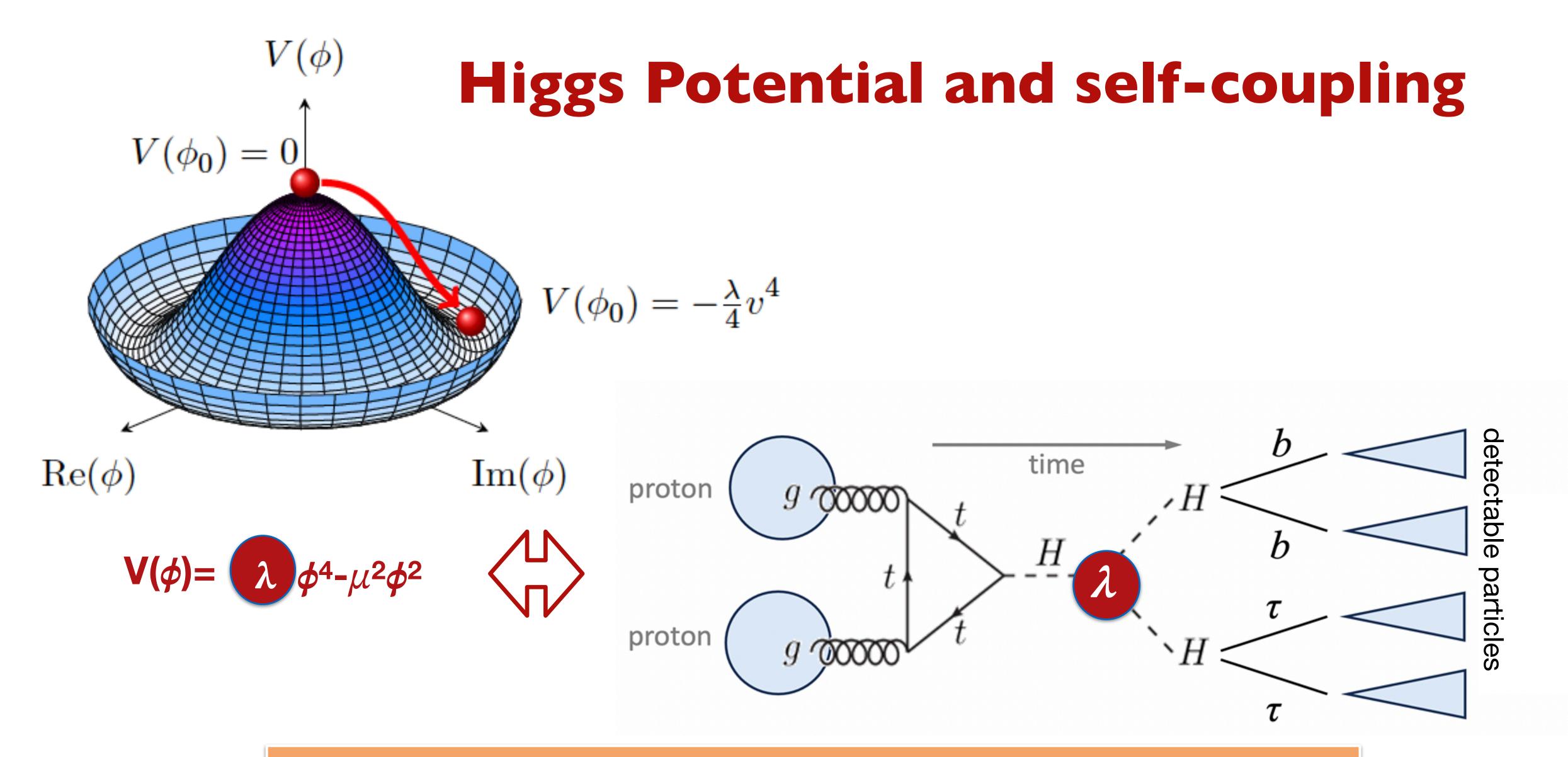






the Higgs potential determines di-Higgs occurrence rate in LHC collisions

6

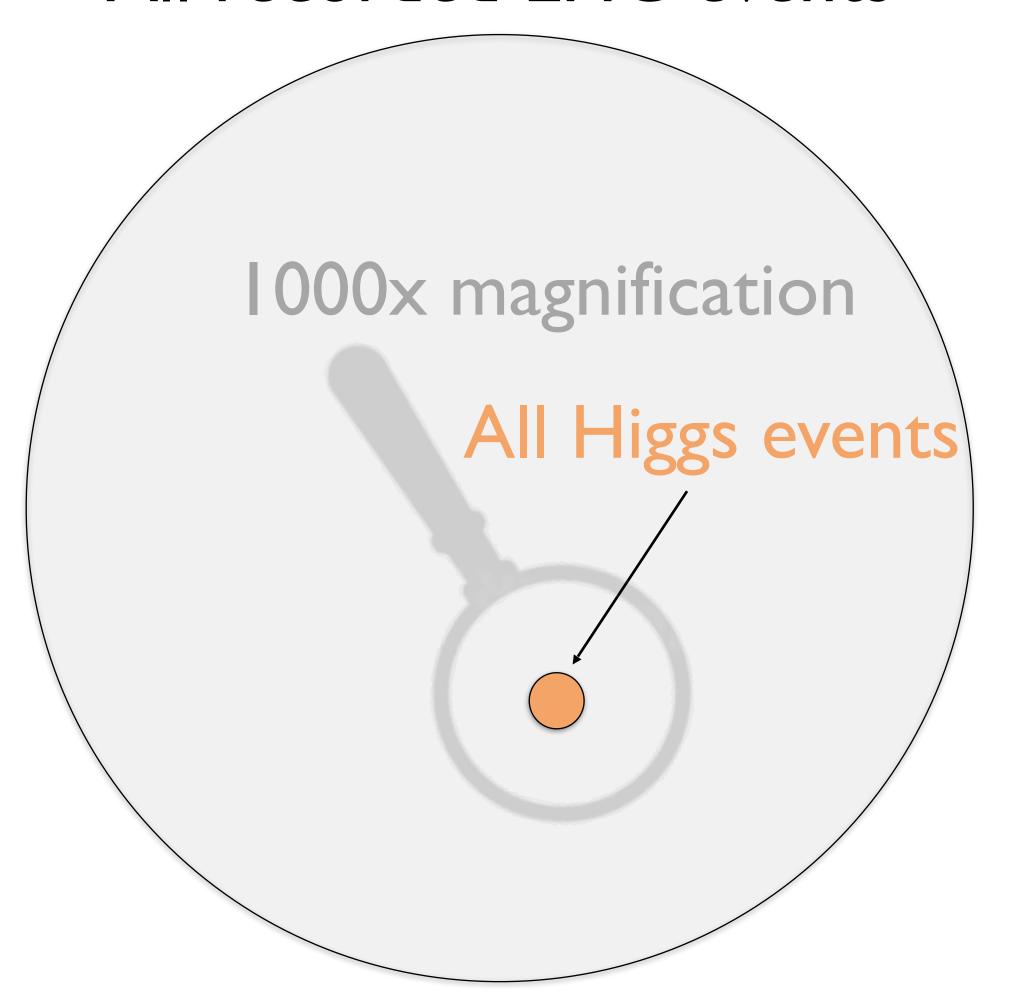


for every di-Higgs event we see a million events that look like it

# Di-Higgs production has not been yet observed because it is very rare.

#### How difficult is it?

#### All recorded LHC events

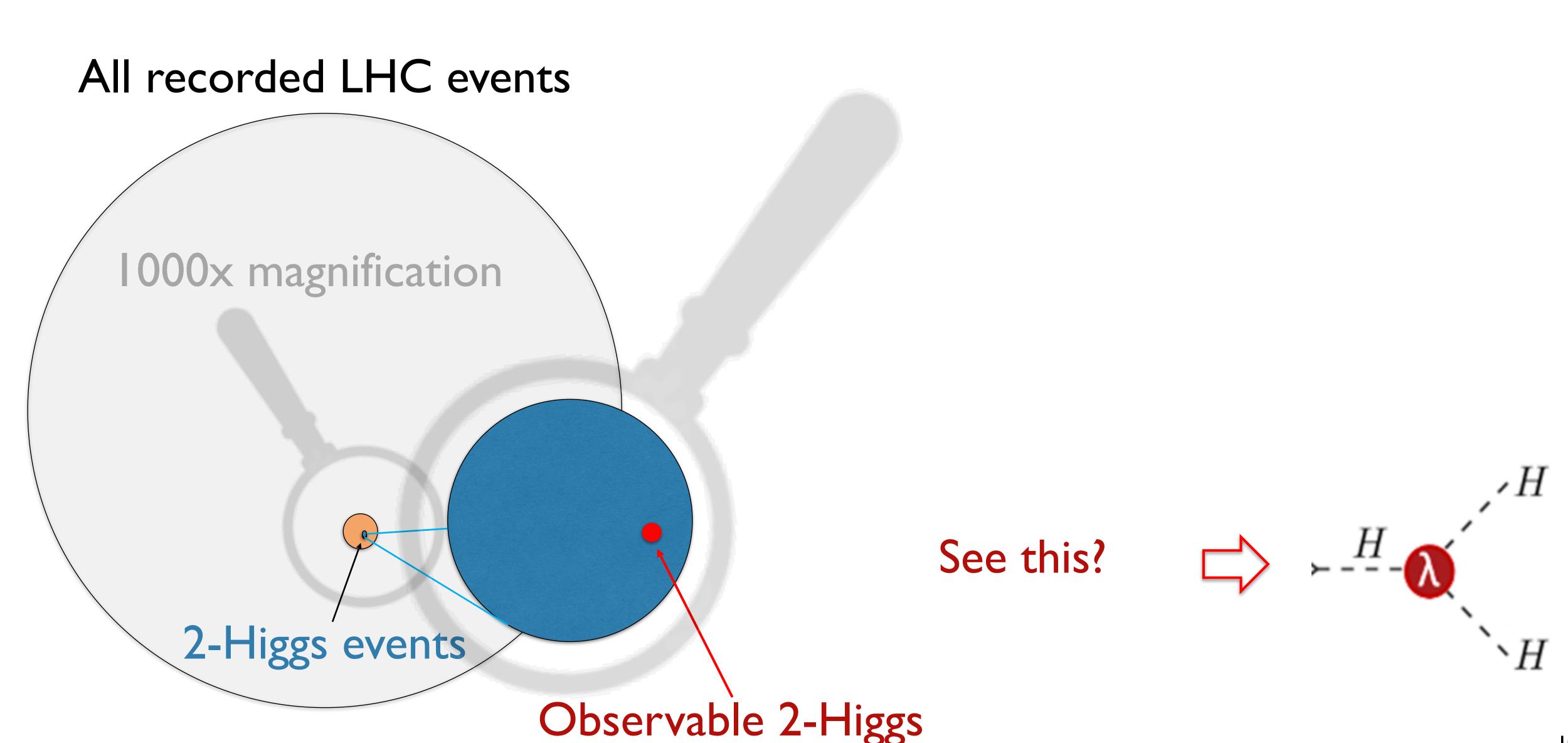


#### How difficult is it?

#### All recorded LHC events

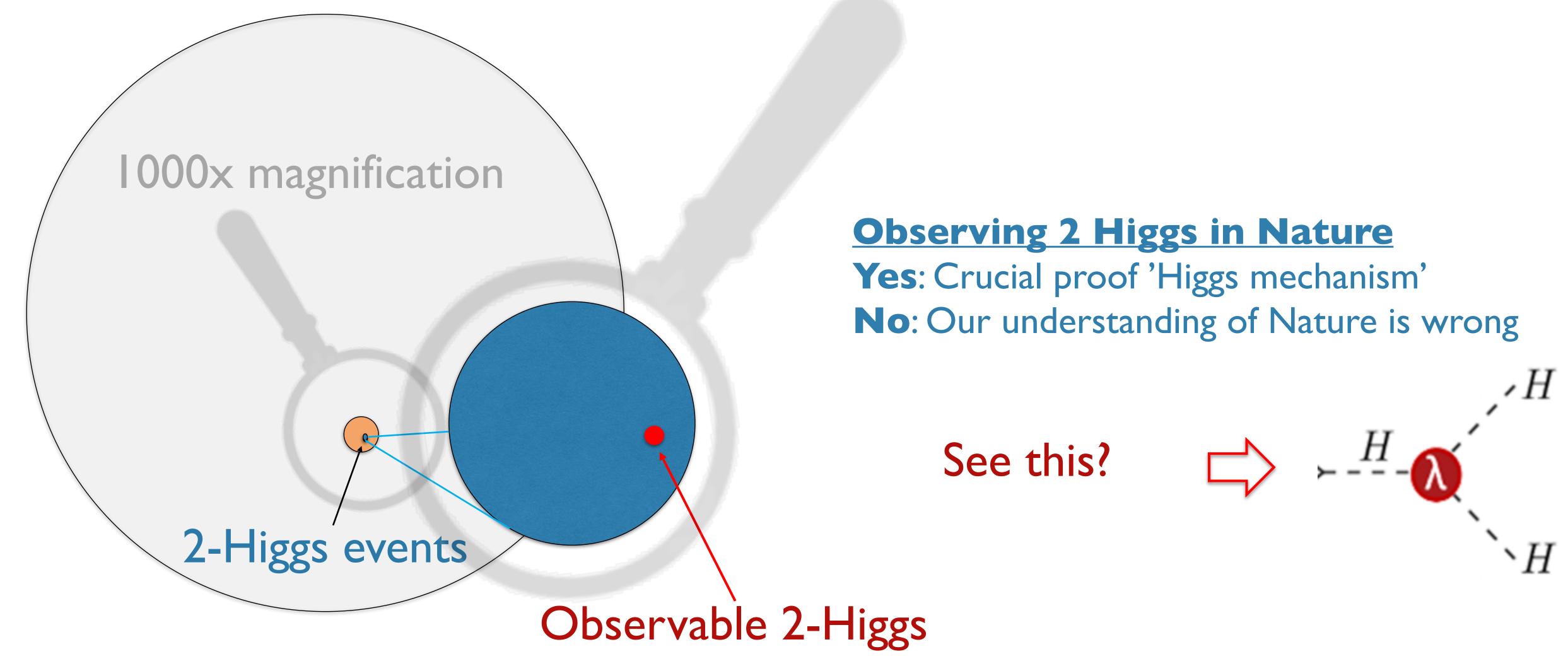


#### How difficult is it?



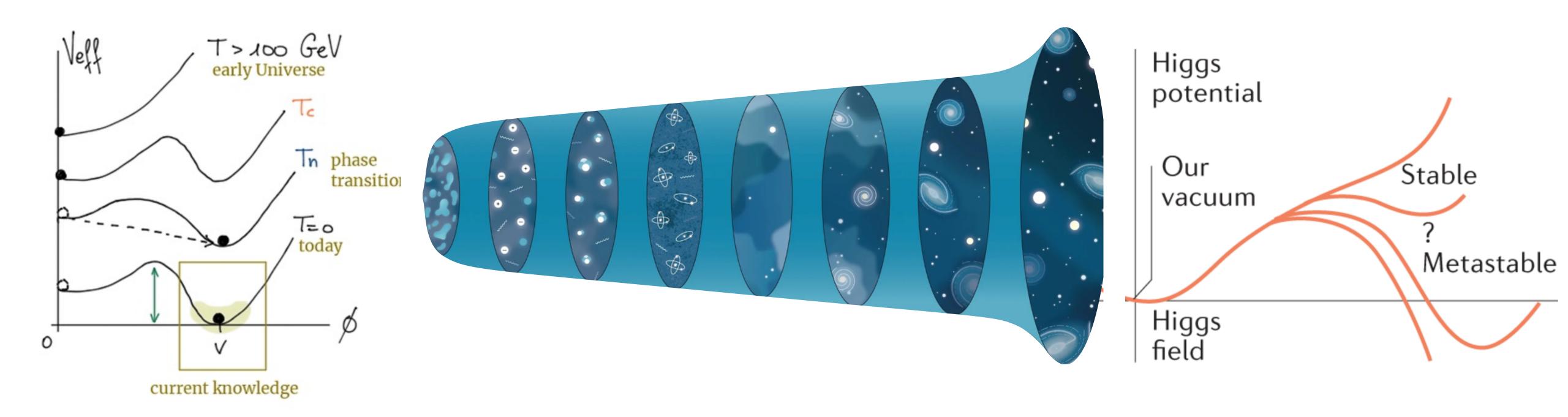
### No evidence of di-Higgs production yet

#### All recorded LHC events



#### Impact of di-higgs observation

# Tremendous impact on cosmology and theoretical physics



disappearance of anti-matter in early universe

stability of universe

#### Understanding the beginning of Universe

A violent EW symmetry breaking

**First Order Phase Transition** 

V = 0  $V \neq 0$  V = 0  $V \neq 0$  V = 0  $V \neq 0$ 

(FOPT)



boiling universe

Why the Universe is made of matter?

The Higgs potential shape is crucial

Ph

ONLY a Strong FOPT can explain a boiling universe

T>100 GeV

In phase

TEO

today

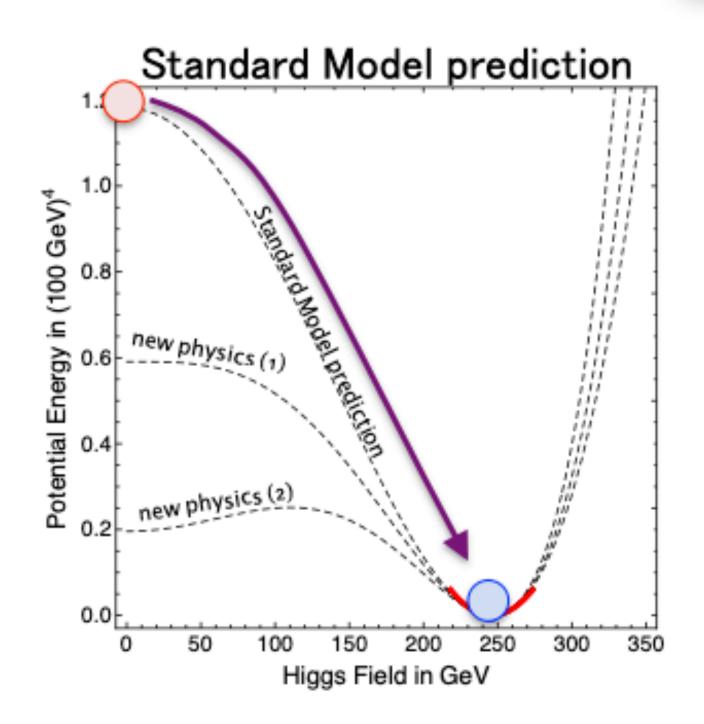
transition

early Universe

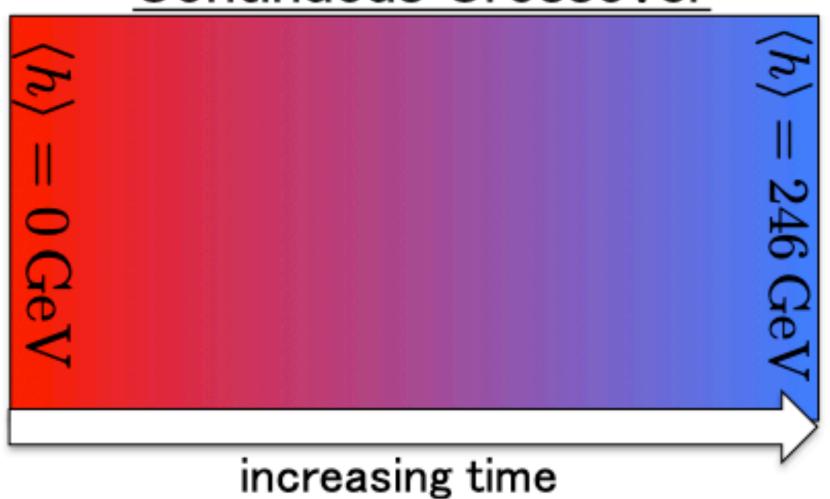
current knowledge

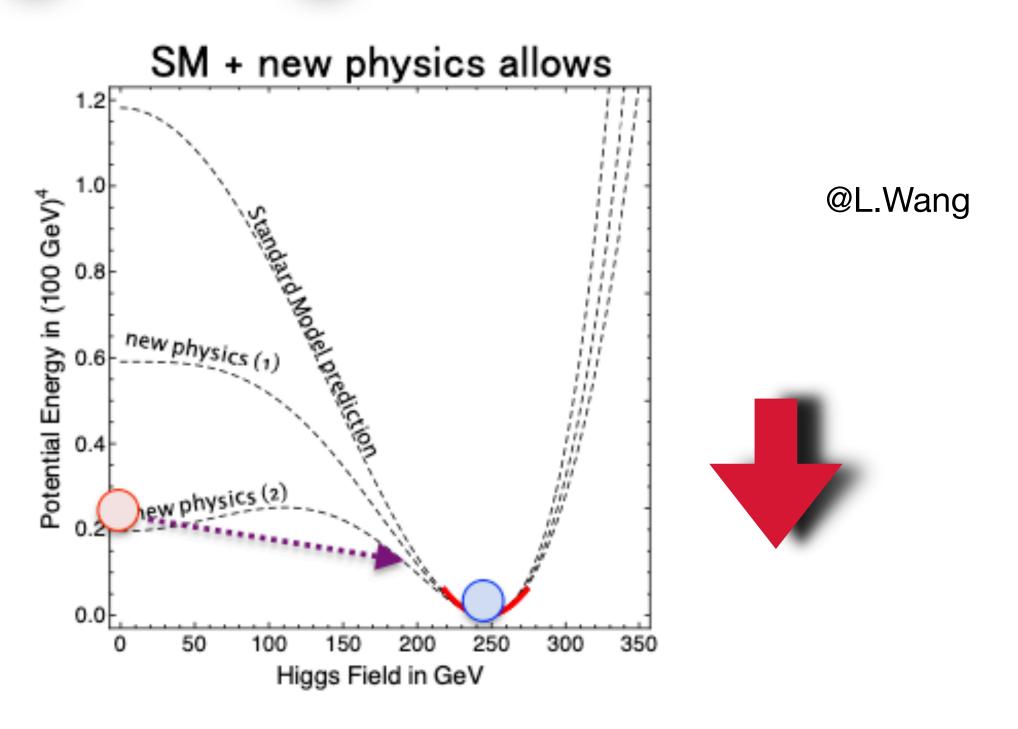
#### Understanding the beginning of Universe

13

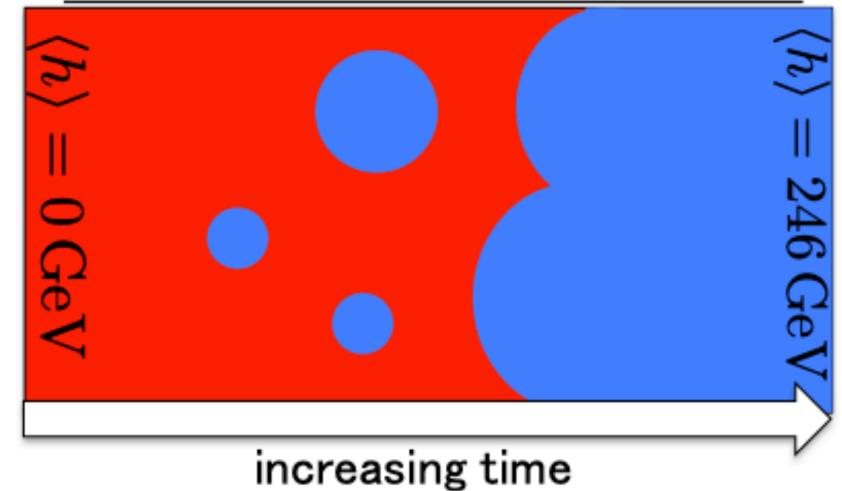


Continuous Crossover





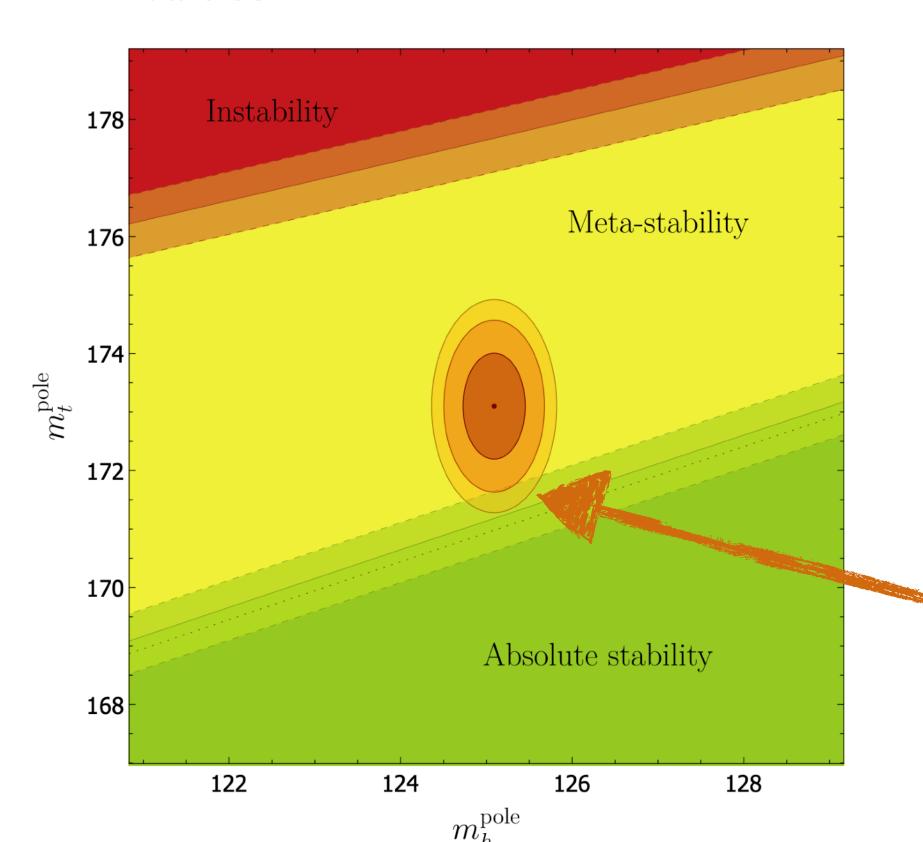
First Order Phase Transition

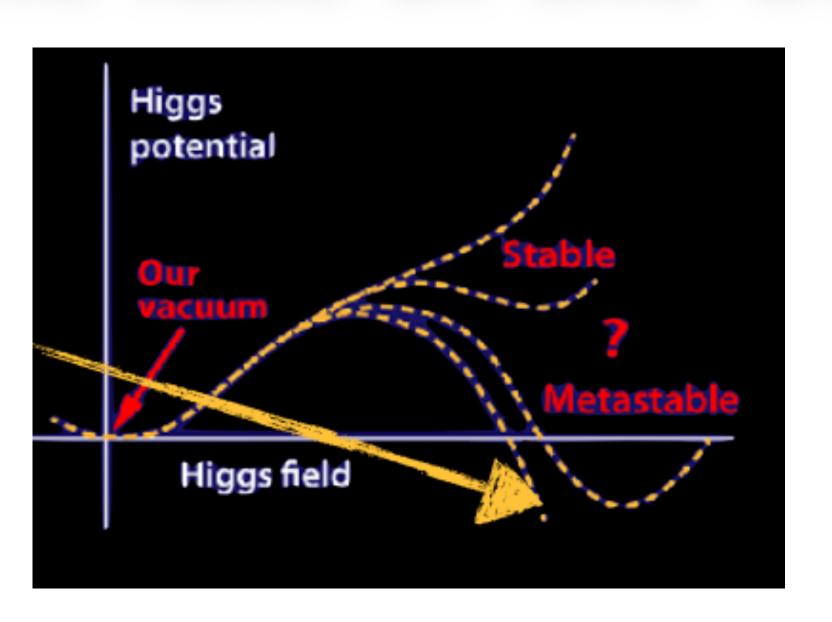


### Understanding the fate of the Universe?

The Higgs field of the SM has a local minimum at a field value of v = 246 GeV.

However, it is possible that a second minimum develops at very large scalar field values





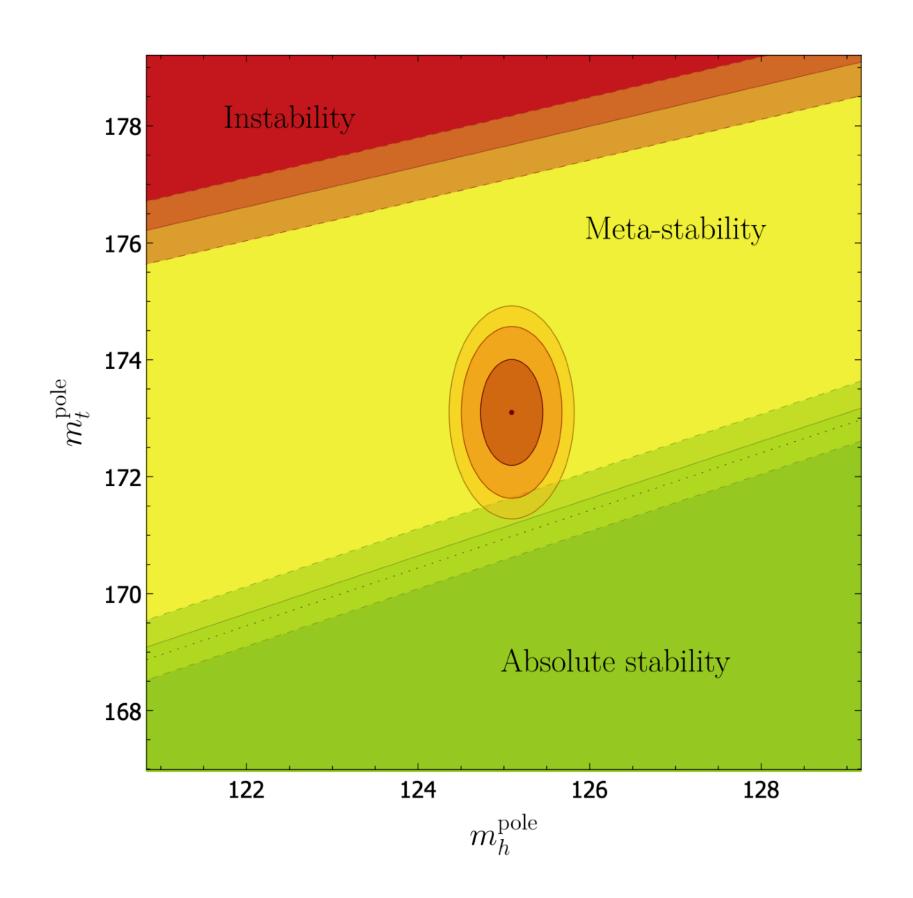
The SM parameters are linked together and from calculations using as input the Higgs mass, top quark mass and the strong coupling constant:

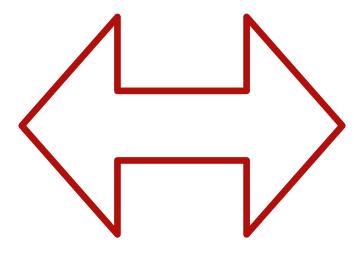
Our Universe appears to be metastable! People often ask me what would happen then?

#### Understanding the fate of the Universe?

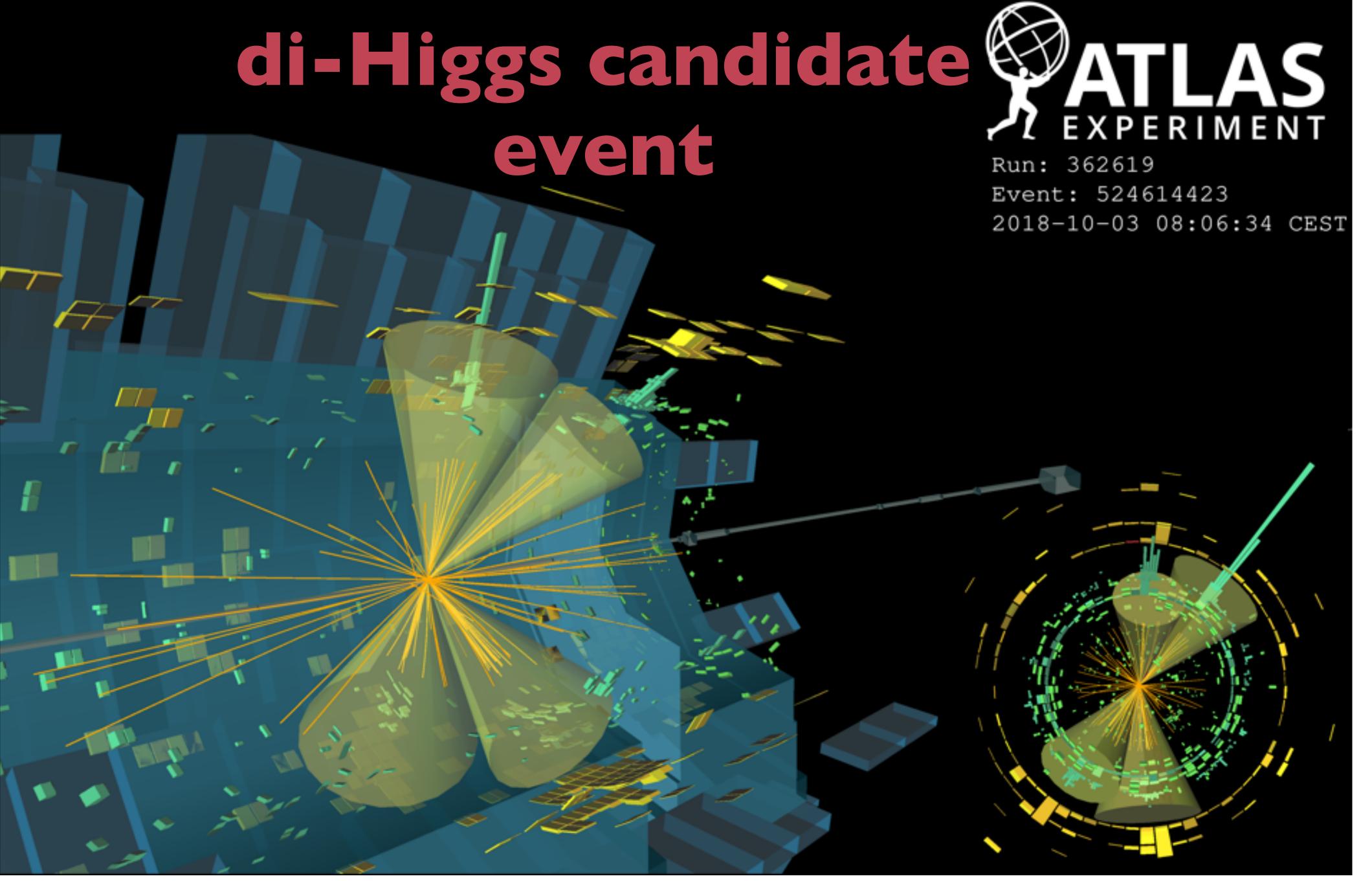
Luckily is the universe will endure a phase transition, this will be in a very distant future much longer than the present life of the universe.

I will be so old by then, that I would prefer this than having to move into a new flat!



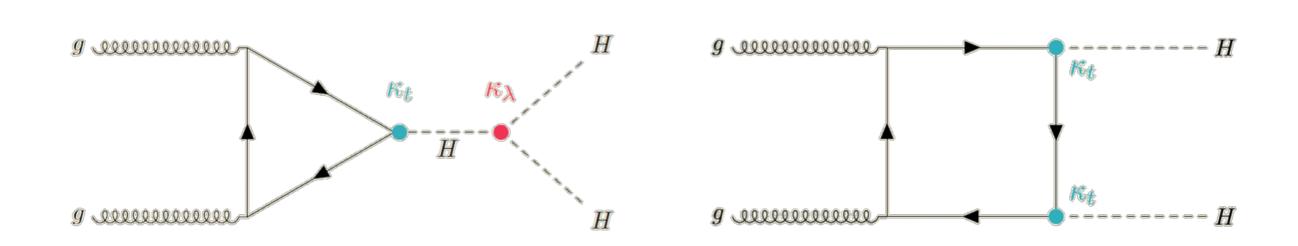




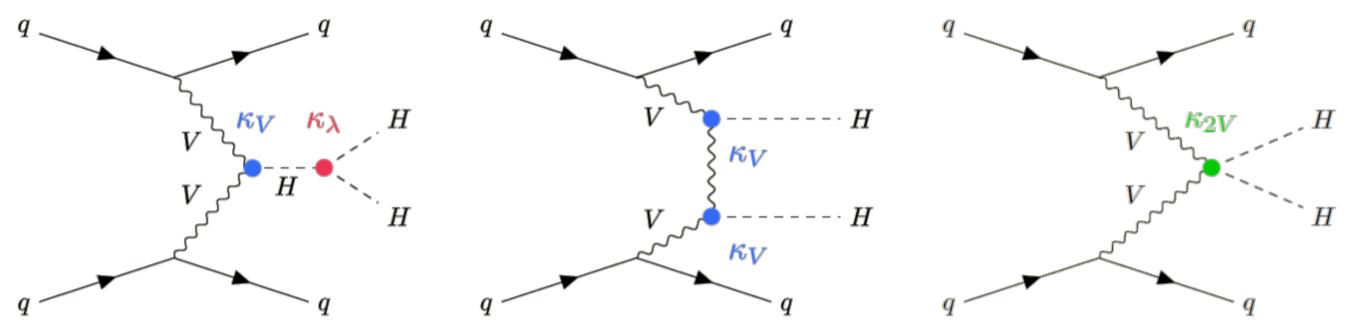


### di-Higgs production at LHC

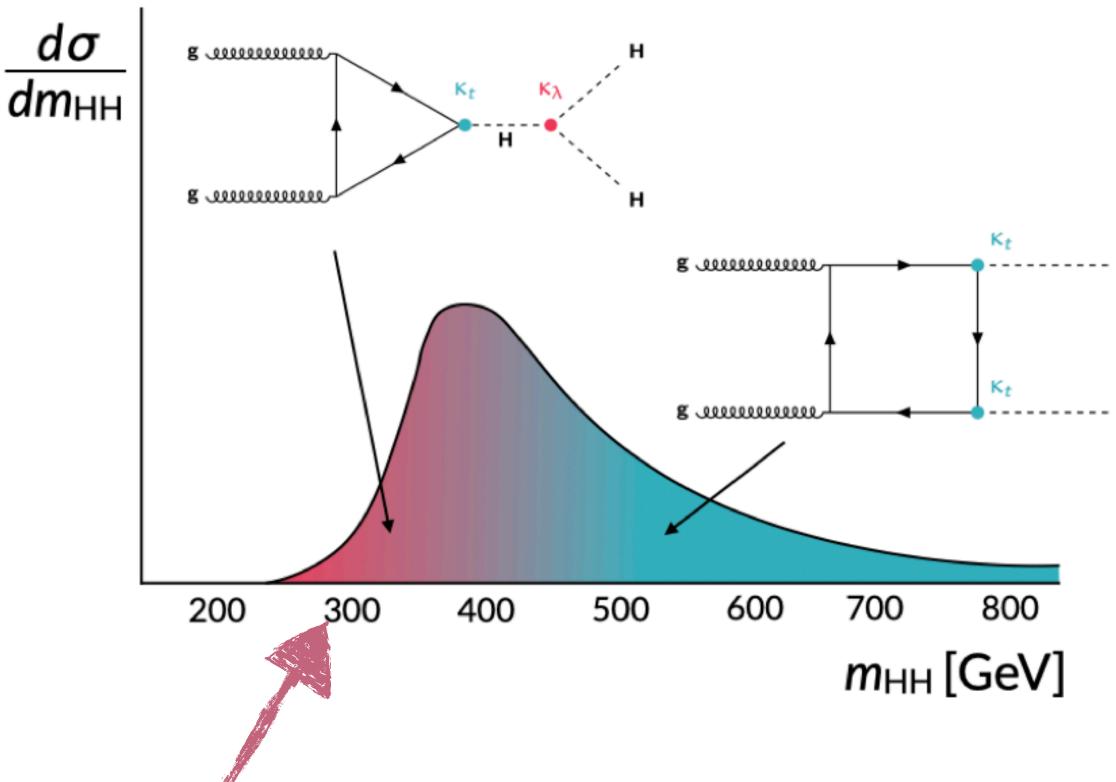
dominant production mode ggF with 2 diagrams that have destructive interference



Vector-boson fusion is the second dominant mode



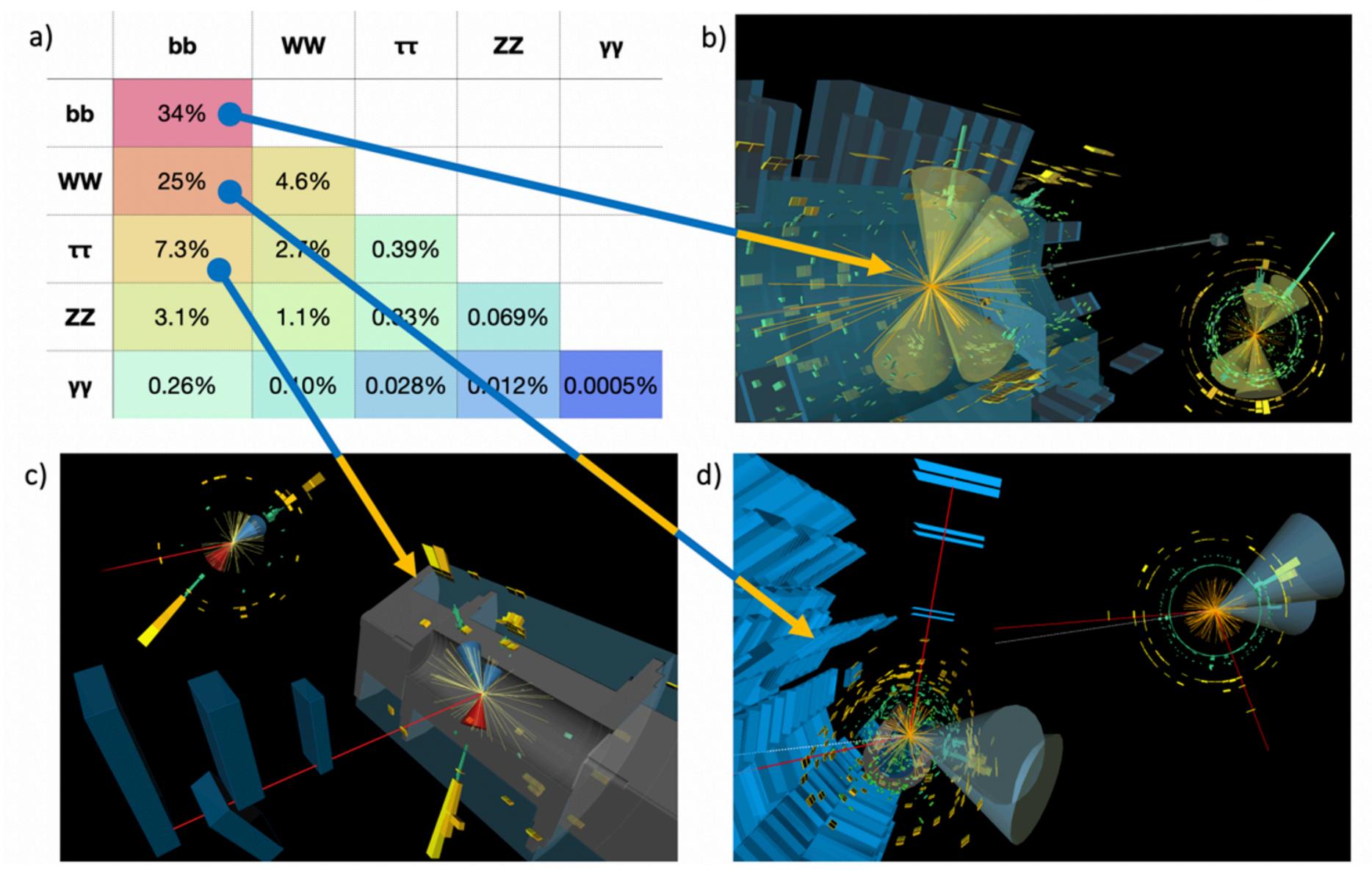
Associated productions, HHV HHtt have much smaller production cross-sections



Triangle diagram contributes mostly at low m<sub>HH</sub>

 $\rightarrow$  Analyses that contribute to that region are most sensitive to  $\kappa_{\lambda} = \lambda/\lambda_{SM}$ 

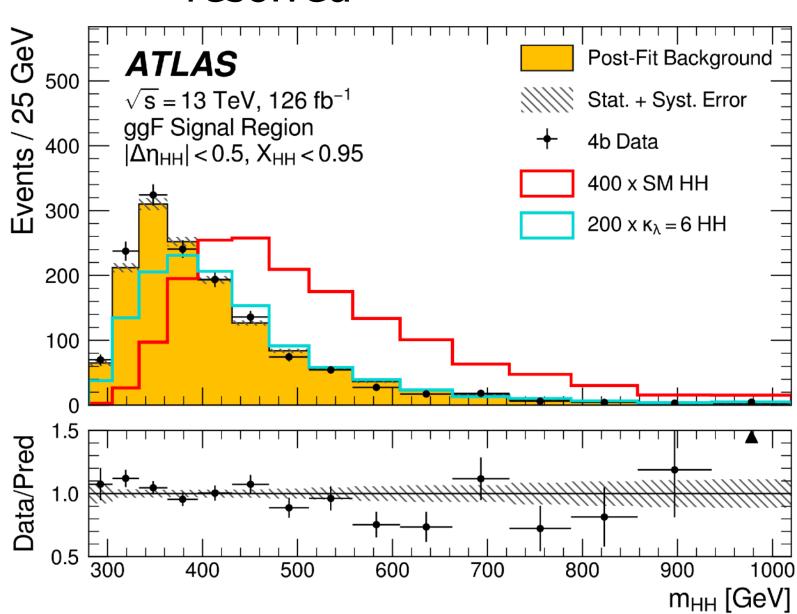
### Different final states investigated

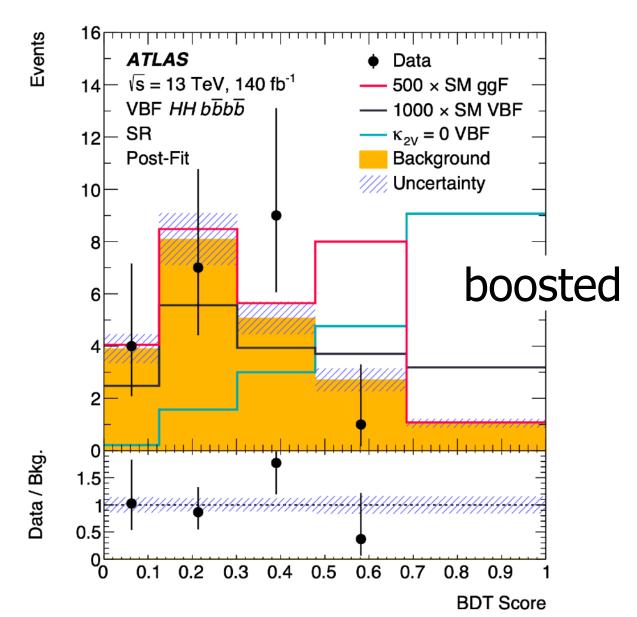


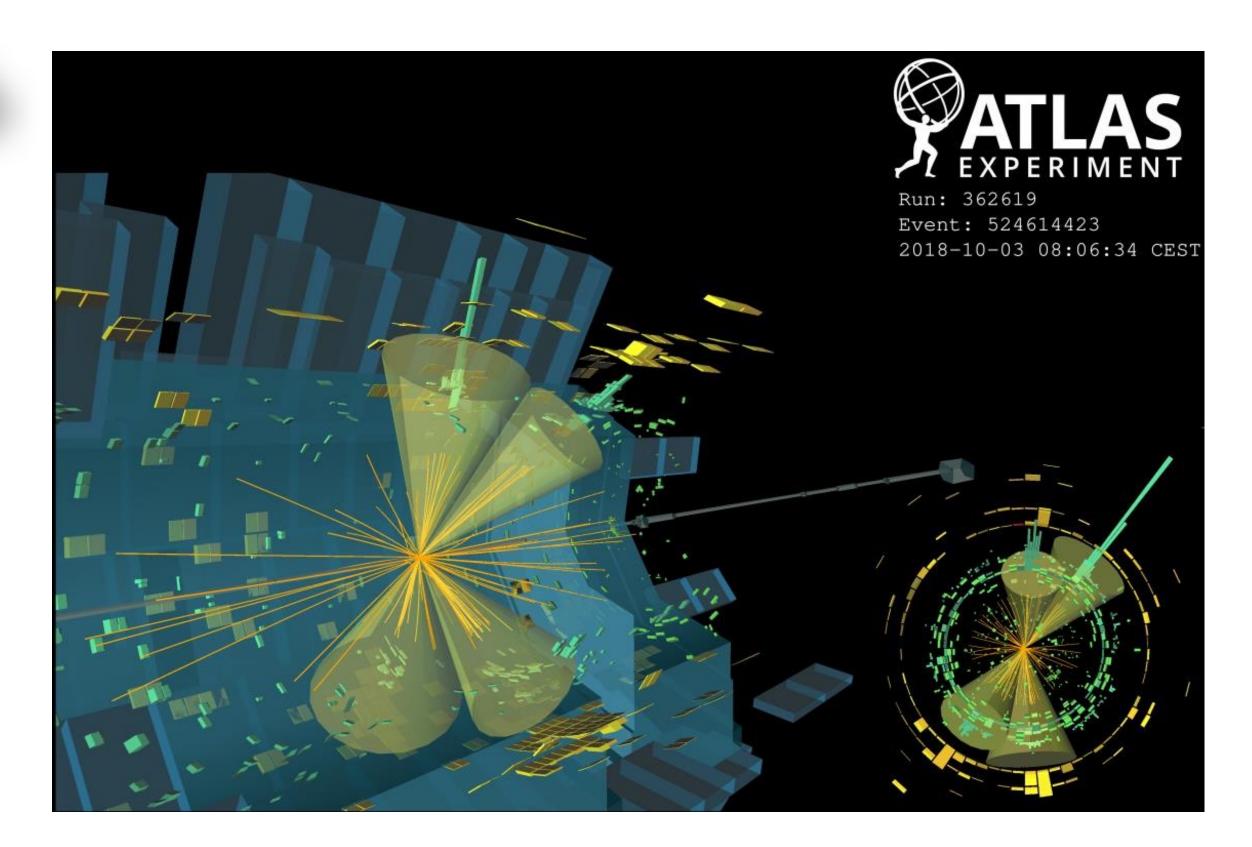


#### $HH \rightarrow bbbb$

#### resolved







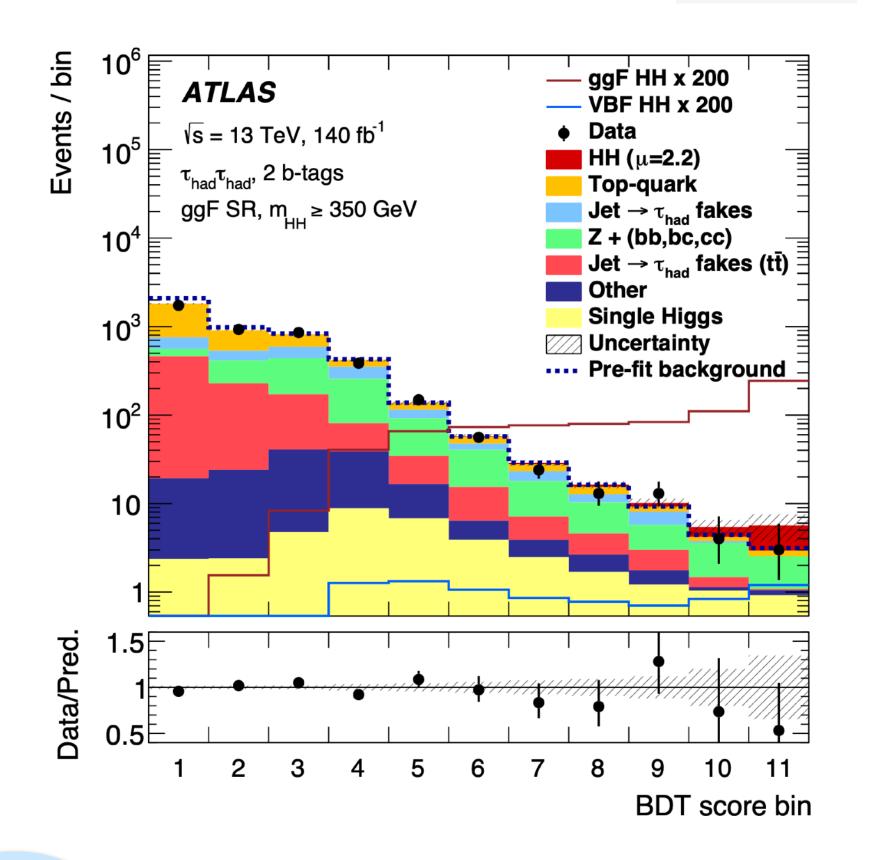
Highest BR, sensitive to high pT of the Higgs

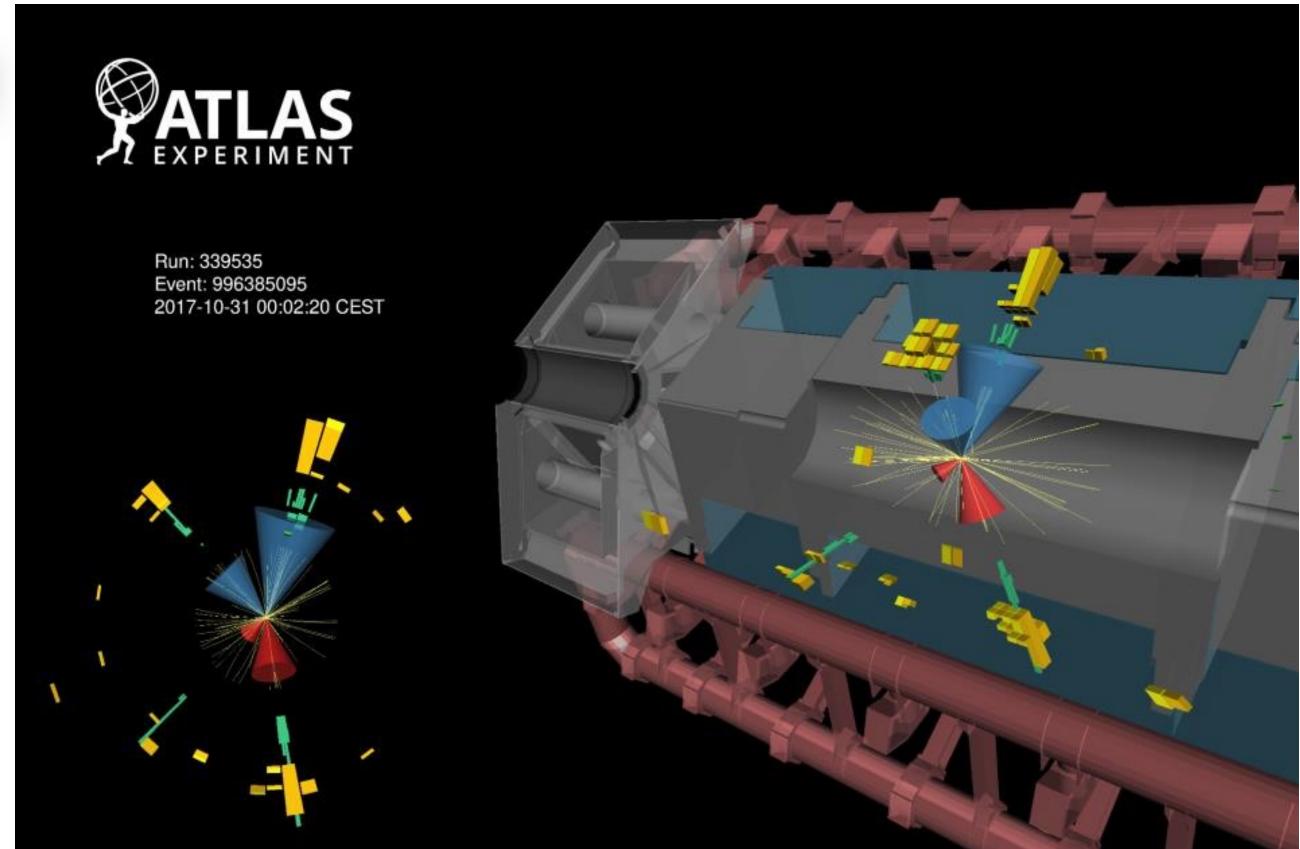
The sensitivity of the analyses is improved relative to previous iterations by using more sophisticated background modeling techniques, event categorization and improved jet reconstruction and flavor identification algorithms, in addition to the increased integrated luminosity of the analyzed data.

both in  $T_{had}T_{had}$  and  $T_{lep}T_{had}$  channel and in ggf+VBF production



arXiv:2209.10910



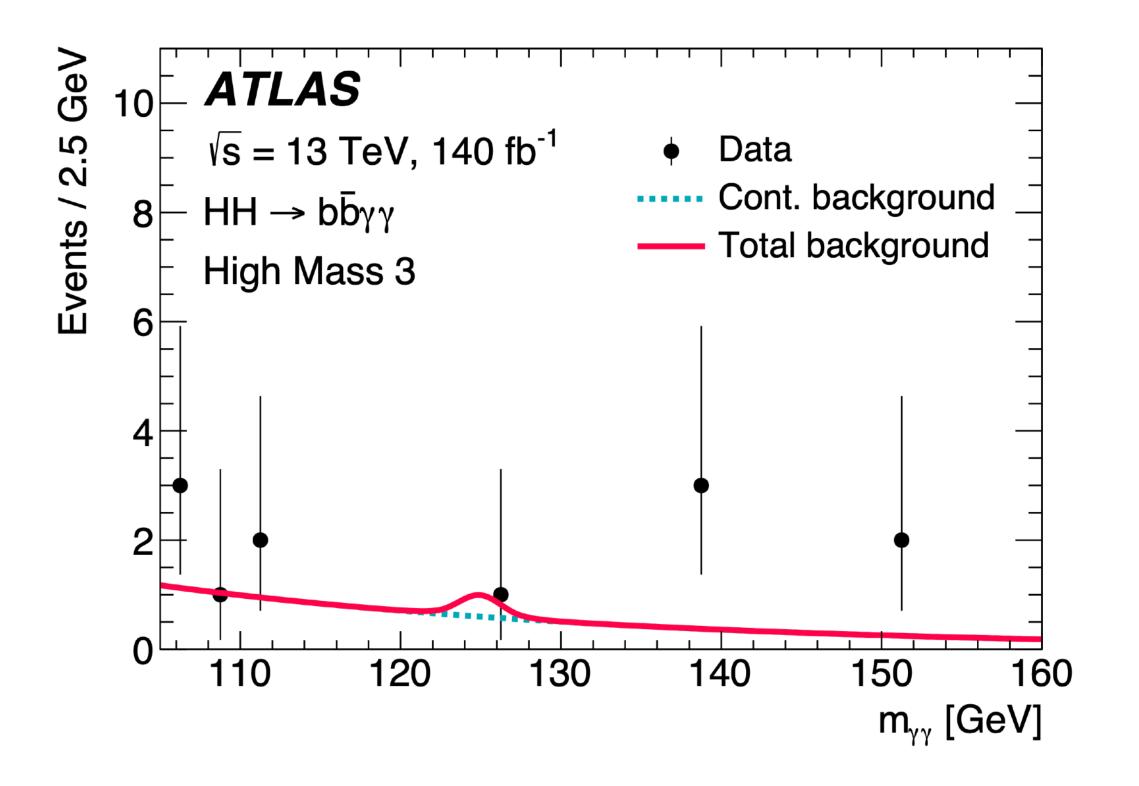


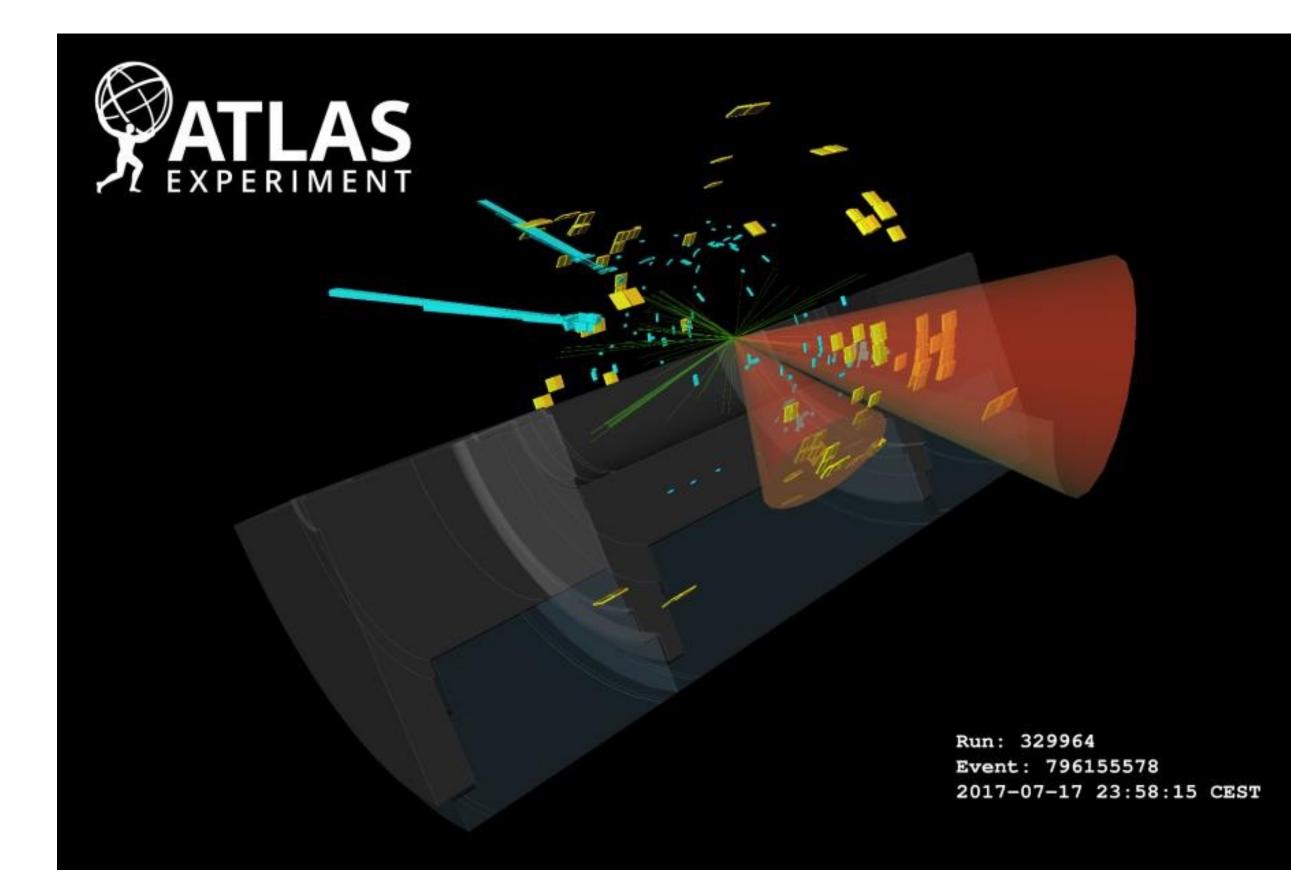


sensitivity gains due to significant improvements in the  $\tau$ had-vis and b-jet reconstruction and identification.



#### $HH \rightarrow bb\gamma\gamma$

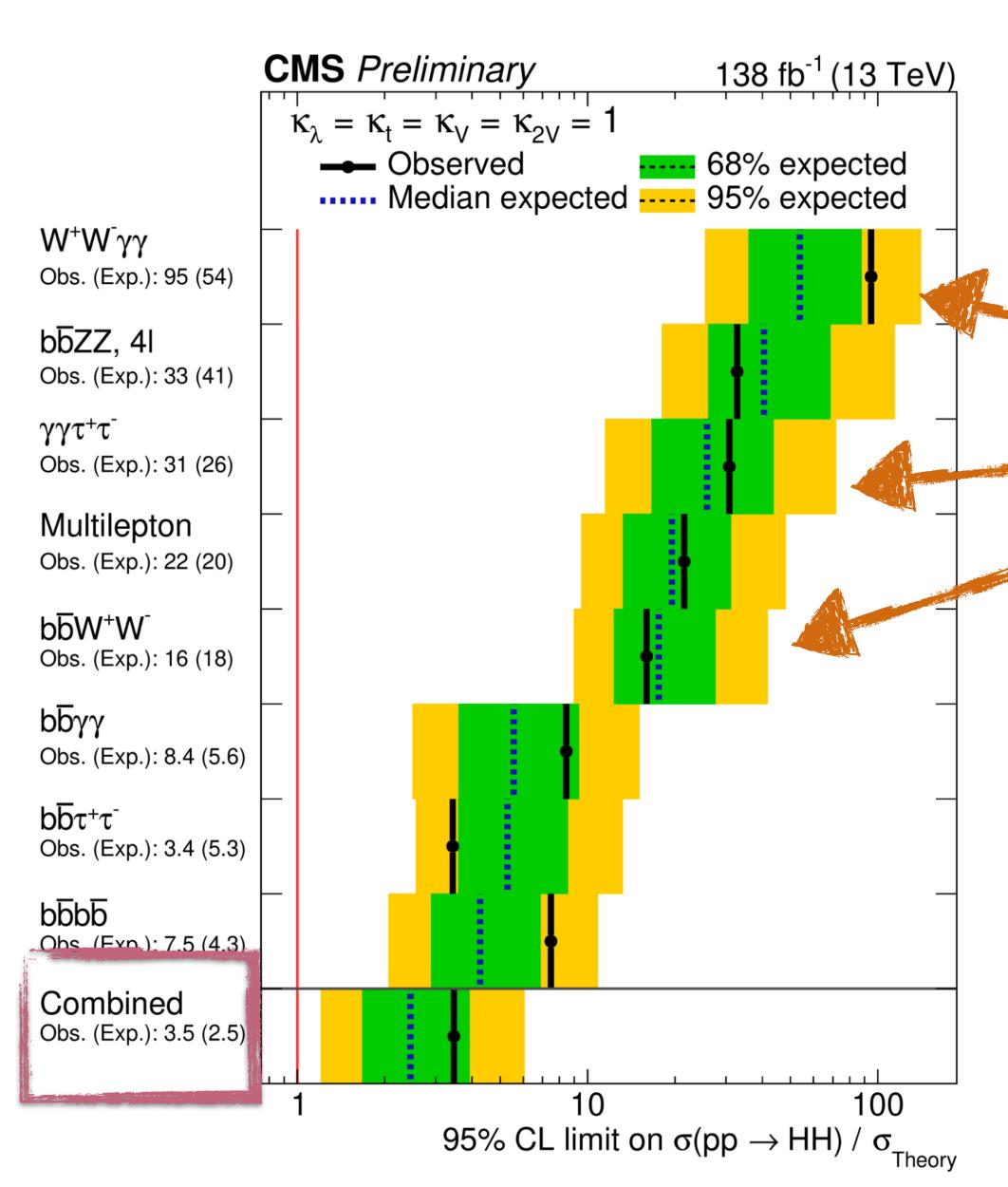




incorporates a categorization based on mbb γγ and multivariate event selections

### di-Higgs in CMS

#### CMS-PAS-HIG-20-011



Inputs: bbbb resolved and boosted, bb $\tau\tau$ , bb $\gamma\gamma$ 

New wrt to Nature combination!

WWγγ

γγττ

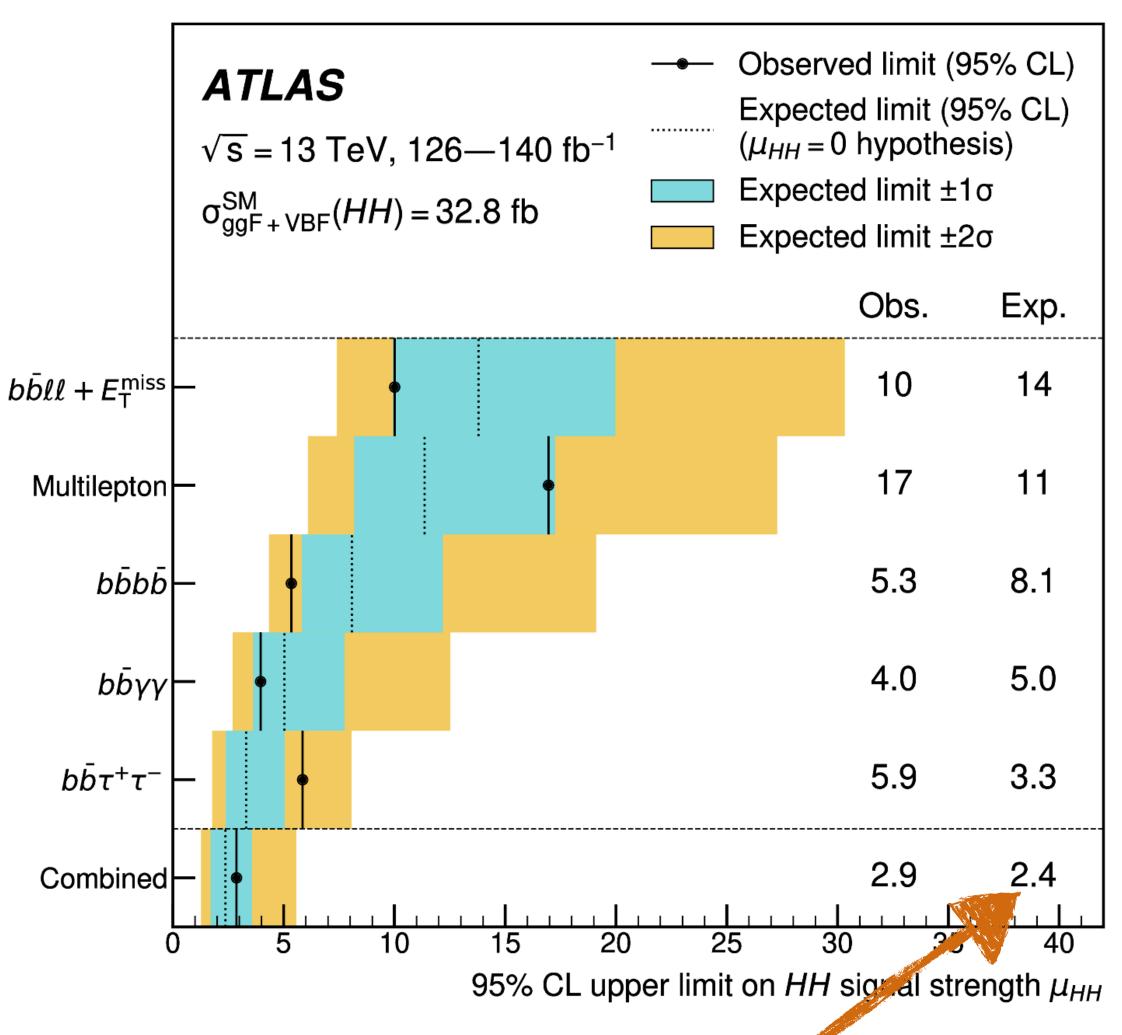
bbWW the only existing analysis in this channel

Most complete HH combination @ CMS x-section sensitivity at 3.5 xSM ( 2.5 exp)

**Observed** :  $\kappa_{\lambda} \in [-1.39, +7.02]$ 

Expected  $\kappa_{\lambda} \in [-1.02, +7.19]$ 

### di-Higgs ATLAS



bbγγ

bbττ

bbbb (resolved jets)

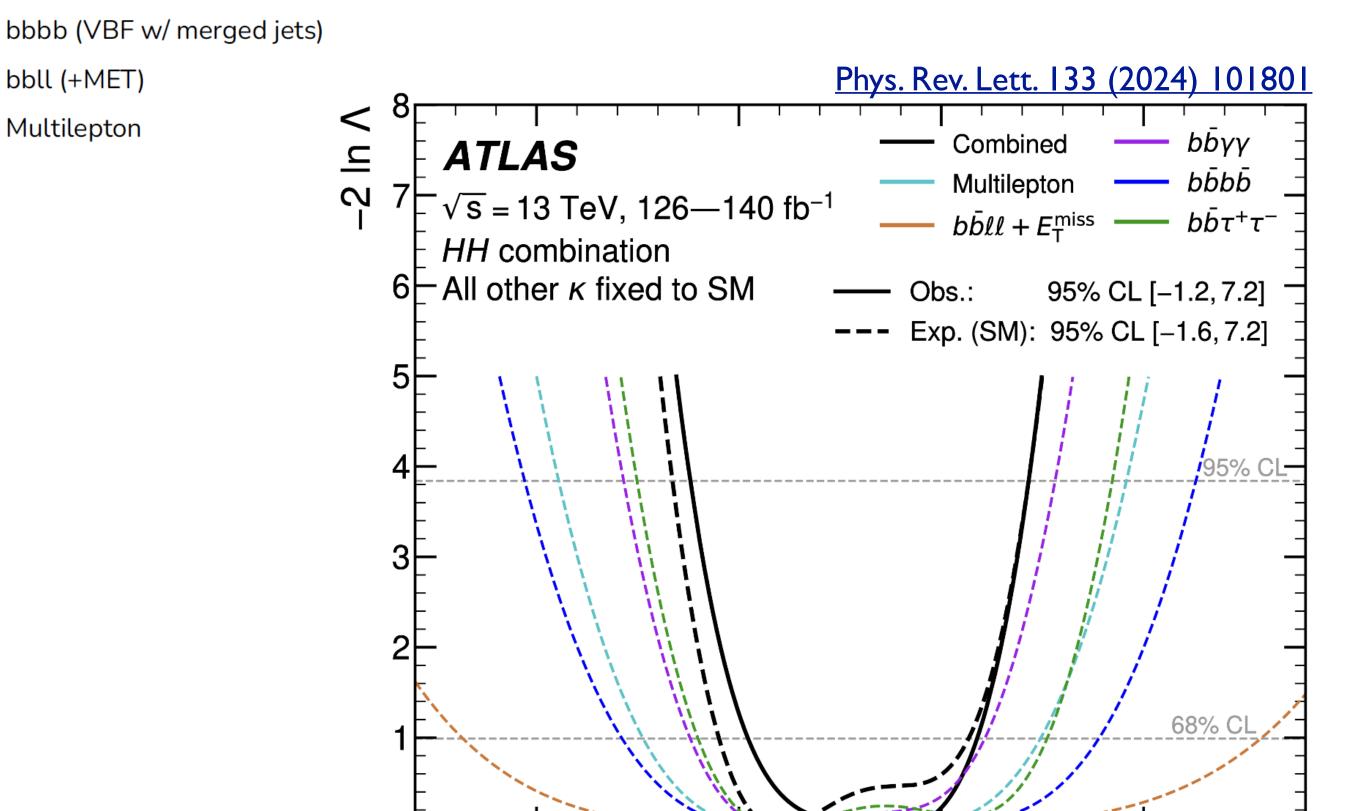
bbll (+MET)

Multilepton

best sensitivity to date

**Observed** :  $K_{\lambda} \in [-1.2, +7.2]$ 

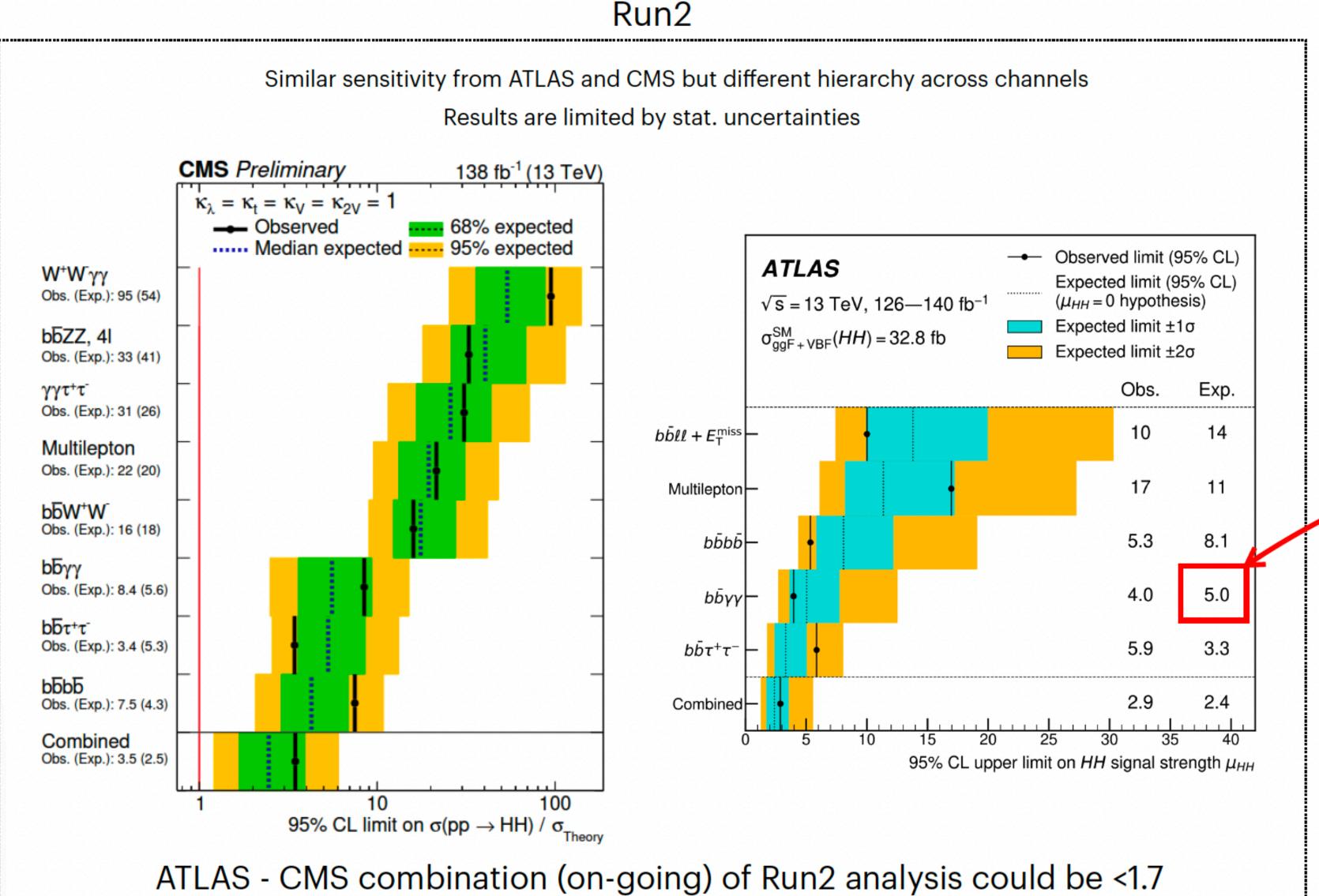
Expected  $\kappa_{\lambda} \in [-1.6, +7.2]$ 



bbyy channel most sensitive, and bbtt at negative  $K_{\lambda}$ 

 $K_{\lambda}$ 

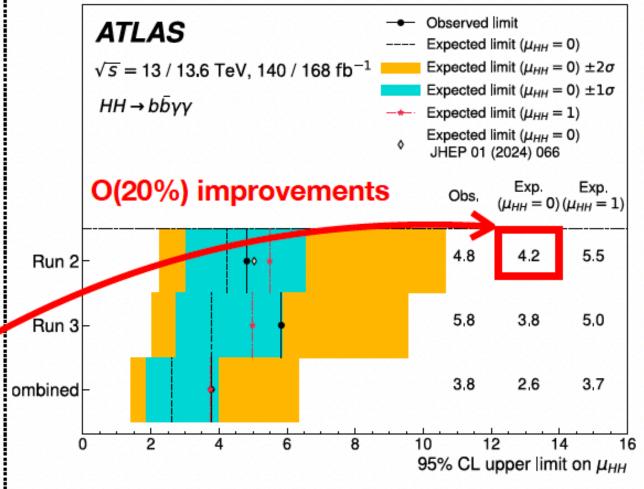
#### ATLAS+CMS



We can expect it to reach ~1 with new improvements, given 500/fb (Run1+Run2+Run3)

#### Run3 getting in the game

Opportunity to maximise the analysis sensitivity (data streams, triggers, object reconstruction, analysis techniques)



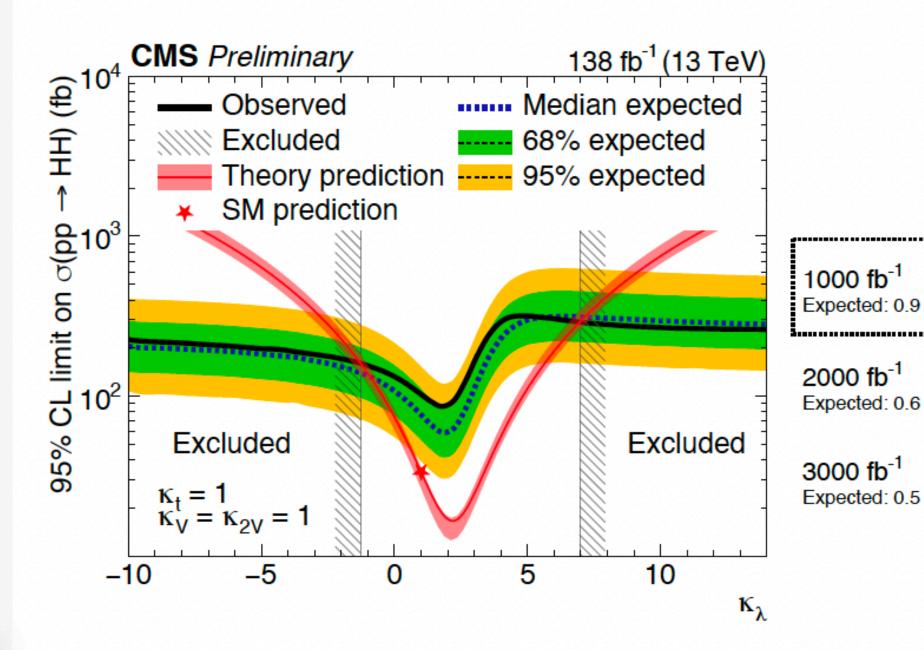
#### ATLAS+CMS

CMS Projections Preliminary

..... Median expected

 $\kappa_{\lambda} = \kappa_{t} = \kappa_{V} = \kappa_{2V} = 1$ 

Sensitivity to exclude of  $\lambda_{HHH}$  = 0 by the the end of Run3

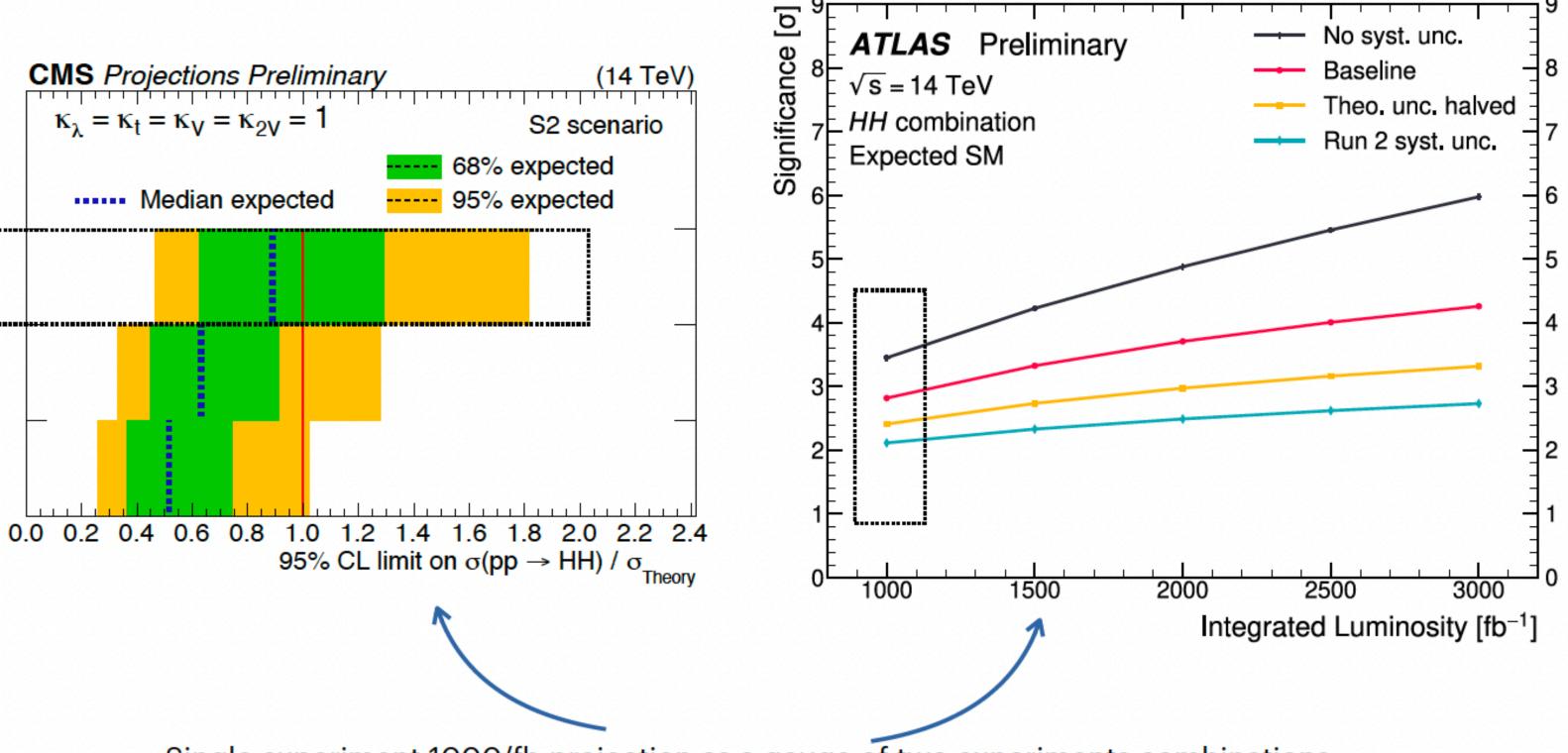


95% CL sensitivity for HH production combining ATLAS and CMS. Sensitivity with single experiment?

68% expected

95% expected

HH evidence at the end of Run3 combining ATLAS and CMS?

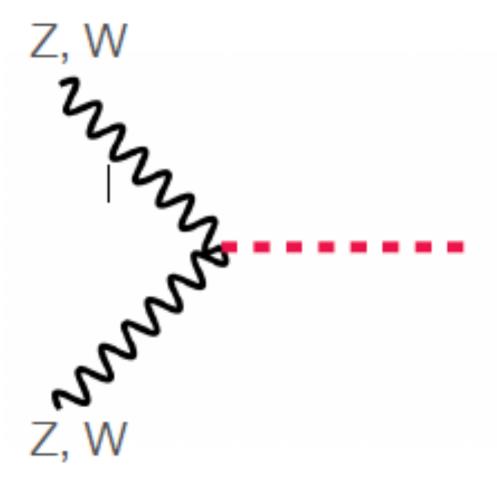


Single experiment 1000/fb projection as a gauge of two experiments combinations

#### Indirect determination of $\lambda_3$ (via single Higgs)

The determination of the Higgs self coupling can be made in a direct way, observing the di-Higgs process (which is the best way of assessing the presence of new physics deviations)

Or indirectly via single Higgs



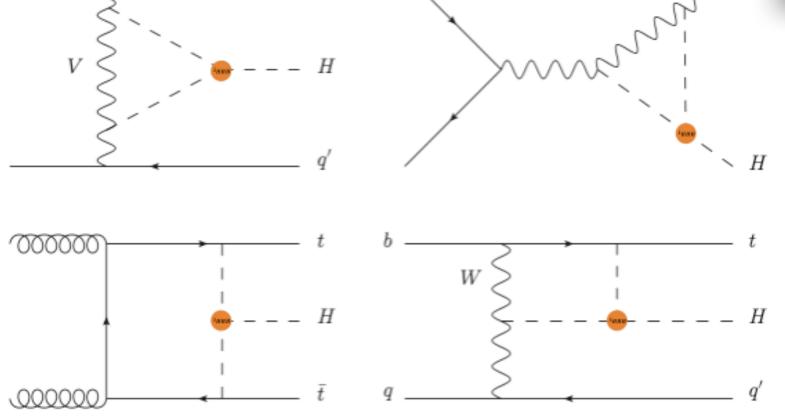
New physics can often induce changes in the hZ coupling

Single Higgs and di-higgs combination

arXiv:2407.13554

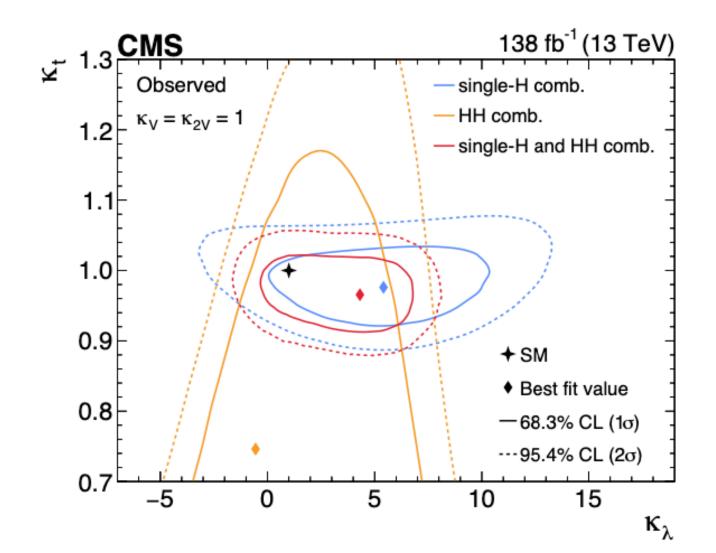
138 fb<sup>-1</sup> (13 TeV)

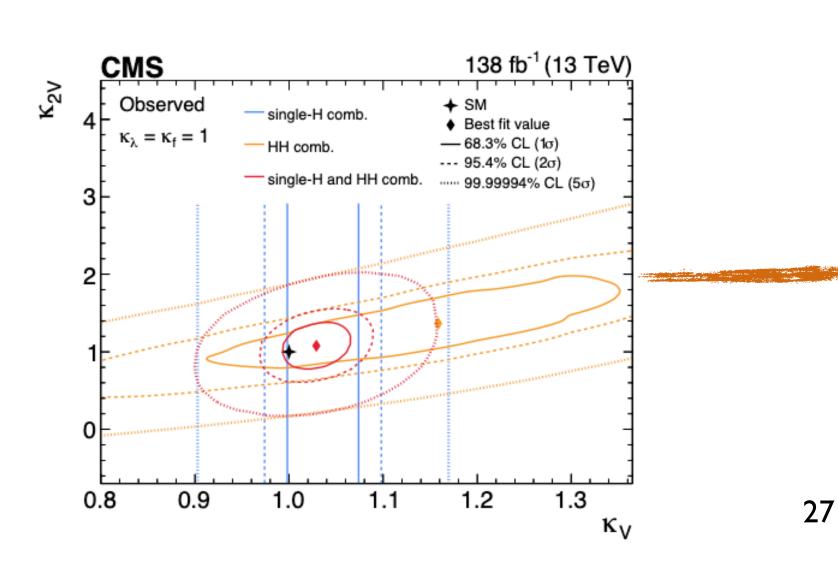
single-H comb., 5.8<sup>+3.5</sup><sub>-4.2</sub>

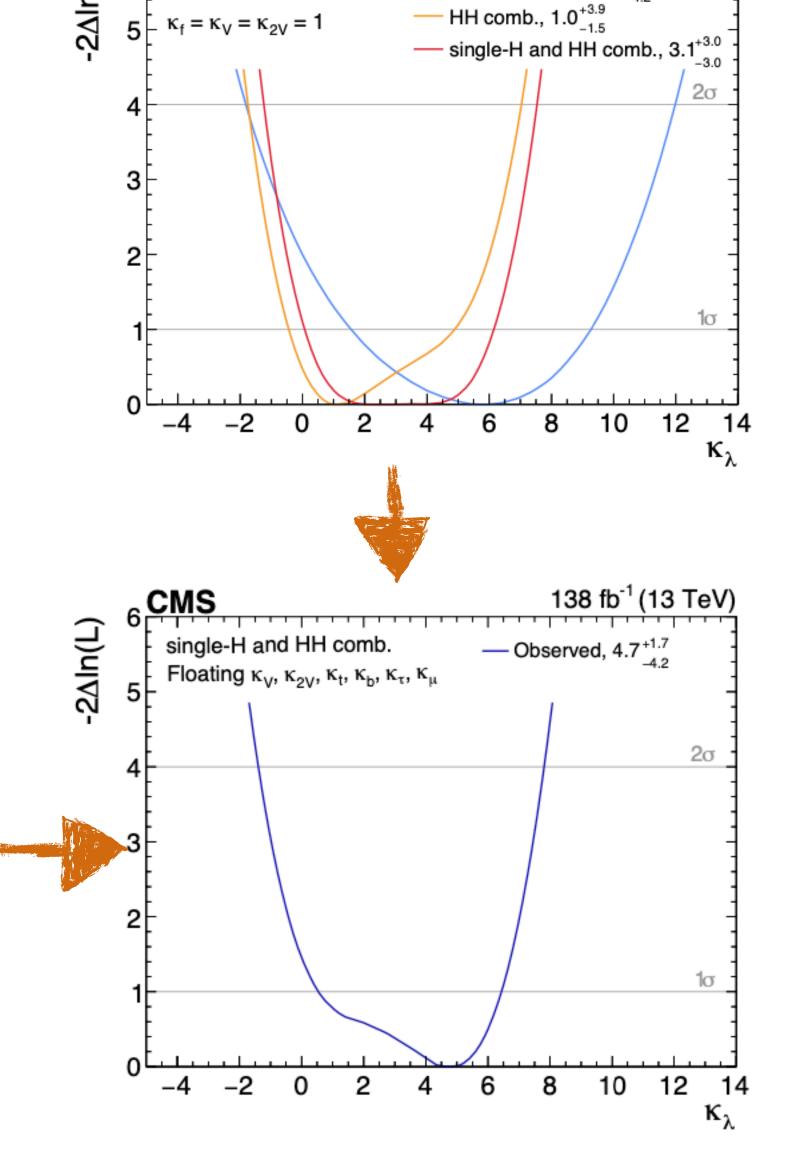


the single-H boson production includes processes sensitive to the higgs self couplings at the NLO EW correction level.

At the LHC, Single Higgs adds little on di-Higgs combination, but it helps because we can constrain the assumptions on  $k_V$ ,  $k_{2V}$ , and  $k_t$  in the fits



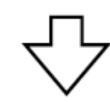




### hZZ vs Higgs self-coupling

$$\frac{1}{\Lambda^2}(H^{\dagger}\partial H)^2$$

 $\frac{1}{\Lambda^2}(H^{\dagger}H)^3$ 



 $\delta_{Zh}$ 

Modify H-Z coupling  $\Rightarrow \delta_{Zh}$ 

Modify Higgs self-coupling  $\Rightarrow \delta_{\lambda_3}$ 

However,  $\delta_{Zh} \propto {
m g}_z$  , while  $\delta_{\lambda_3}$  is not related to  $\lambda_{3,{
m SM}}$ 

With some tuning, one can find models in which  $\delta_{\lambda_3} > \delta_{Zh}$ 

gz = coupling strength of theZ boson to other particles.

$$g_Z = rac{g}{\cos heta_W}\,,$$

g is the weak isospin coupling (the  $SU(2)_L$  coupling),  $\theta_W$  is the weak mixing angle (Weinberg angle).

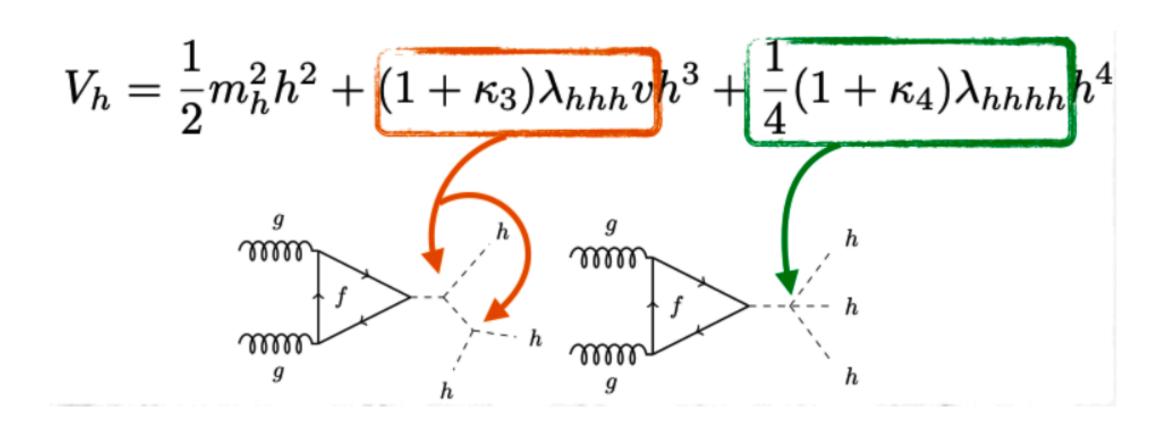
# 3 Higgses



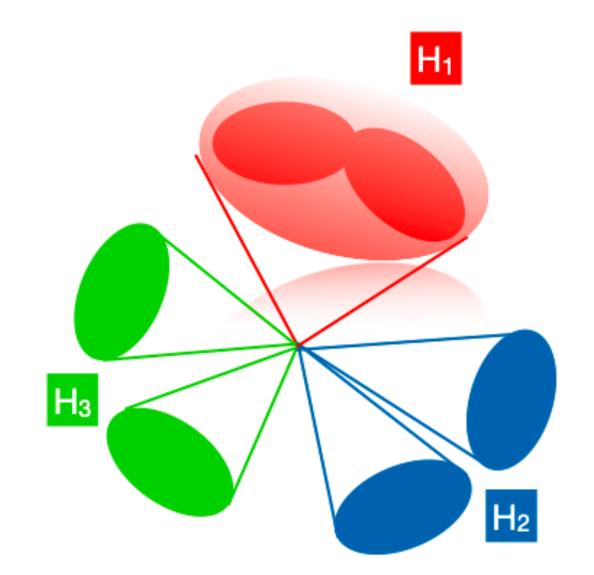


#### Quartic Higgs couplings HHH-6b

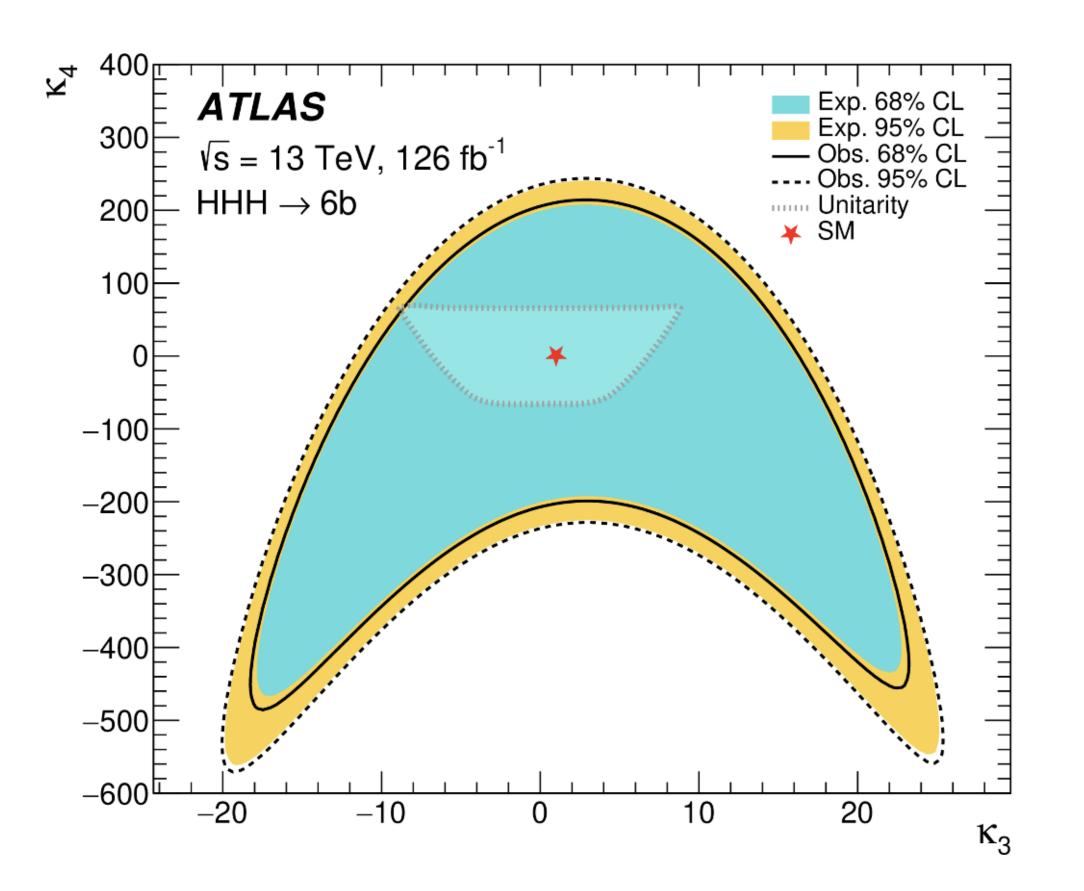
arXiv:2411.02040



First time that the constraints are set on the quartic Higgs coupling.



Data consistent with SM. x-sec(pp→HHH) < 59.4 fb @ 95% CL (µ < 747)



# Invisible Higgs

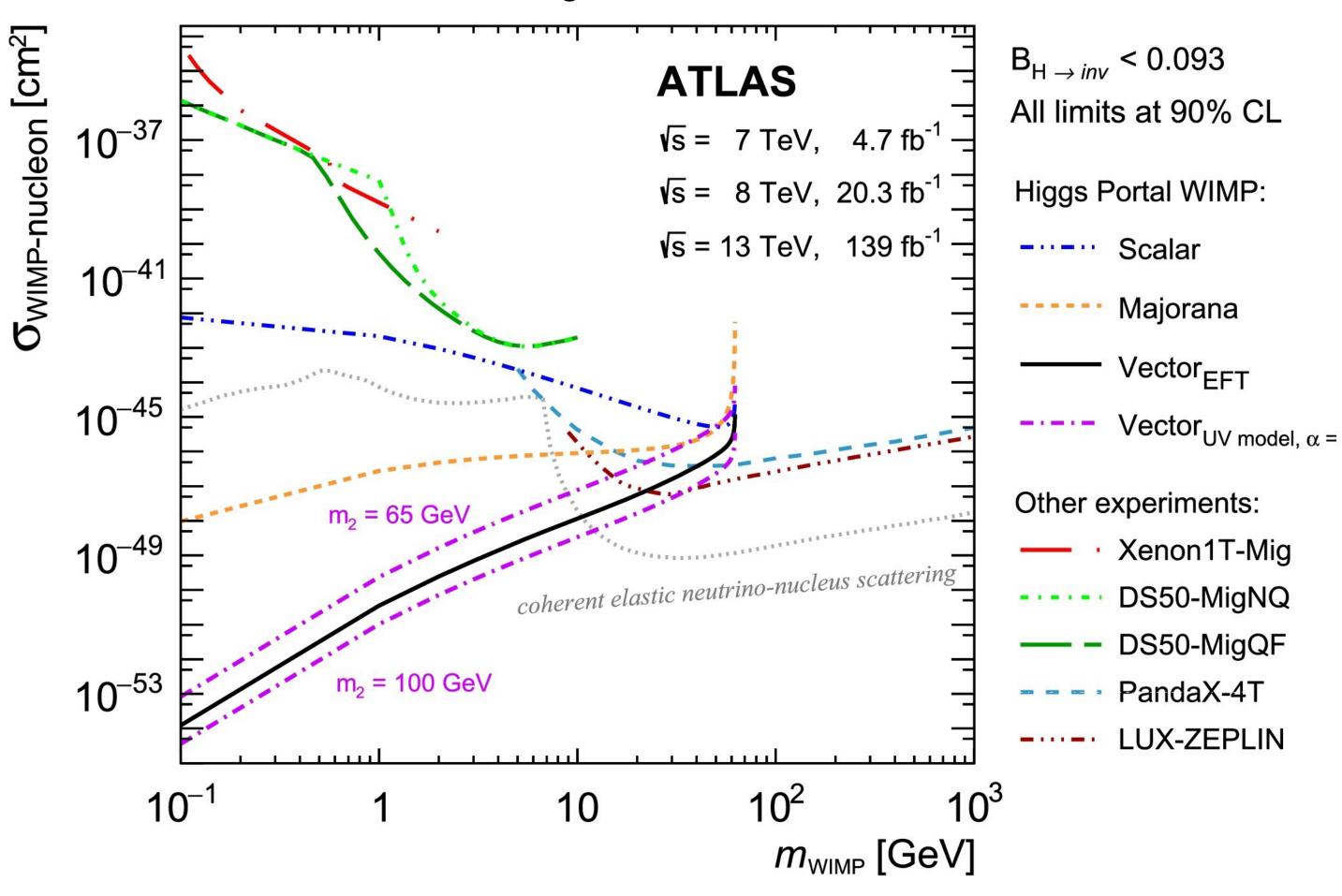


## Invisible Higgs width

- SM particles get mass through the Higgs. Dark matter could behave the same way and be produced in Higgs decays
- SM Higgs invisible decays are <0.1%</li>
- The analysis:  $B(H \rightarrow inv)_{obs} < 10.7\% @95\%CL$   $B(H \rightarrow inv)_{exp} < (7.7\%) @95\%CL$

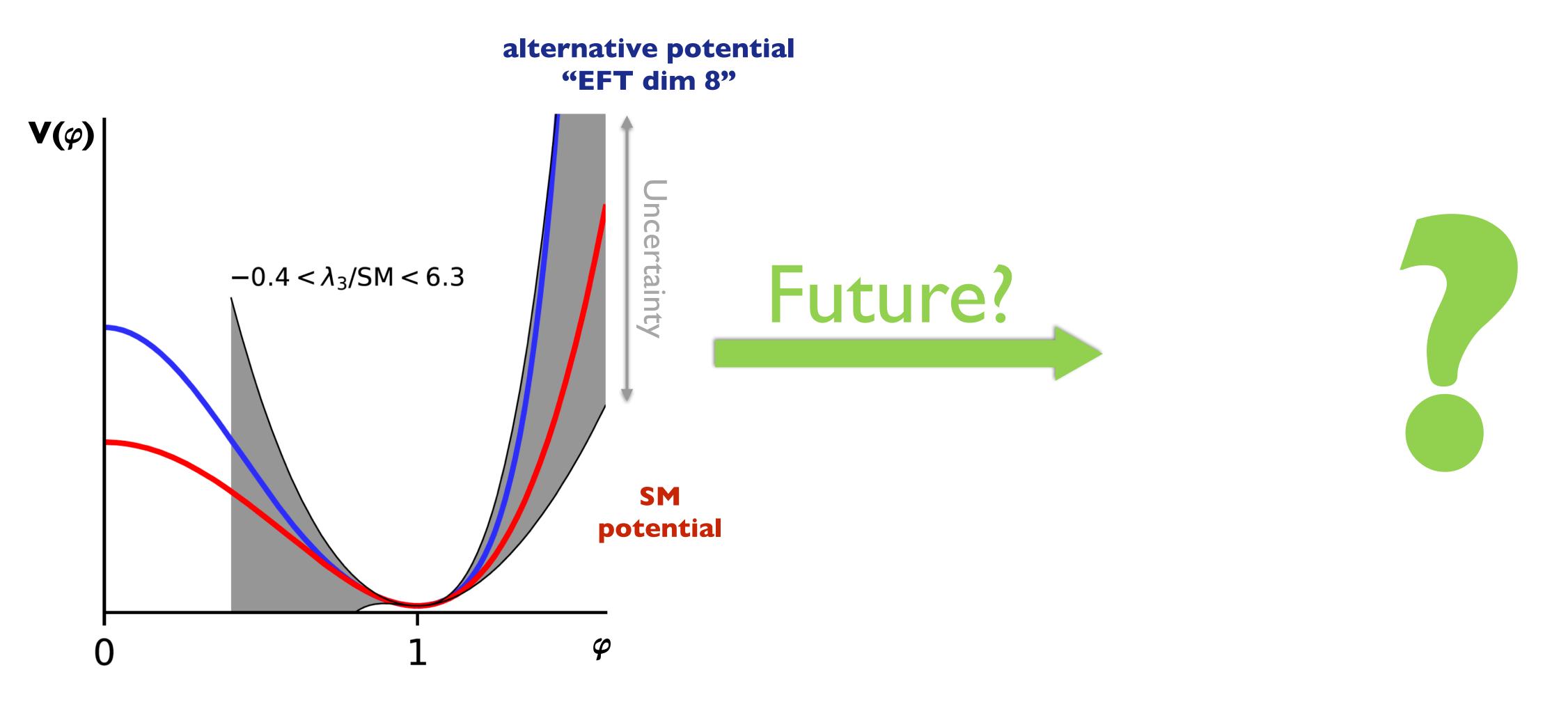
$$ext{BR}(H o ext{inv}) = rac{\Gamma(H o ext{inv})}{\Gamma_{ ext{total}}}$$

90% CL on the spin-independent WIMP-nucleon scattering cross-section

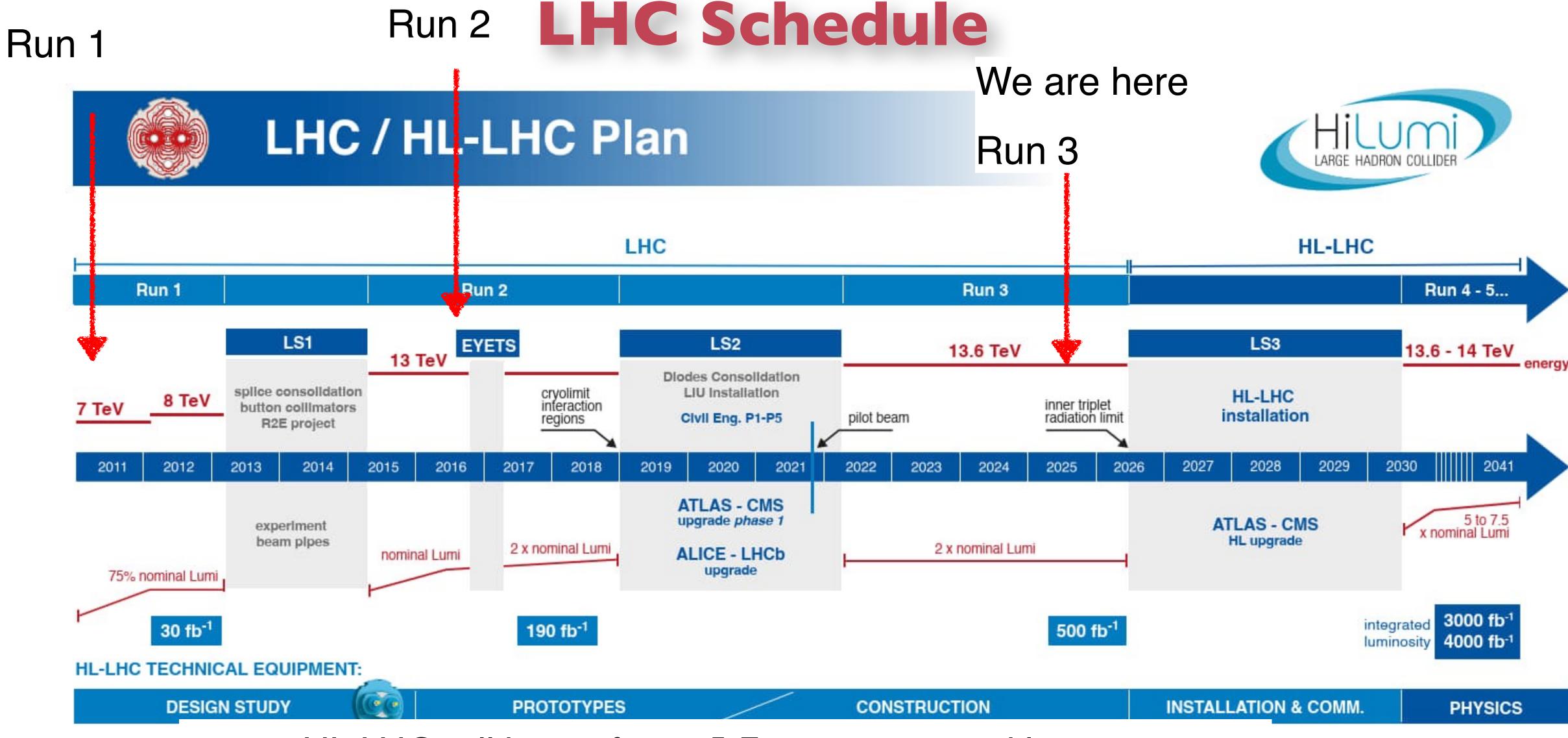


The regions above the limit contours are excluded

# What's the future for the Higgs?







HL-LHC will have a factor 5-7 more integrated luminosity

DEFINITION EXCAVATION BUILDINGS

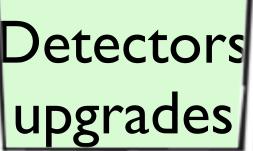
# HL-LHC upgrade: The challenges

### Unprecedented opportunities come with great challenges

HL-LHC promises to provide 3000 fb-1

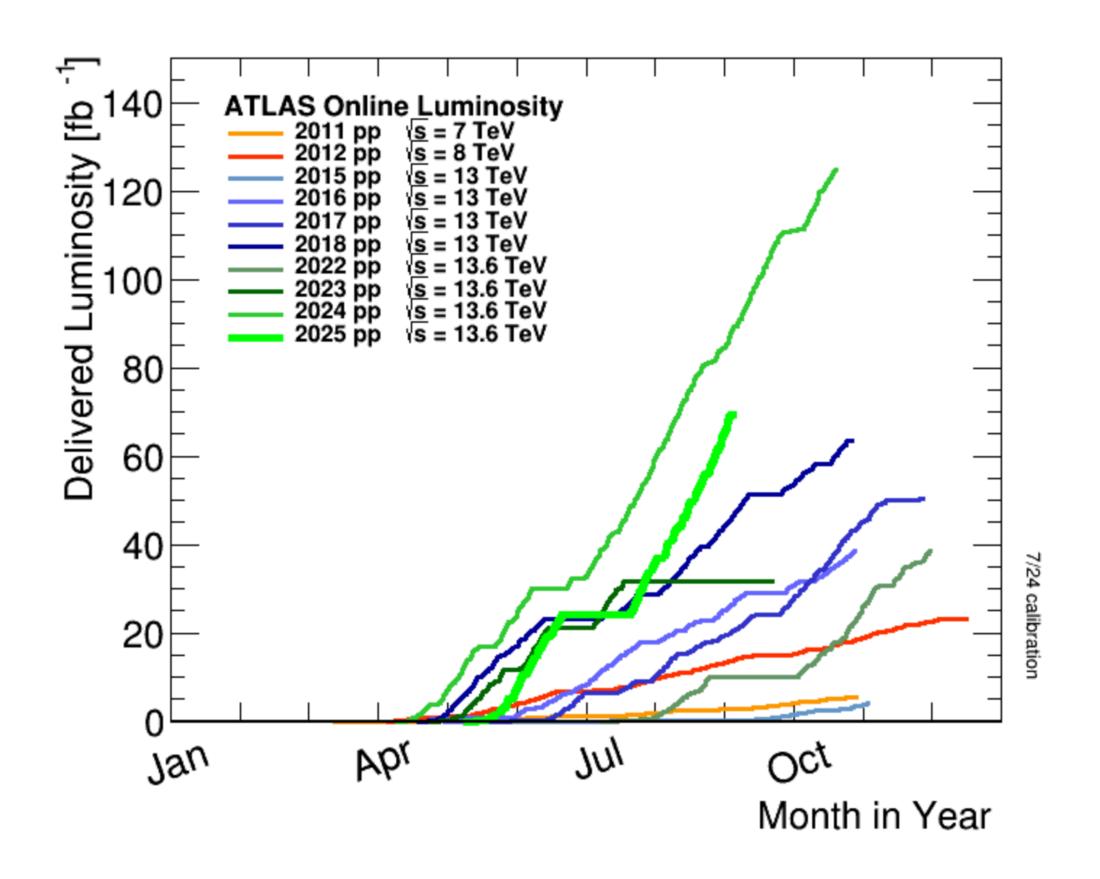
instantaneous luminosity a factor of 5-7 larger than LHC nominal value.

Up to 200 p-p interactions per bunch crossing!

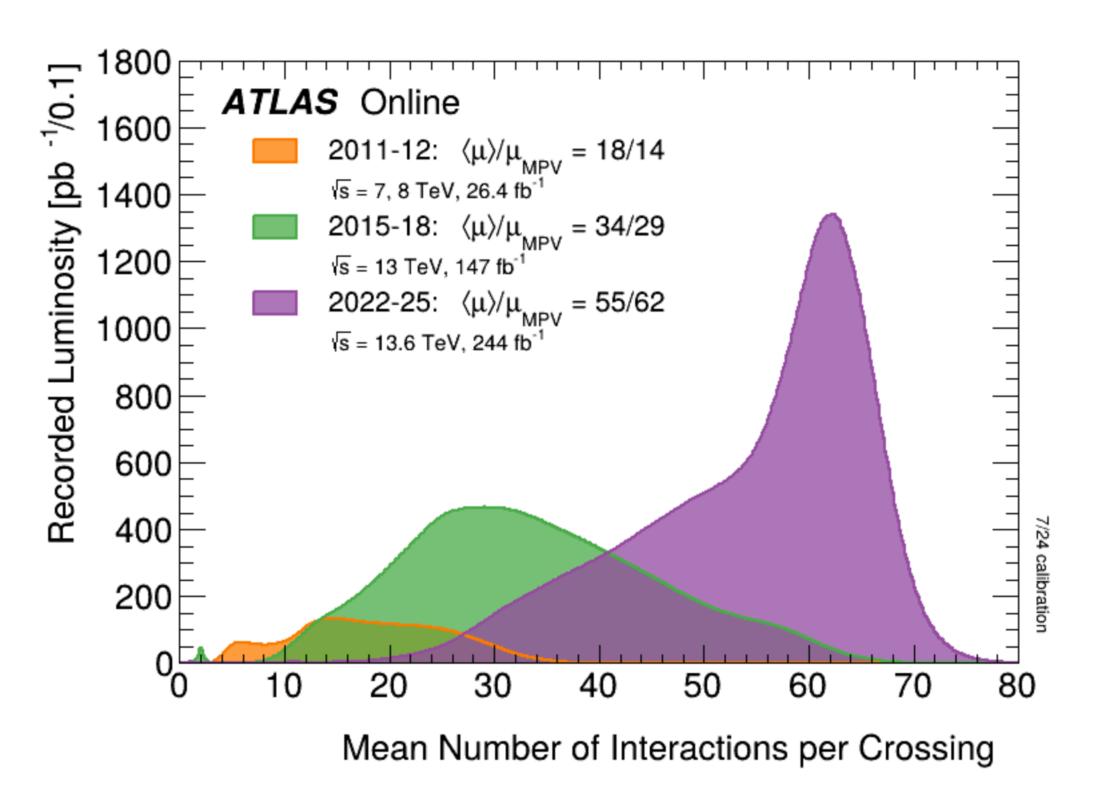


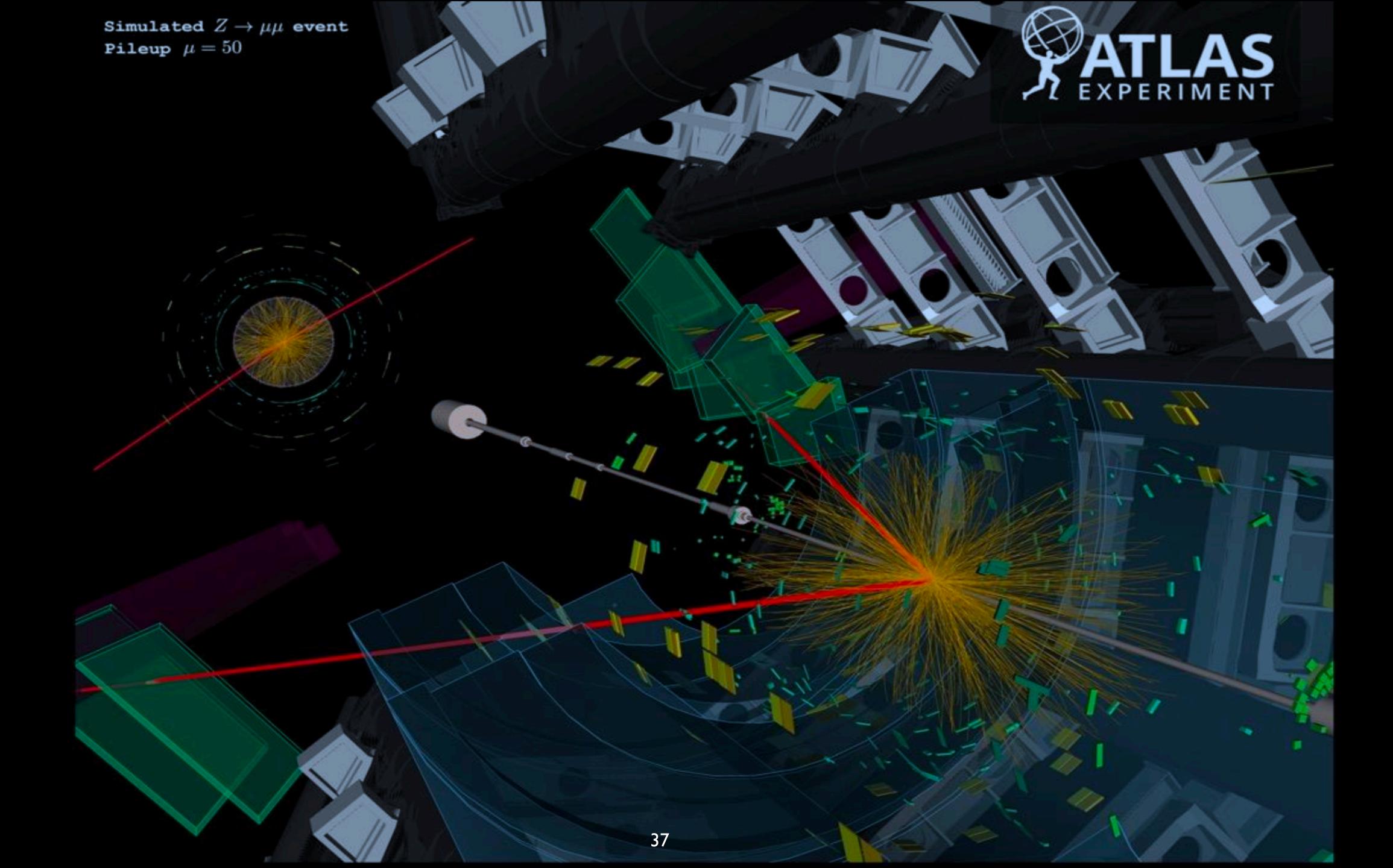
# HL-LHC upgrade: The challenges

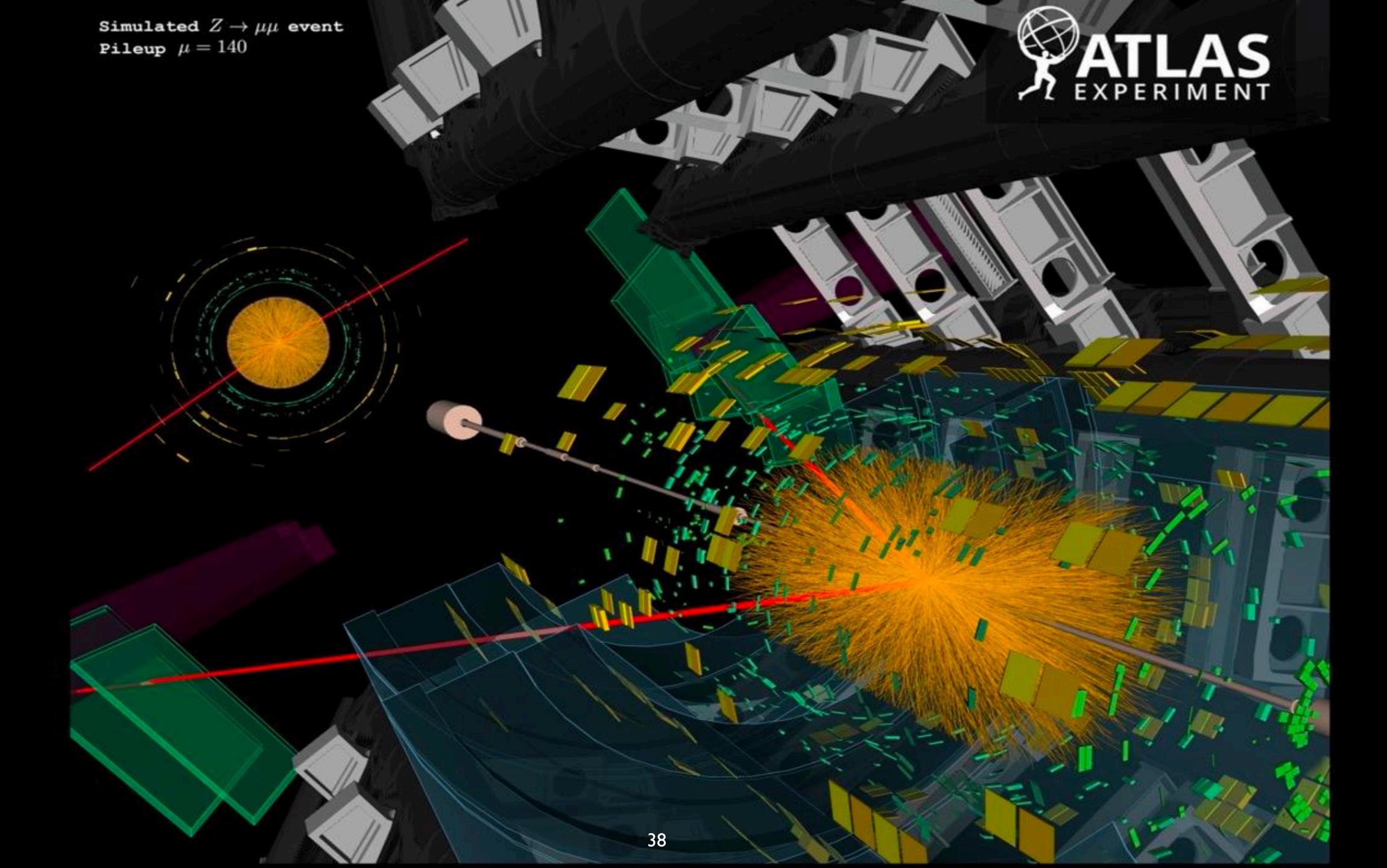
Luminosity-> delivered number of interactions has never had such a steep rise!



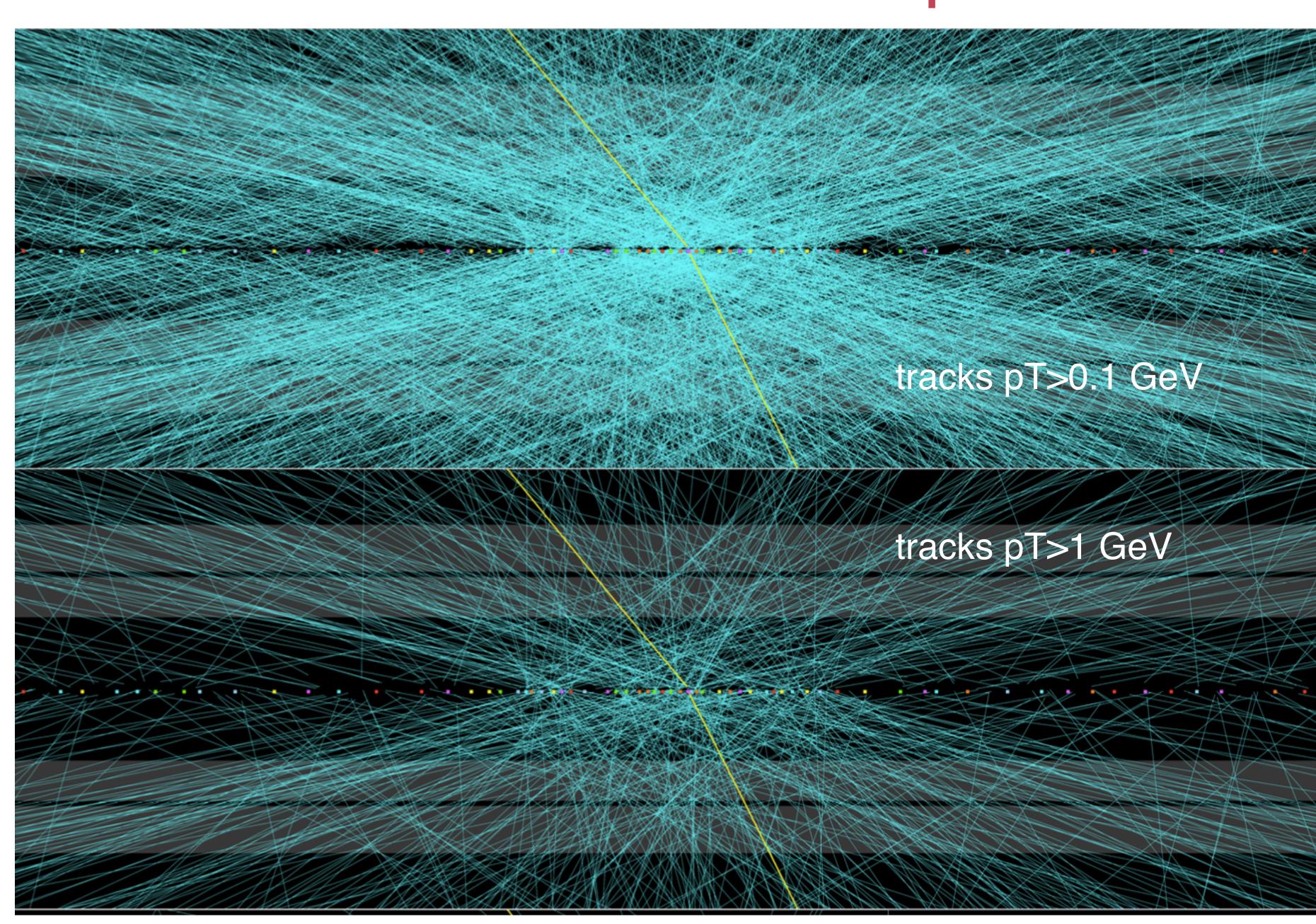
Luminosity->comes at the price of increased number of interactions per bunch crossing (that are a background).  $\mu$ = number of interactions per Crossing







# Pile-up



1)Z->μμ event with 65 interaction vertices

## New Muon Chambers Inner barrel region with new

**RPC and sMDT detectors** 

new and upgraded forward and luminosity detectors



### **Upgraded Trigger and Data Acquisition system**

Level-0 Trigger at 1 MHz, Full-feature global trigger Improved High-LevelTrigger (150 kHz full-scan tracking)

### **Electronics Upgrades**

- LAr Calorimeter
- Tile Calorimeter
- Muon system

### **High Granularity Timing**

Detector (HGTD) Forward region (2.4 <  $|\eta|$  < 4.0) to reduce

Pile-up

ITK: All silicon, up to  $l\eta l = 4$ 

strongly augmented tracking acceptance, 50x present channels → to cope with high occupancy

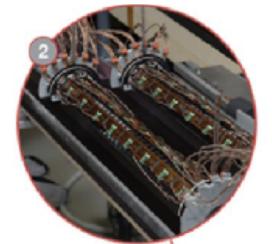
## CMS upgrade

During Long Shutdown 2 (2018-2022), CMS completed the Phase 1 upgrades and started the Phase 2 upgrades. Some highlights:

- Phase 1: HCAL barrel readout, new barrel inner pixel (layer 1)
- Phase 2: First of GEM chambers installed, upgraded CSC electronics for HL-LHC, new beam pipe.
- GPU at HLT and transitioned to a hybrid CPU + GPU in trigger software (HLT nodes) : A Graphics Processing Unit (GPU) is a programmable architecture, offering large number of parallel independent streams of instructions, originally designed for image processing. Accelerate online processing

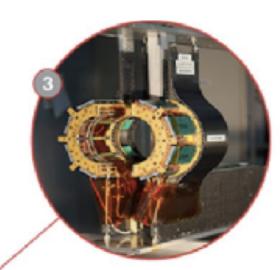
#### **BEAM PIPE**

Replaced with an entirely new one compatible with the future tracker upgrade for HL-LHC, improving the vacuum and reducing activation.



#### PIXEL TRACKER

All-new innermost barrel pixel layer, in addition to maintenance and repair work and other upgrades.



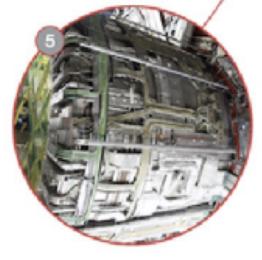
#### BRIL

New generation of detectors for monitoring LHC beam conditions and luminosity.



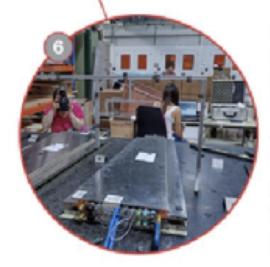
### HADRON CALORIMETER

New on-detector electronics installed to reduce noise and improve energy measurement in the calorimeter.



### SOLENOID MAGNET

New powering system to prevent full power cycles in the event of powering problems, saving valuable time for physics during collisions and extending the magnet lifetime.

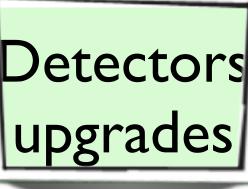


### CATHODE STRIP CHAMBERS (CSC)

Read-out electronics upgraded on all the 180 CSC muon chambers allowing performance to be maintained in HL-LHC conditions.

## GAS ELECTRON MULTIPLIER (GEM) DETECTORS

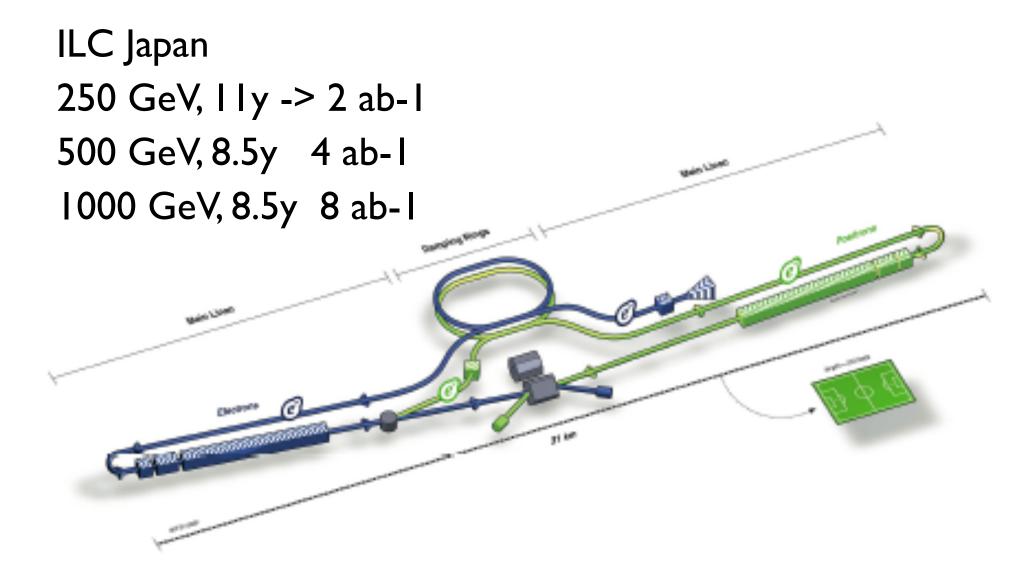
An entire new station of detectors installed in the endcap-muon system to provide precise muon tracking despite higher particle rates of HL-LHC.



## Future Colliders after LHC

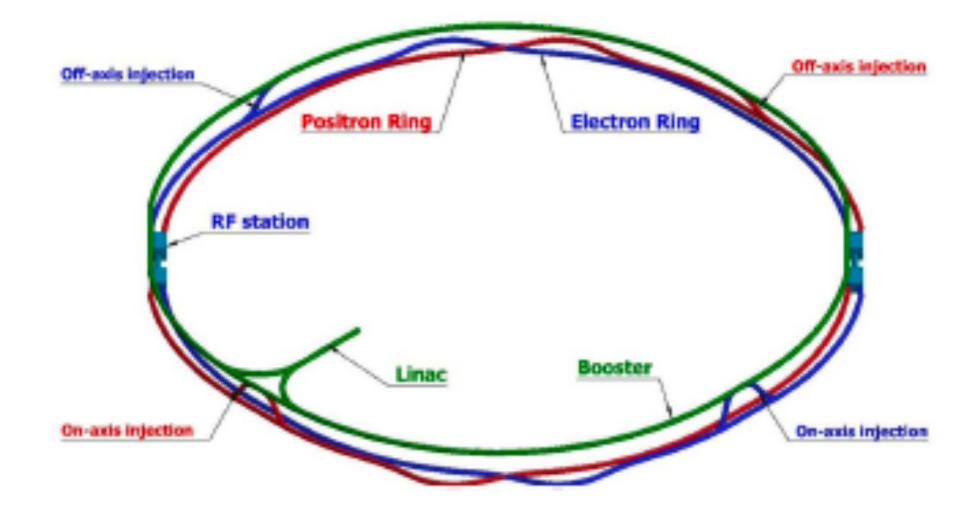
### Future e+e- Colliders in Asia

### Linear e+e-



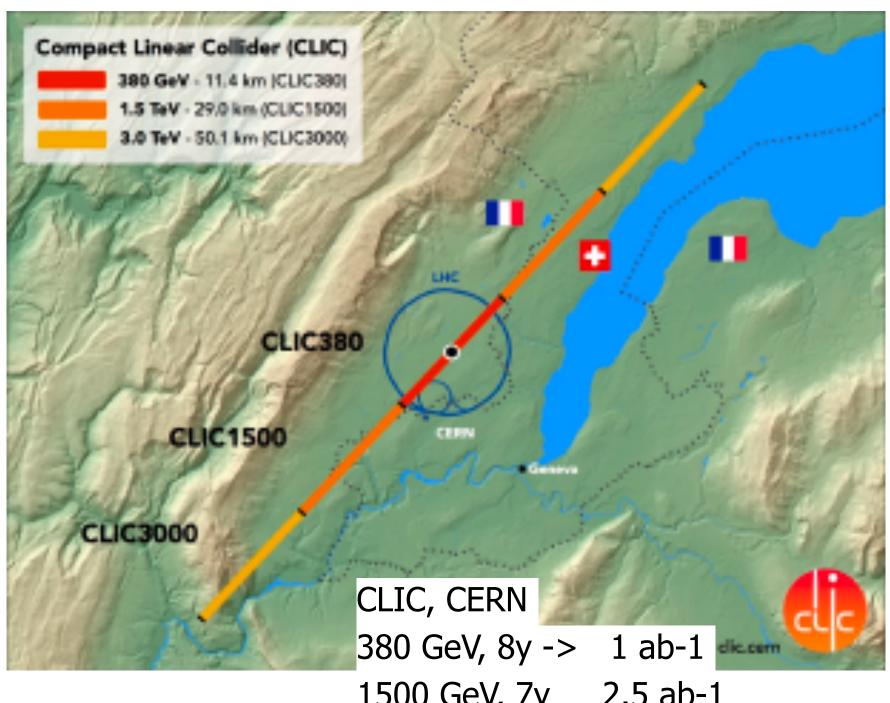
### Circular e+e-

CEPC, China 100 Km tunnel mZ, 2y -> 16 ab-1 2 mW, 1y 2.6 ab-1 240 GeV, 7y 5.6 ab-1



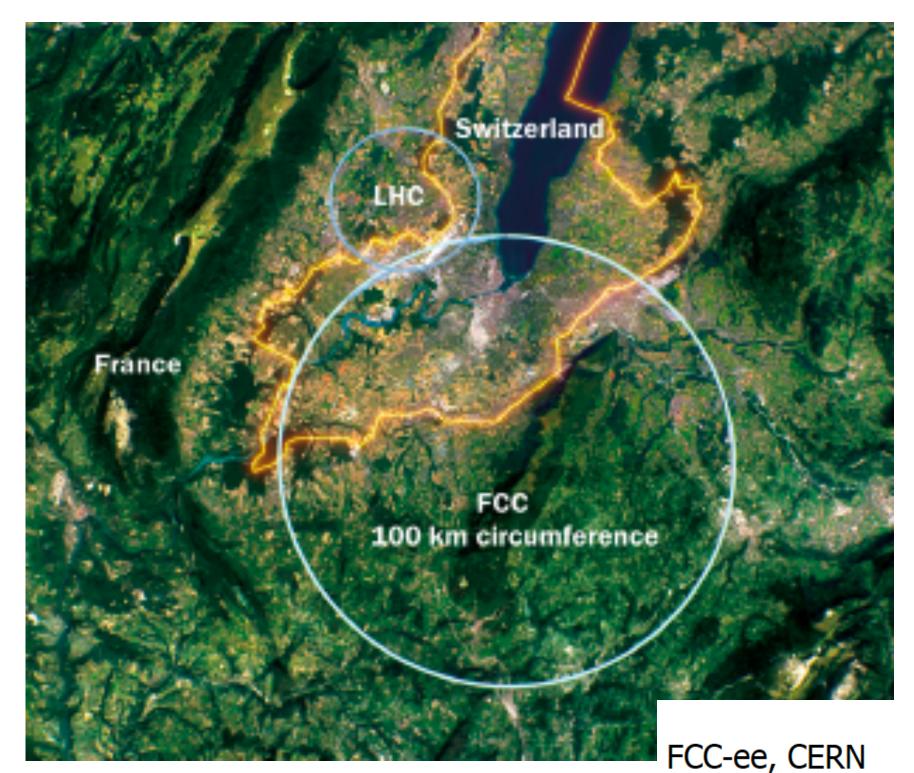
## Future e+e- Colliders in Europe

### Linear e+e-: CLIC



1500 GeV, 7y 2.5 ab-1 3000 GeV, 8.5y 5 ab-1

### Circular e+e-: Fcc-ee



100 Km tunnel in Geneva

m<sub>z</sub>, **4y-> 150 ab**-1 2 m<sub>W</sub>, 1-2y 10 ab<sup>-1</sup> 240 GeV, 3y 5 ab<sup>-1</sup> 2 m<sub>top</sub>, 5y 1.5 ab<sup>-1</sup>

### Can we reuse the LHC tunnel?

LEP3 e+e- machine in the LHC tunnel

**LEP3** (e<sup>+</sup>e<sup>-</sup>, circular, 91 – 230 GeV)



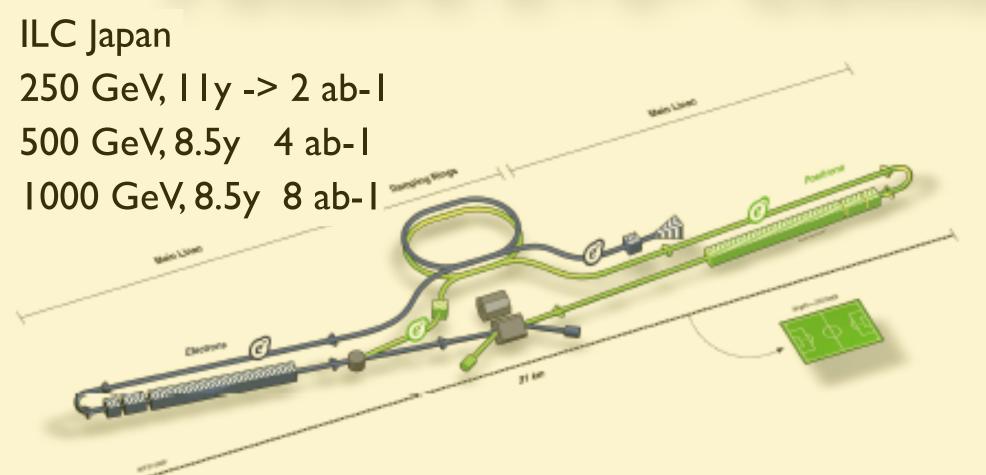
Contra: Lower centre of mass energy wrt to FCC-ee due to synchrotron radiation. Limited number of bunches (tens at most), constrained by RF power and beam-beam effects in the smaller ring.

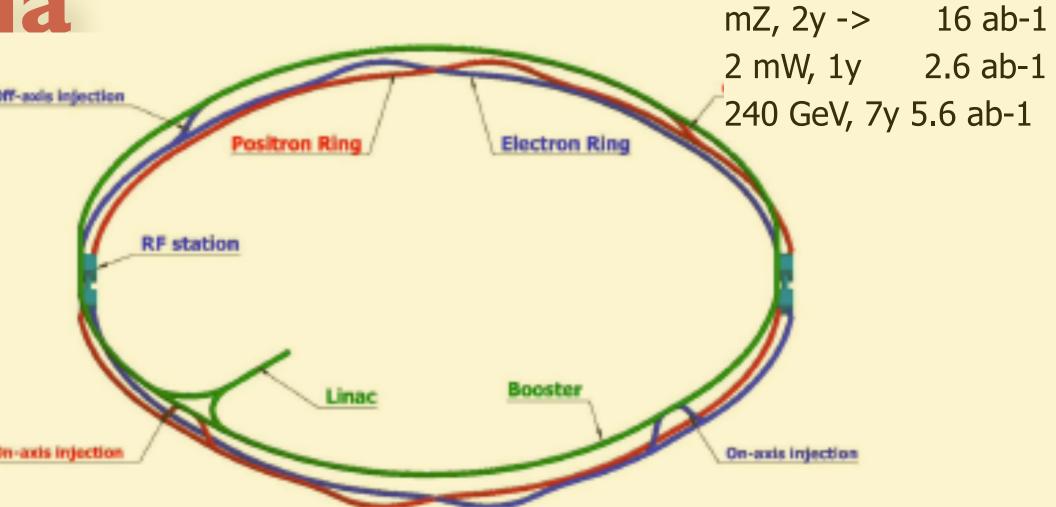
Pros: Lower cost, Build new hadronic (p-p) machine in parallel.

### Do we need an e+e- machine?

e+e- Could measure directly Higgs width and provide precision measurements as input to pp program that would increase the pp measurement precision

Future e+e- Colliders: Asia

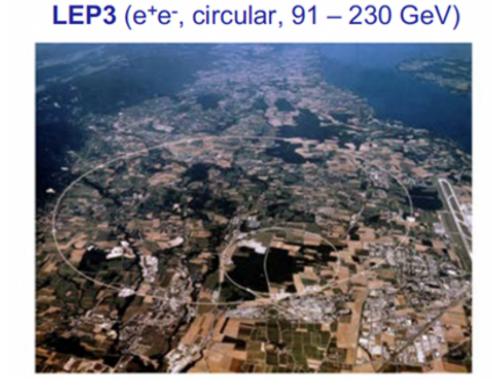


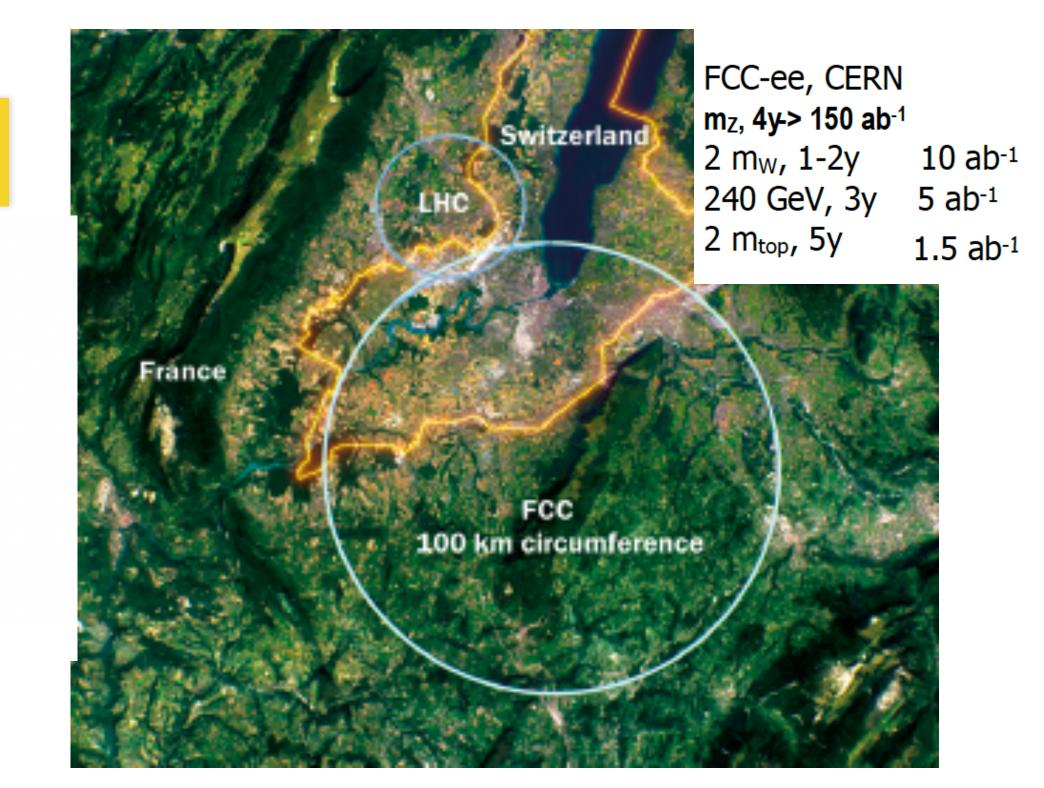


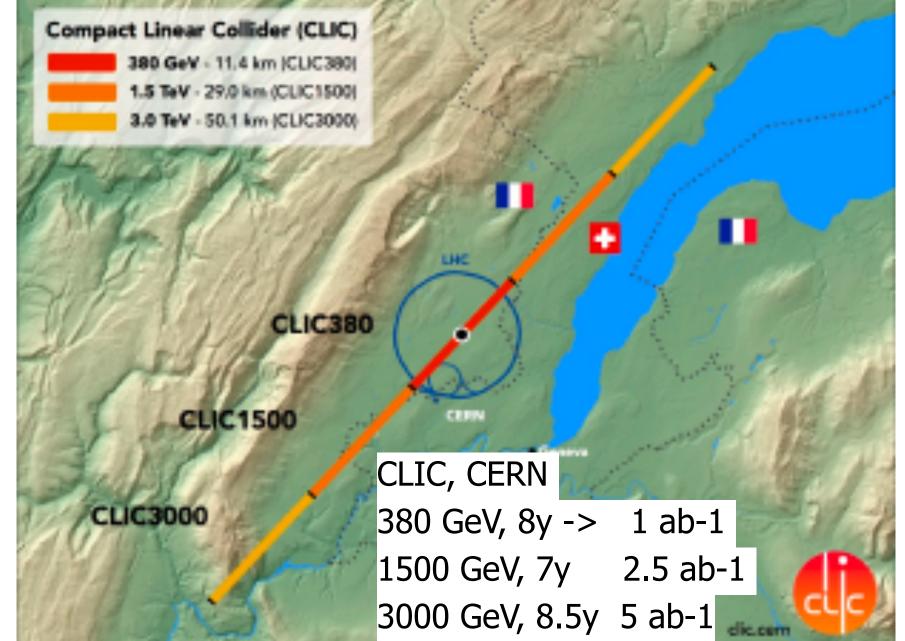
CEPC, China

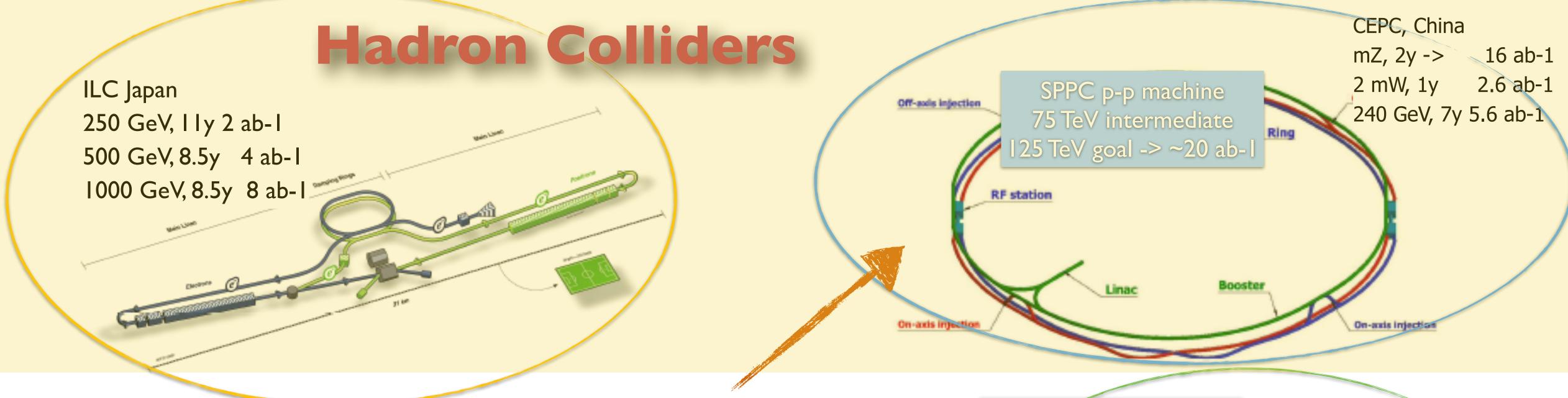
# The soul by 10,000 to 10,0

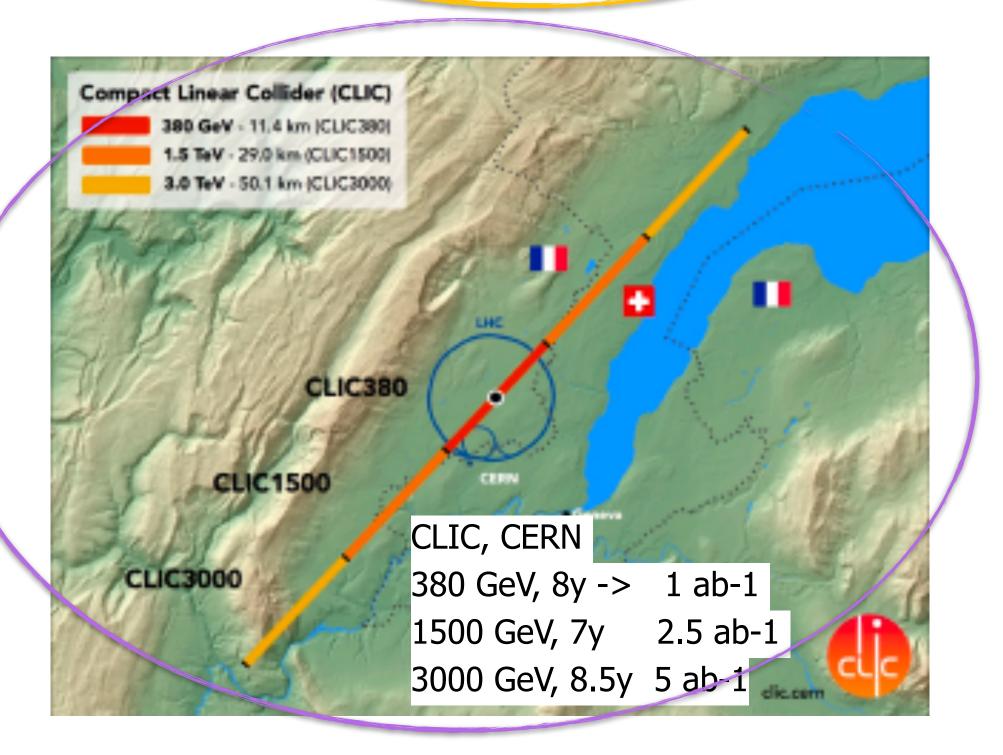
LED2 (ata- aircular 01 220 Ca)/)



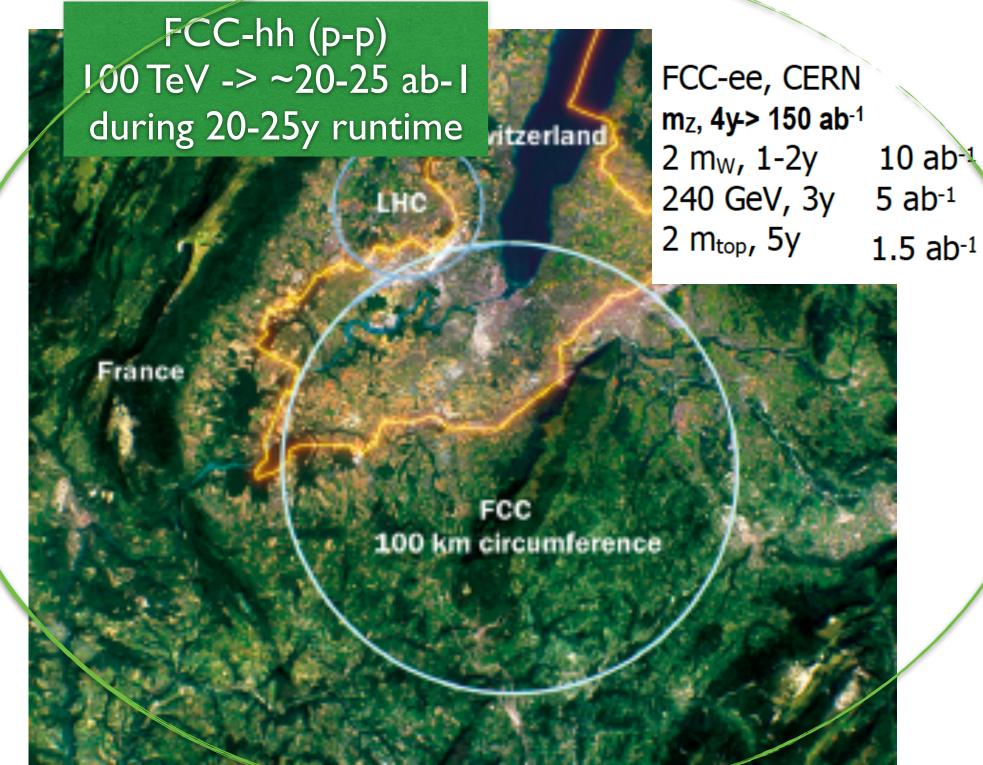






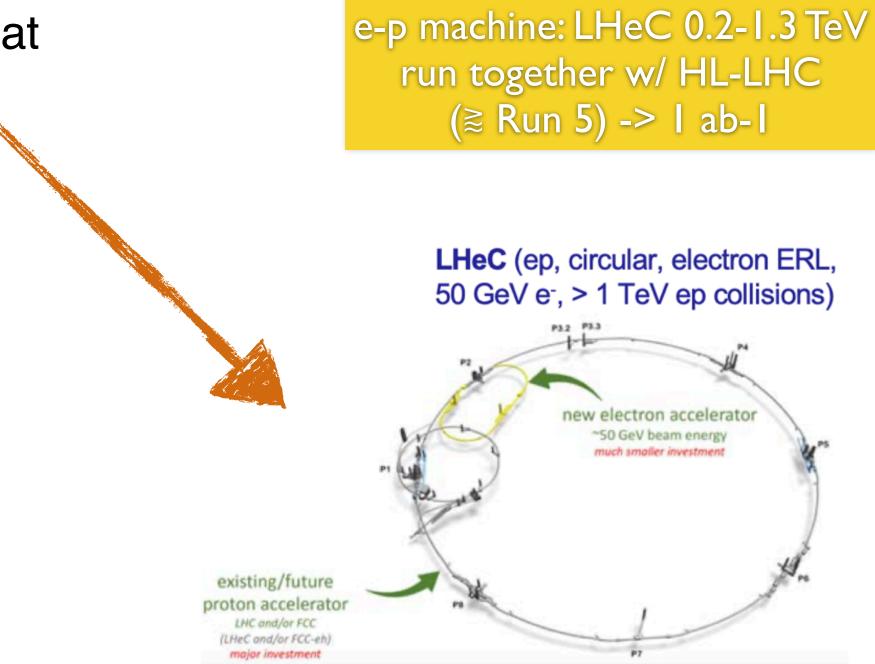


p-p machines in the 100 km tunnels

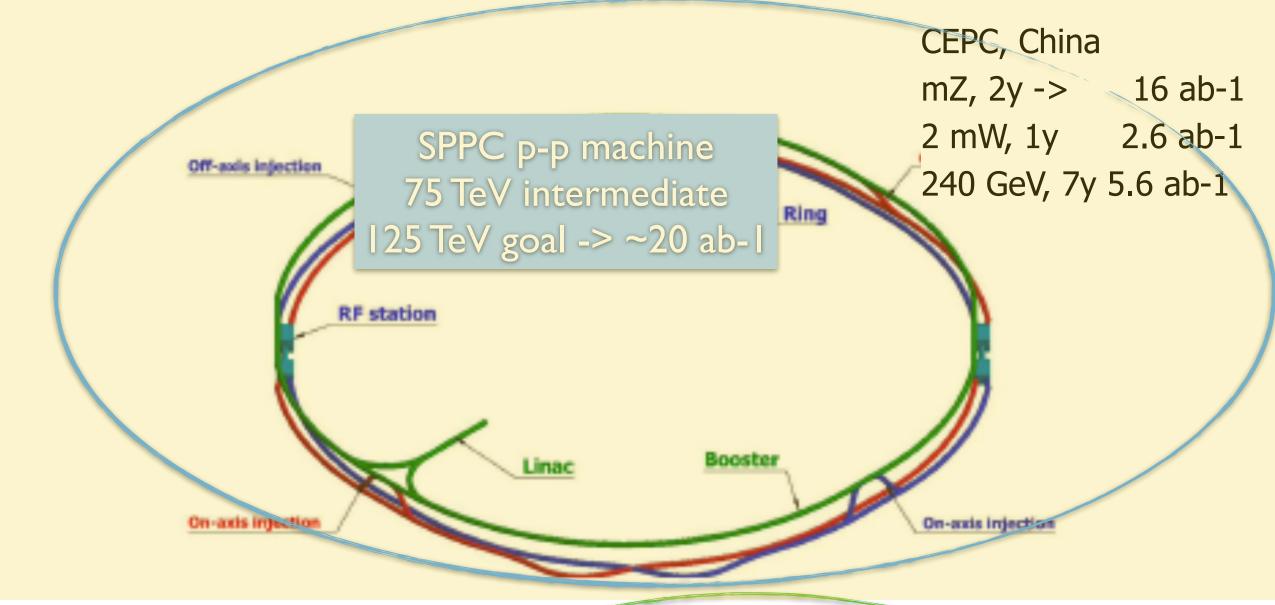


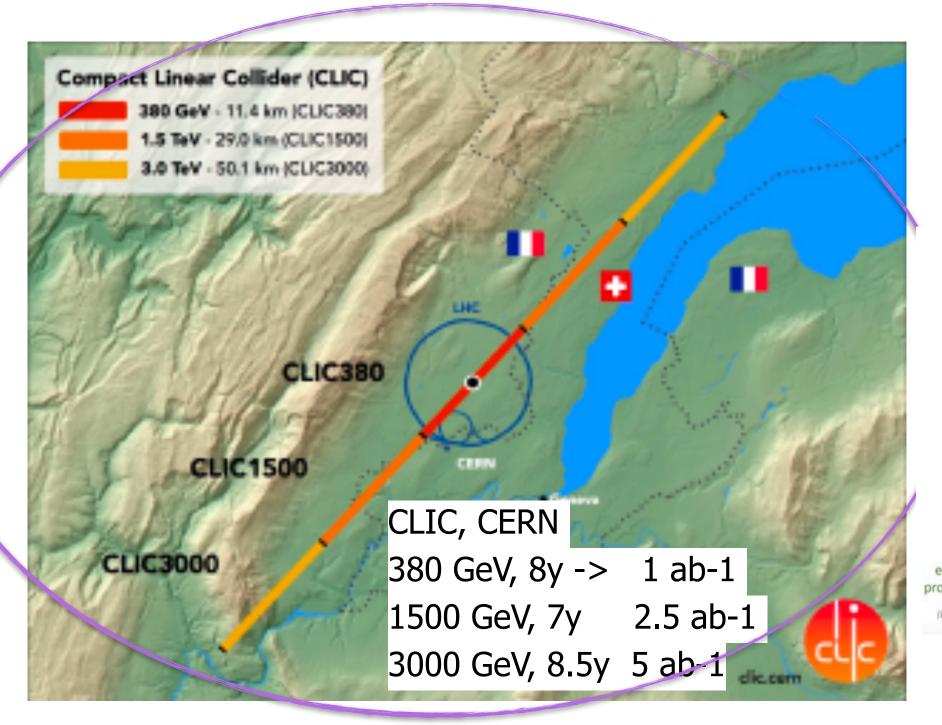
## Do we need an e-p machine?

an electron proton accelerator could also be very useful as it would provide measuerements of the PDFs of the proton, that are going to make errors smaller at a proton proton collider. There is an option called LHe-c, that could run in parallel to LHC, using its tunnel





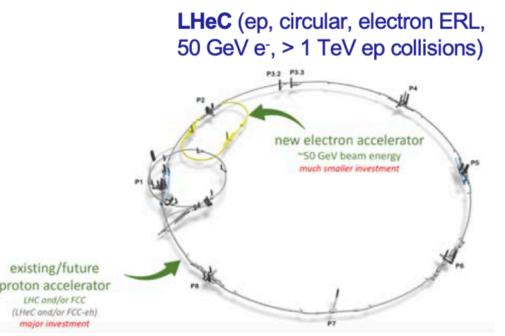


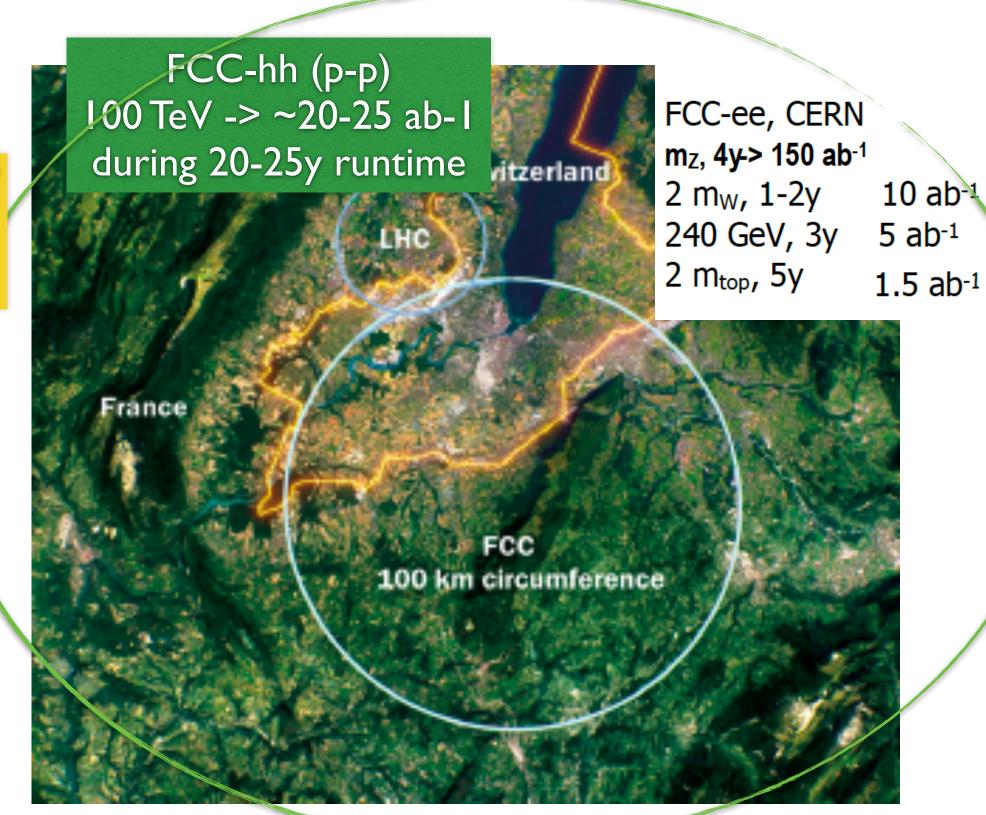


e-p machine: LHeC 0.2-I.3 TeV run together w/ HL-LHC (

(

Run 5) -> I ab-I



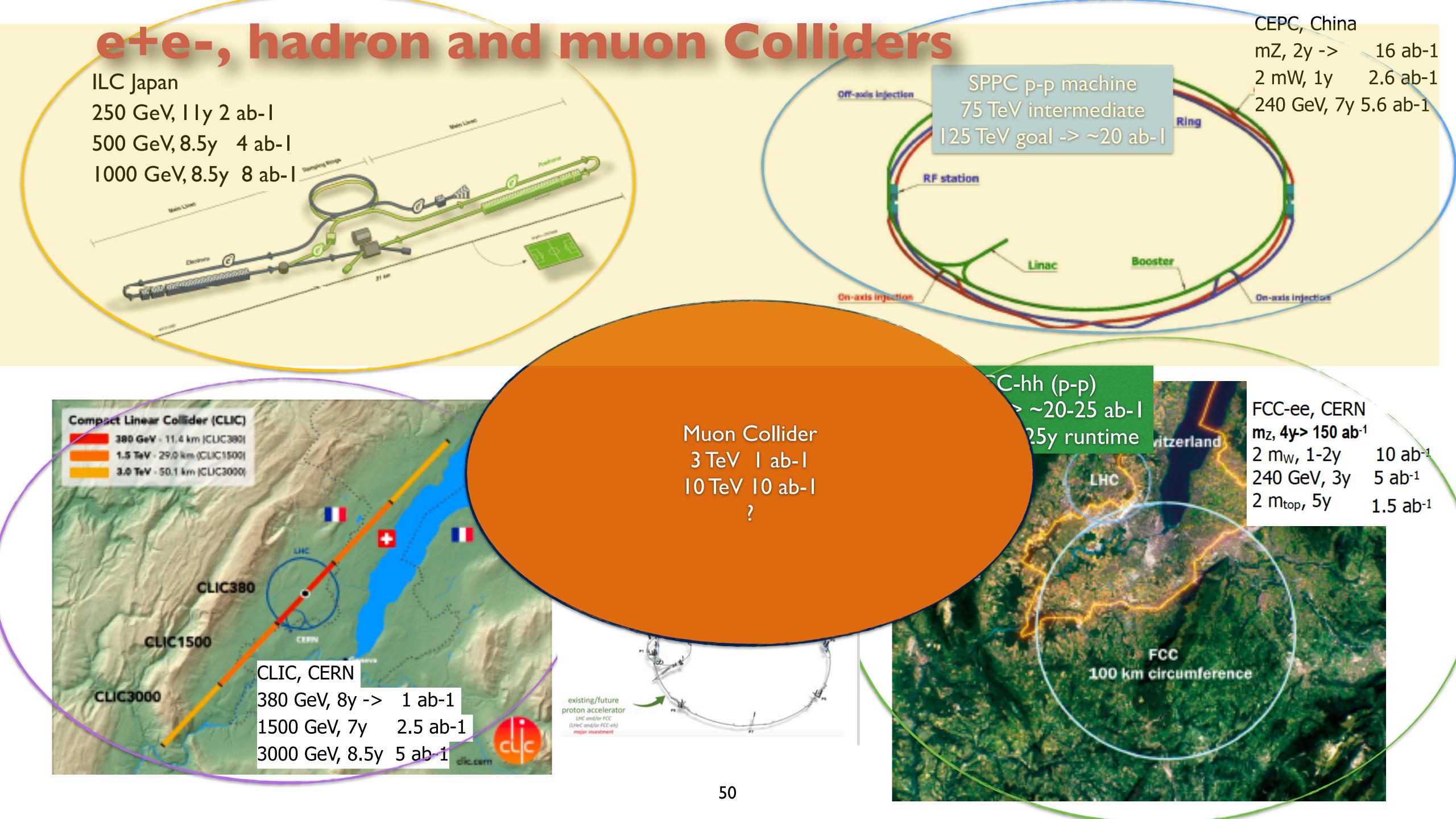


## Do we need an µ+µ- machine?

A muon Collider, is a new adventure. It combines the advantages of a hadron machine and a lepton machine. The muon mass is 207 times the electron mass->Negligible synchrotron radiation emission It does't have PDFs (substructure) and loss of energy as the protons. It can probe precision physics at much lower center of mass energies IP 1 500 Accelerator ring  $\mu$  injector **∑** 200  $\beta = 10$ 50  $\mu$  cooling 10 25 30

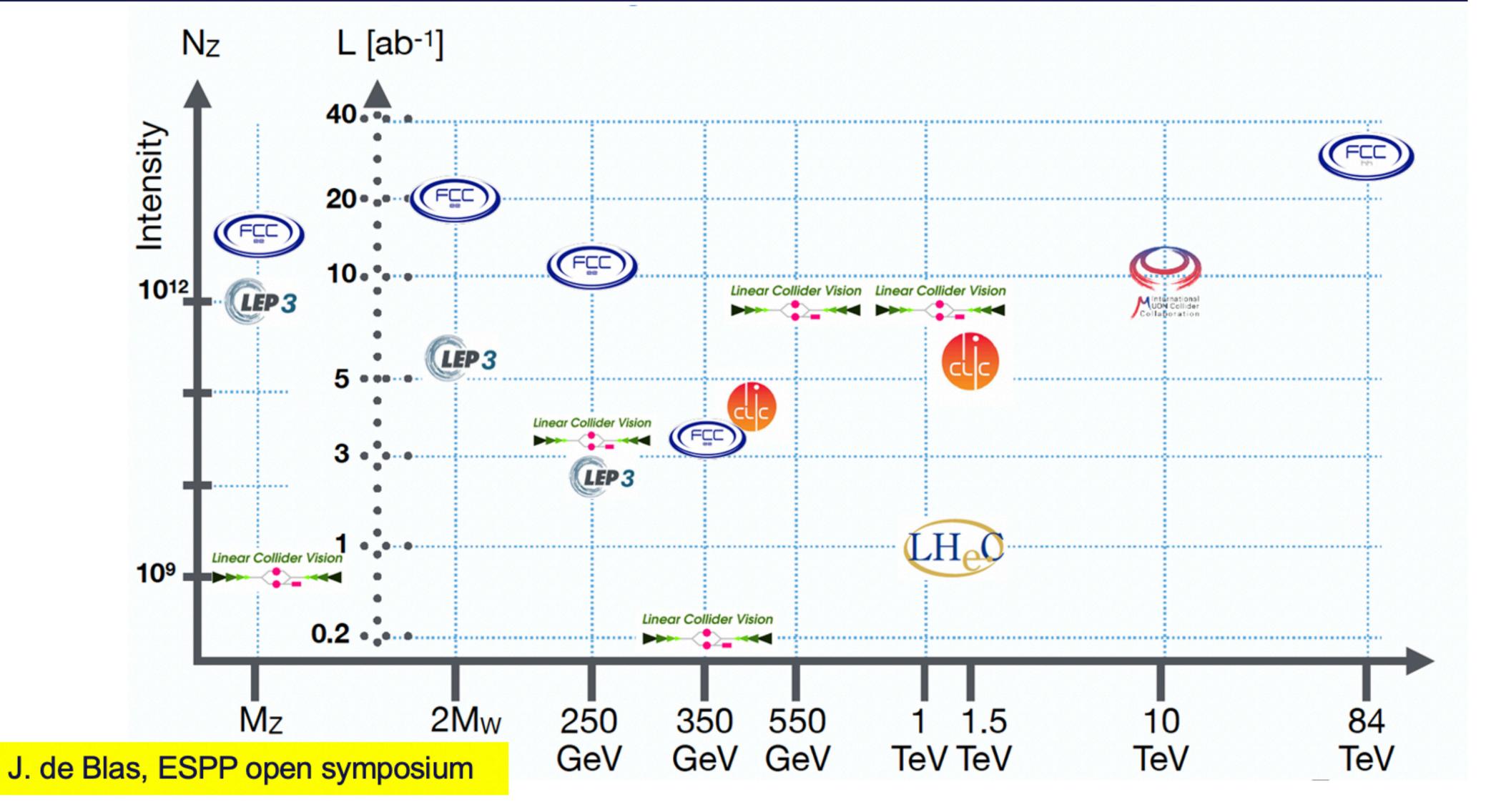
Challenge: muons have limited lifetime: 2.2  $\mu$ s Technology challenge!

**Proton-driven scheme** 

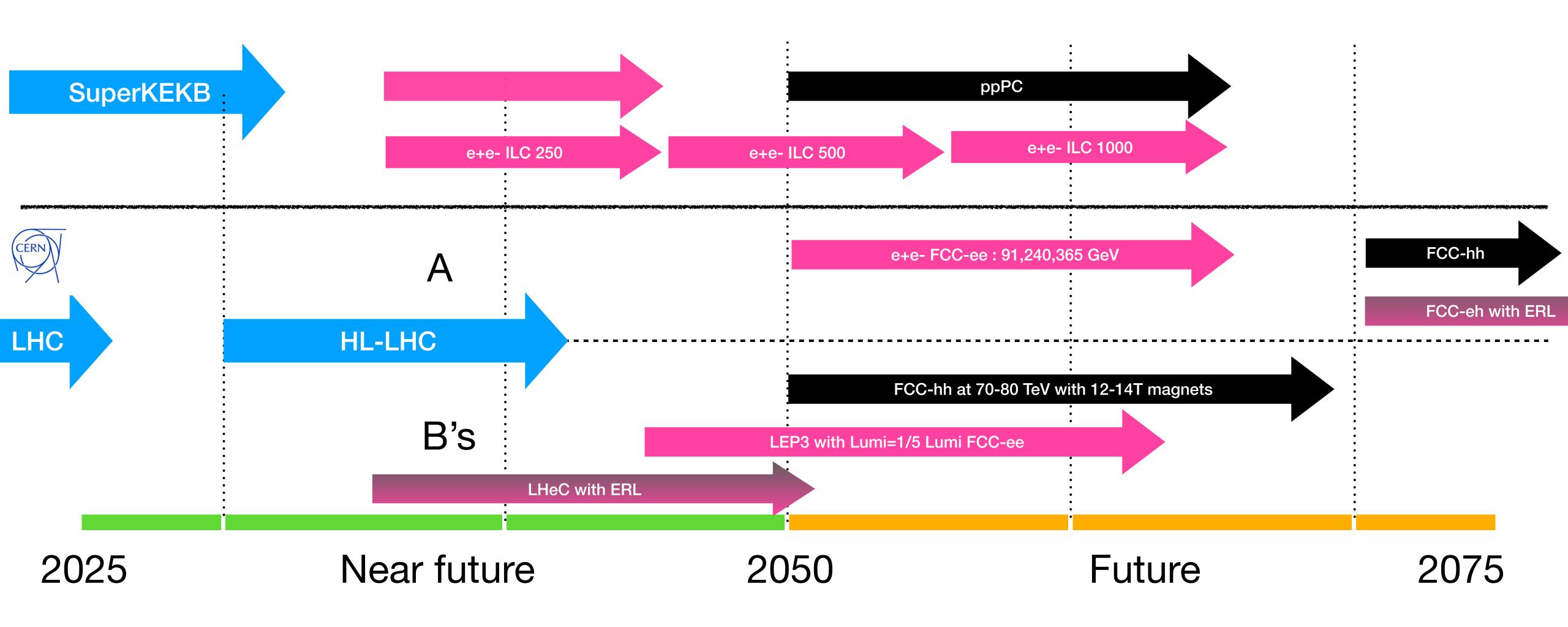


# Lecture 3: measurements in the future

## comparing future collider capabilities

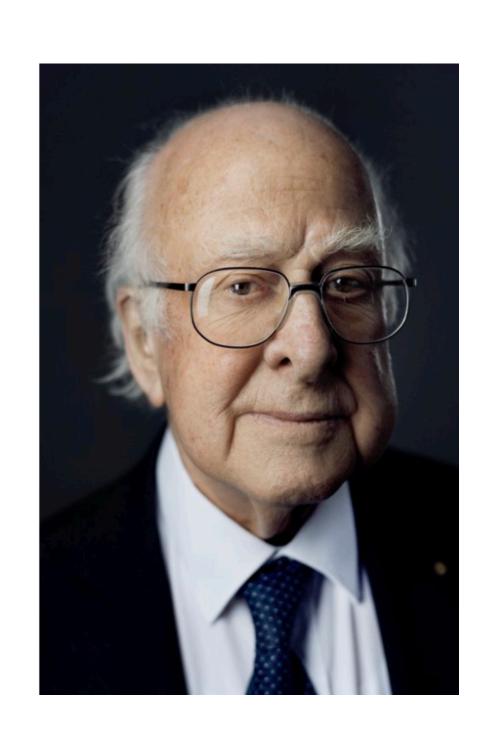


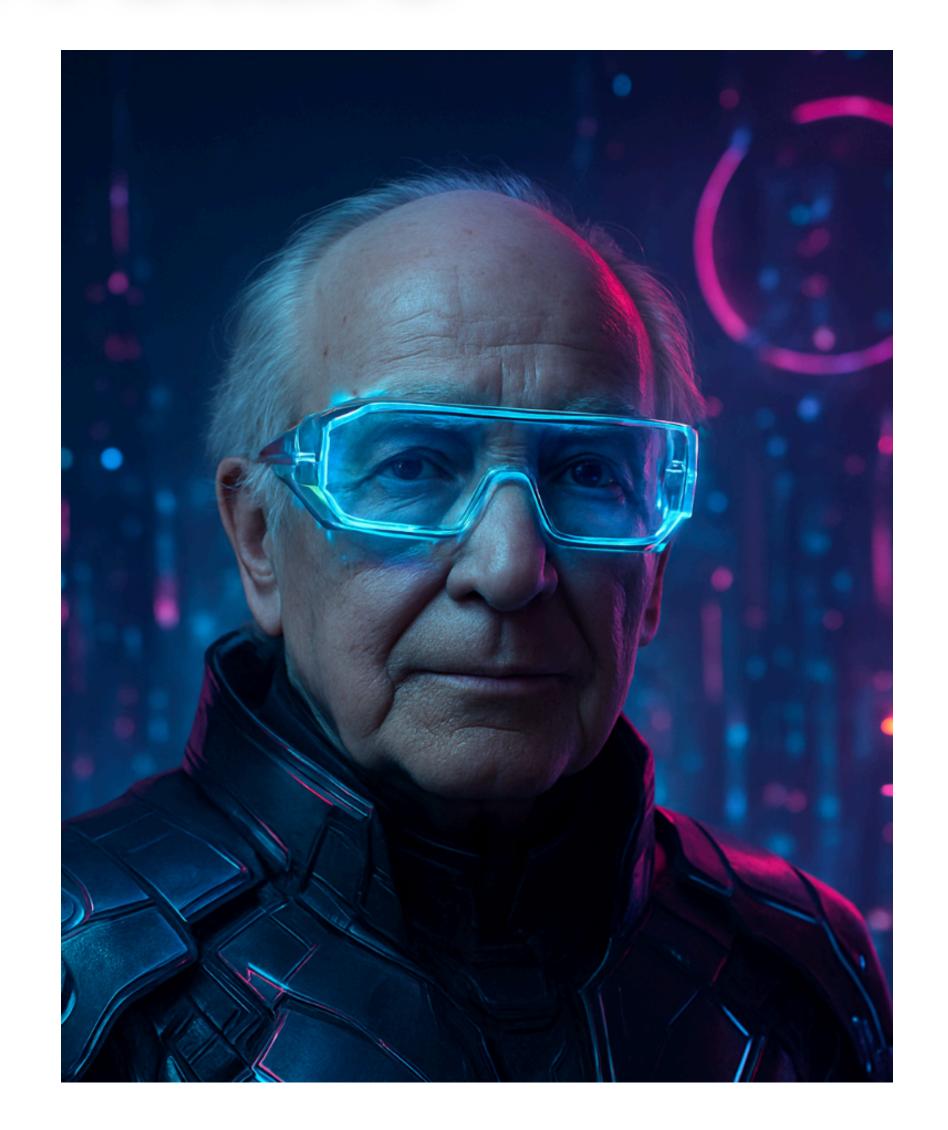
## Future experiments



@F.Maltoni

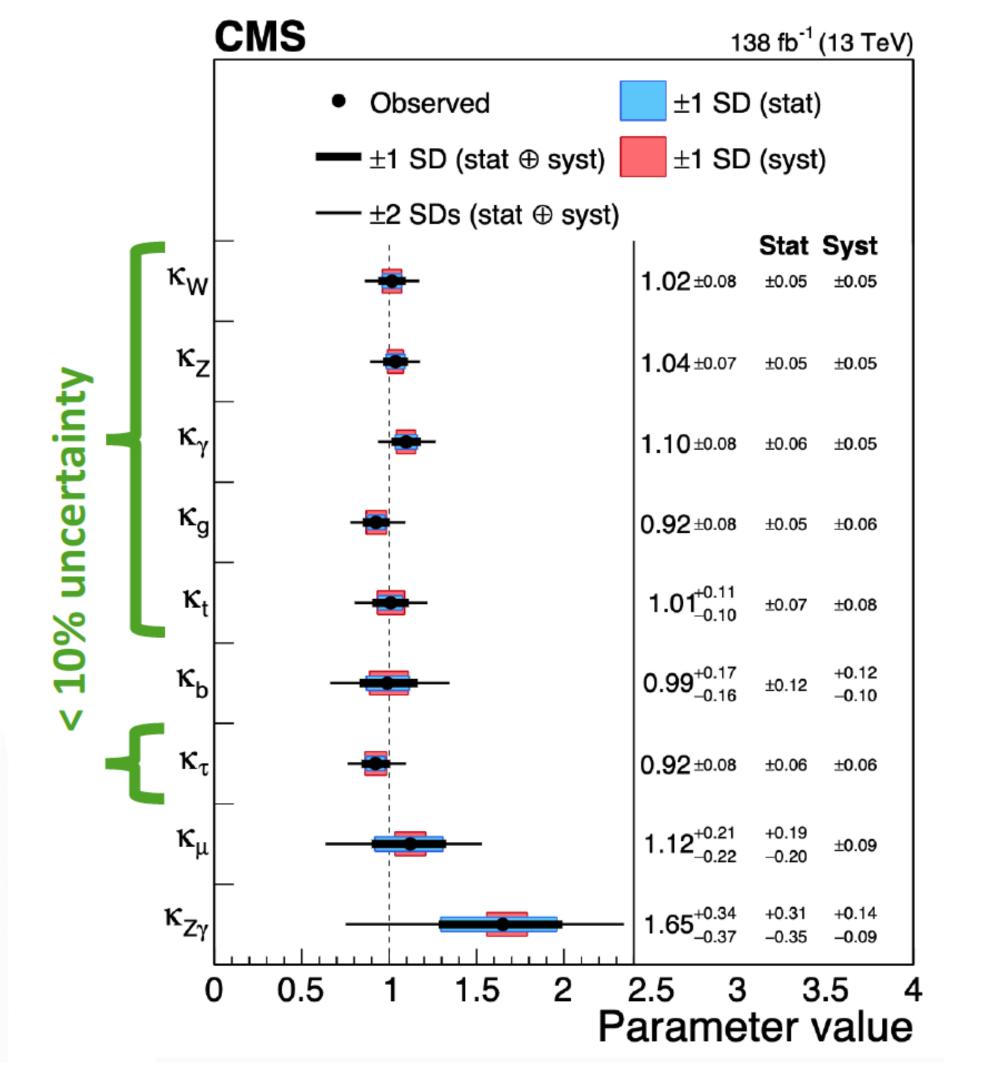
# Single Higgs Future





# Higgs Couplings HL-LHC

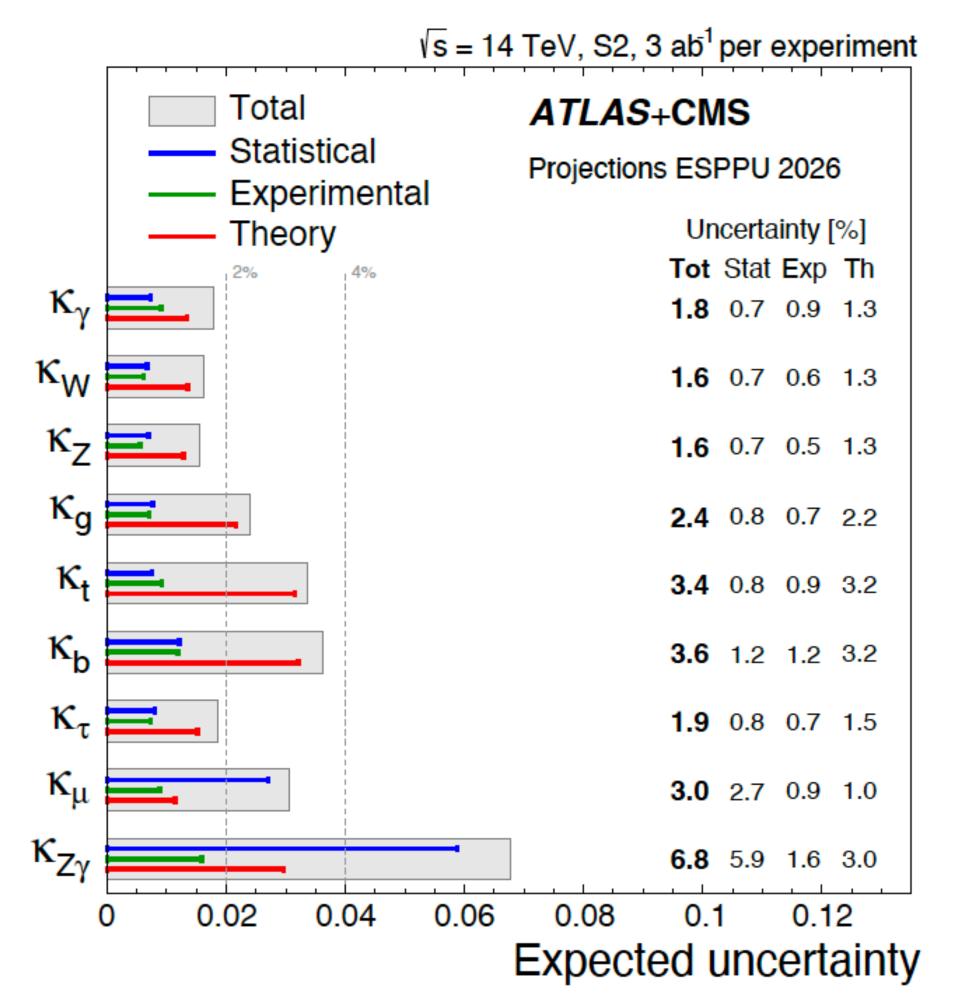




from ~10% to 2-3% precision



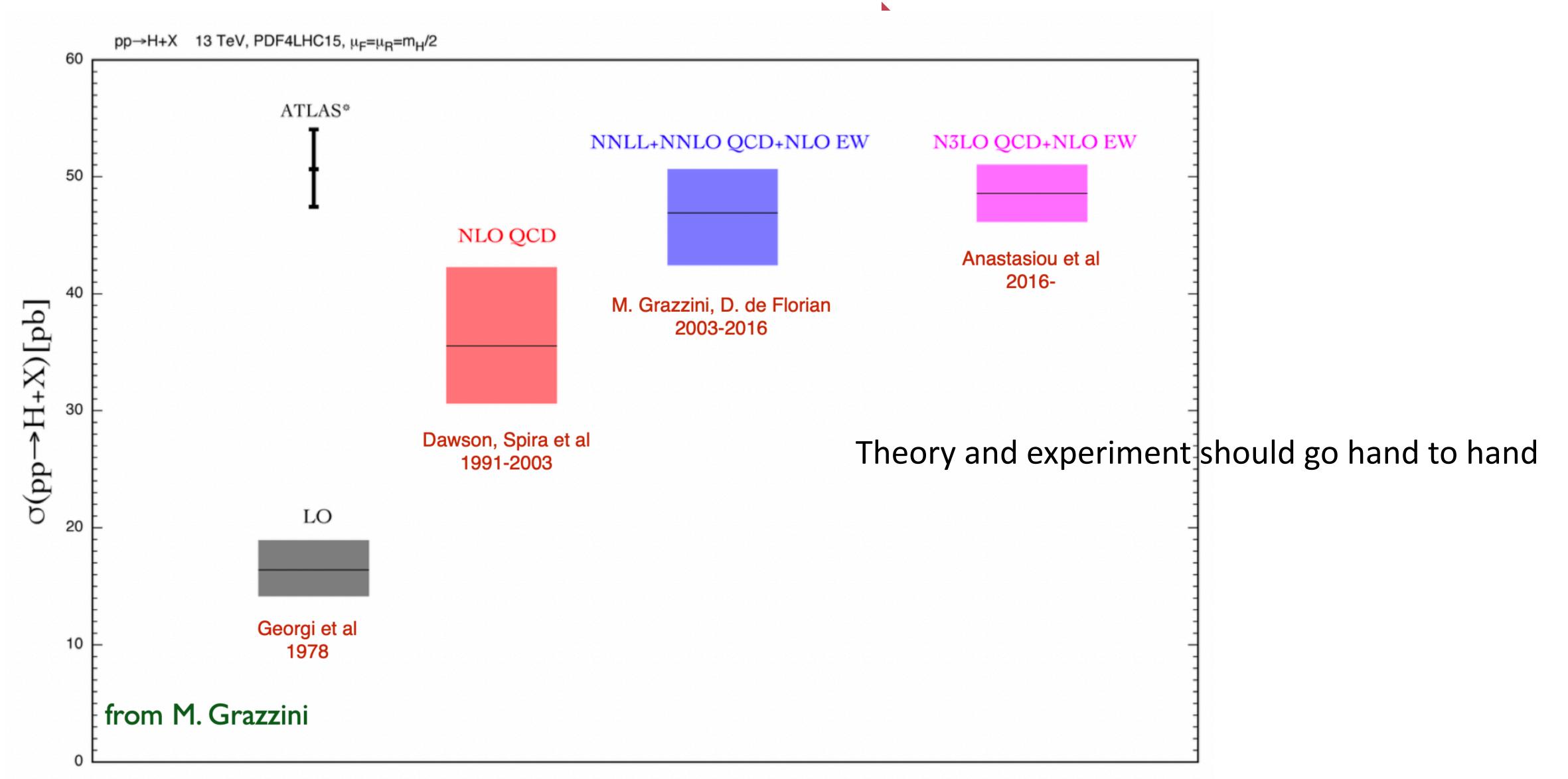




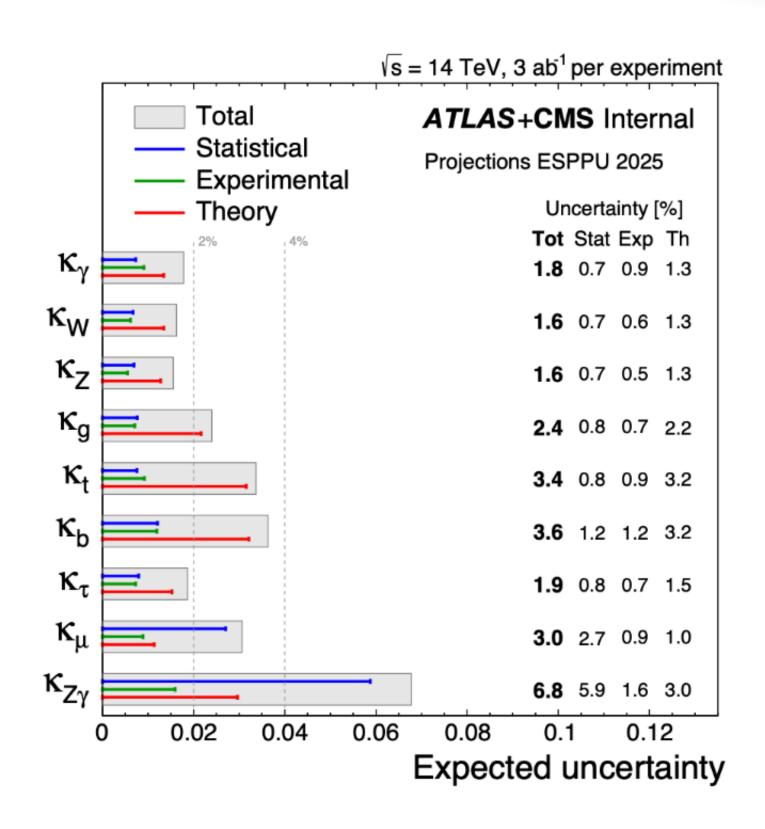
dominated by theory errors!

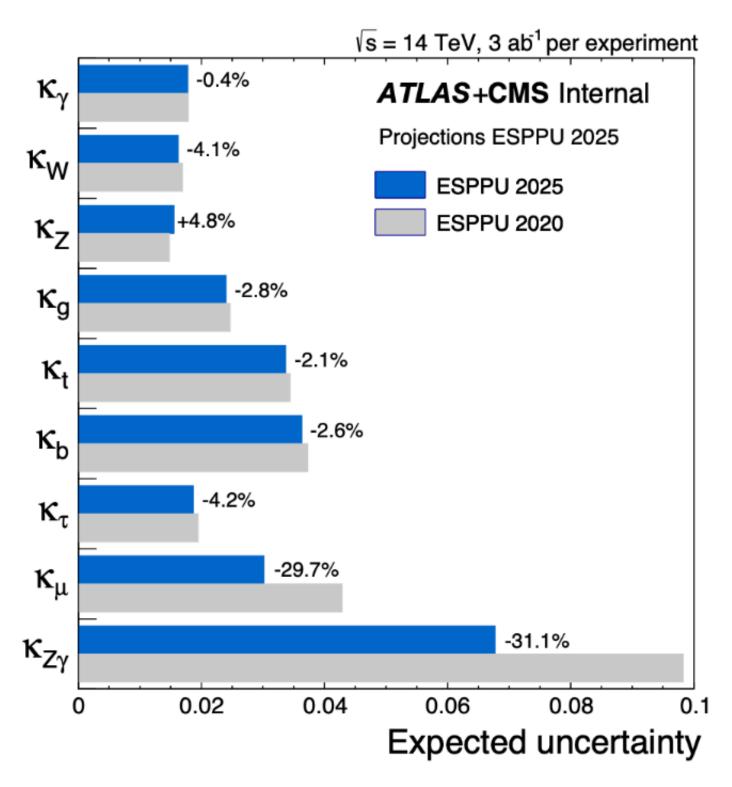


## Inclusive Higgs x-section theoretical improvements



## New Higgs Couplings Projection





### ATLAS Combination,

using most recent Run-2 analyses for

 $H\rightarrow \gamma\gamma$ ,  $Z\gamma$ ,  $\mu\mu$ ,  $VH\rightarrow bb$ ,  $\tau\tau$ 

+ ESU 2020 projections for

H→ZZ, WW ,ttH→bb ,ttH→multileptons

CMS Combination from ESU2020 except  $k_{\mu}$  and  $k_{Z\gamma}$  recent Full Run2 results

k<sub>z</sub> worse than ESU2020 due to refinement of ZH theory uncertainty

# HL-LHC results will stay for years to come and some will be reference until Fcc-hh

[De Blas et al., 2020]

kappa-0	HL-LHC	LHeC	HE-	LHC		ILC			CLIC		CEPC	FC	C-ee	FCC-ee/eh/hh
			S2	S2'	250	500	1000	380	15000	3000		240	365	
κ <sub>W</sub> [%]	1.7	0.75	1.4	0.98	1.8	0.29	0.24	0.86	0.16	0.11	1.3	1.3	0.43	0.14
$\kappa_{Z}$ [%]	1.5	1.2	1.3	0.9	0.29	0.23	0.22	0.5	0.26	0.23	0.14	0.20	0.17	0.12
$\kappa_{\rm g}$ [%]	2.3	3.6	1.9	1.2	2.3	0.97	0.66	2.5	1.3	0.9	1.5	1.7	1.0	0.49
κ <sub>γ</sub> [%]	1.9	7.6	1.6	1.2	6.7	3.4	1.9	98*	5.0	2.2	3.7	4.7	3.9	0.29
$\kappa_{Z\gamma}$ [%]	10.	_	5.7	3.8	99*	86*	85 <b>*</b>	120*	15	6.9	8.2	81*	<b>75</b> ★	0.69
Kc [%]	_	4.1	_	_	2.5	1.3	0.9	4.3	1.8	1.4	2.2	1.8	1.3	0.95
$\kappa_t$ [%]	3.3	_	2.8	1.7	_	6.9	1.6	_	_	2.7	_	_	_	1.0
$\kappa_b  [\%]$	3.6	2.1	3.2	2.3	1.8	0.58	0.48	1.9	0.46	0.37	1.2	1.3	0.67	0.43
$\kappa_{\mu}$ [%]	4.6	_	2.5	1.7	15	9.4	6.2	320*	13	5.8	8.9	10	8.9	0.41
κ <sub>τ</sub> [%]	1.9	3.3	1.5	1.1	1.9	0.70	0.57	3.0	1.3	0.88	1.3	1.4	0.73	0.44

 $k_{\gamma}$ ,  $k_{z_{\gamma}}$  will be the most precise measurements for a long time

## New European Strategy numbers

https://indico.cern.ch/event/1439855/contributions/6461657/

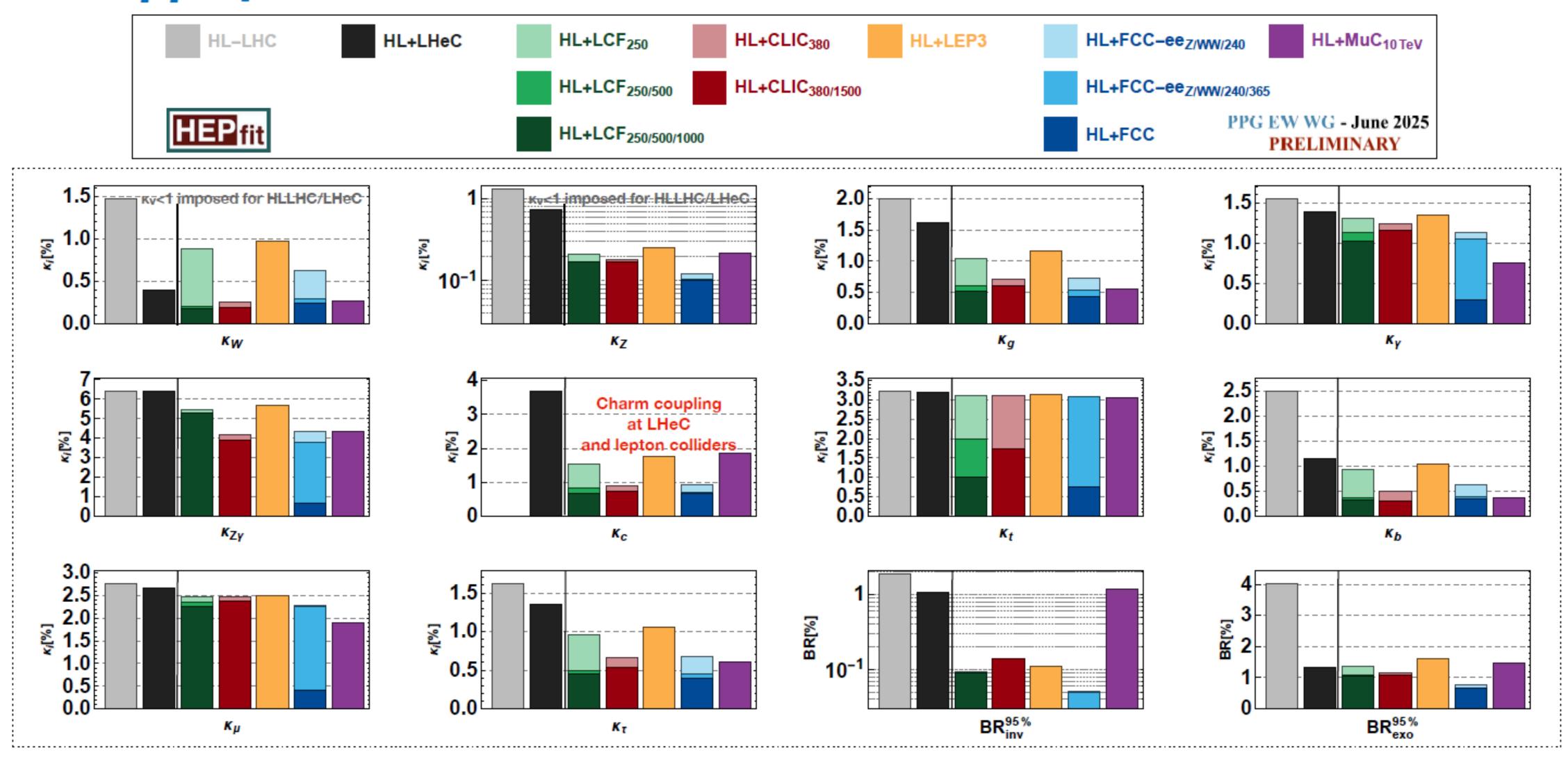
Coupling	HL-LHC	FCC-ee
κ <sub>Z</sub> (%)	1.3* :10	0.10
$\kappa_{\mathrm{W}}$ (%)	1 5*	0.29
$\kappa_{\mathrm{b}}$ (%)	2.5*	0.38 / 0.49
$\kappa_{\rm g}$ (%)	2* :4	0.49 / 0.54
$\kappa_{\tau}$ (%)	1.6*	0.46
$\kappa_{\rm c}$ (%)	_	0.70 / 0.87
$\kappa_{\gamma}$ (%)	1.6*	1.1
$\kappa_{\mathrm{Z}\gamma}$ (%)	10*	4.3
$\kappa_{\mathrm{t}}$ (%)	3.2*	3.1
$\kappa_{\mu}$ (%)	4.4*	3.3
ns  (%)	_	$^{+29}_{-67}$
Γ <sub>H</sub> (%)	- 1	0.78
$\mathcal{B}_{inv}$ (<, 95% CL)	$1.9 \times 10^{-120}$	$5 \times 10^{-4}$
B <sub>unt</sub> (<, 95% CL)	$4 \times 10^{-2}$ *	$6.8 \times 10^{-3}$

The Fcc-ee projections include HL\_LHC, difficult to understand...UNFAIR comparison

In addition results of 4 experiments are combined while only 2 are officially envisaged for the moment.

## New European Strategy numbers

### Kappa framework results



# Higgs Couplings at Fcc-hh at lower energies

Mangano Eurpean Strategy kickoff meeting

Coupling precision	100 TeV CDR baseline	80 TeV	120 TeV
δg <sub>Hγγ</sub> / g <sub>Hγγ</sub> (%)	0.4	0.4	0.4
$\delta g_{H\mu\mu}$ / $g_{H\mu\mu}$ (%)	0.65	0.7	0.6
$\delta g_{HZ_Y} / g_{HZ_Y} (\%)$	0.9	1.0	0.8

A lower energy Fcc-hh can deliver results that are very close to baseline option

## rare Higgs Couplings HL-LHC projections

ATLAS+CMS

		H-	$ ightarrow Z\gamma$	$H  o \mu \mu$		
${\cal L}$		S1	S2	S1	S2	
		$\delta \mu [\%]$	$\delta\mu [\%]$	$\delta \mu [\%]$	$\delta\mu[\%]$	
	ATLAS	31	21	16	13	
$2 {\rm \ ab}^{-1}$	CMS	24	23	9.8	8.4	
	ATLAS+CMS	20	15	7.5	5.9	
	ATLAS	29	17	15	11	
$3 \text{ ab}^{-1}$	CMS	20	19	8.5	7.0	
	ATLAS+CMS	17	14	7.5	5.9	

High pT bins statistically limited, new physics is likely to appear at high pT, large gains still with HL-LHC!

### ATLAS+CMS $H\rightarrow \gamma \gamma p_T > 650 \text{ GeV}$

### ATLASx2 VH→ bb p<sub>T</sub> > 600 GeV

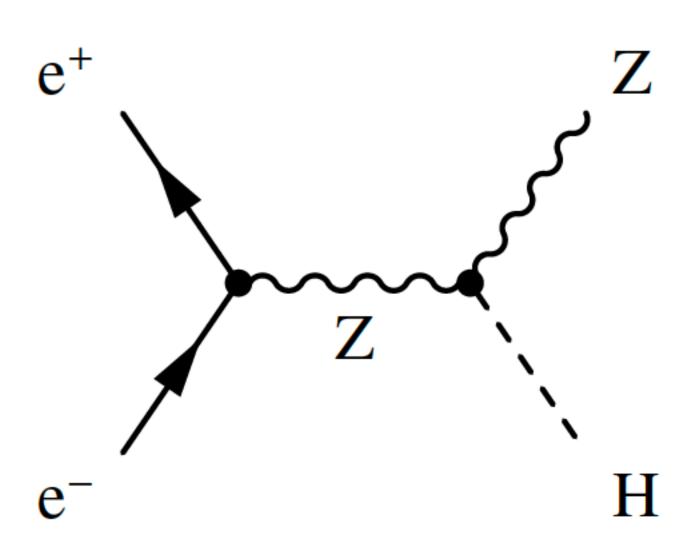
		$\delta\sigma_{\mathrm{ggF}p_{\mathrm{T}}^{H}>65}^{\gamma\gamma}$	50 GeV [%]	$\delta\sigma_{\mathrm{WH}p_{\mathrm{T}}^{W}>600\mathrm{GeV}}^{b\overline{b}}[\%]$	$\delta \sigma_{\mathrm{ZH}p_{\mathrm{T}}^{Z}>600\mathrm{GeV}}^{b\overline{b}}[\%]$
$\mathcal{L}$		SI	<i>S</i> 2	S1	S2
$2 \text{ ab}^{-1}$	ATLAS+CMS	$1.2 \times 10^2$	97	-	-
$3 \text{ ab}^{-1}$	ATLAS+CMS	$1.0 \times 10^2$	79	19	22

# Higgs Mass Future

HL-LHC: 21 MeV FCC-ee: 3 MeV Linear collider: 12 MeV LEP3: 15 MeV

Higgs mass estimates

# Higgs Width @ e+e- Colliders



Precision of Fcc-ee 4%
Precision of ILC(250)~11%
Precision of ILC(500) 5%
LEP3 (H->ZZ) 21%
LEP3(WWW fusion) 12%

Measuring the inclusive production cross-section of the ZH process without presuming specific Higgs boson decay modes allows the determination of the absolute coupling between the Higgs and Z bosons. Combining this measurement with data on the ZH,  $H \rightarrow ZZ*$  process, along with the  $H \rightarrow ZZ*$  branching ratio, allows for the calculation of the total Higgs boson width without any additional theoretical assumptions:

$$\Gamma_H \propto \frac{(\sigma_{e^+e^- o ZH} \times BR_{H o ZZ^*})^2}{\sigma_{e^+e^- o ZH}}$$

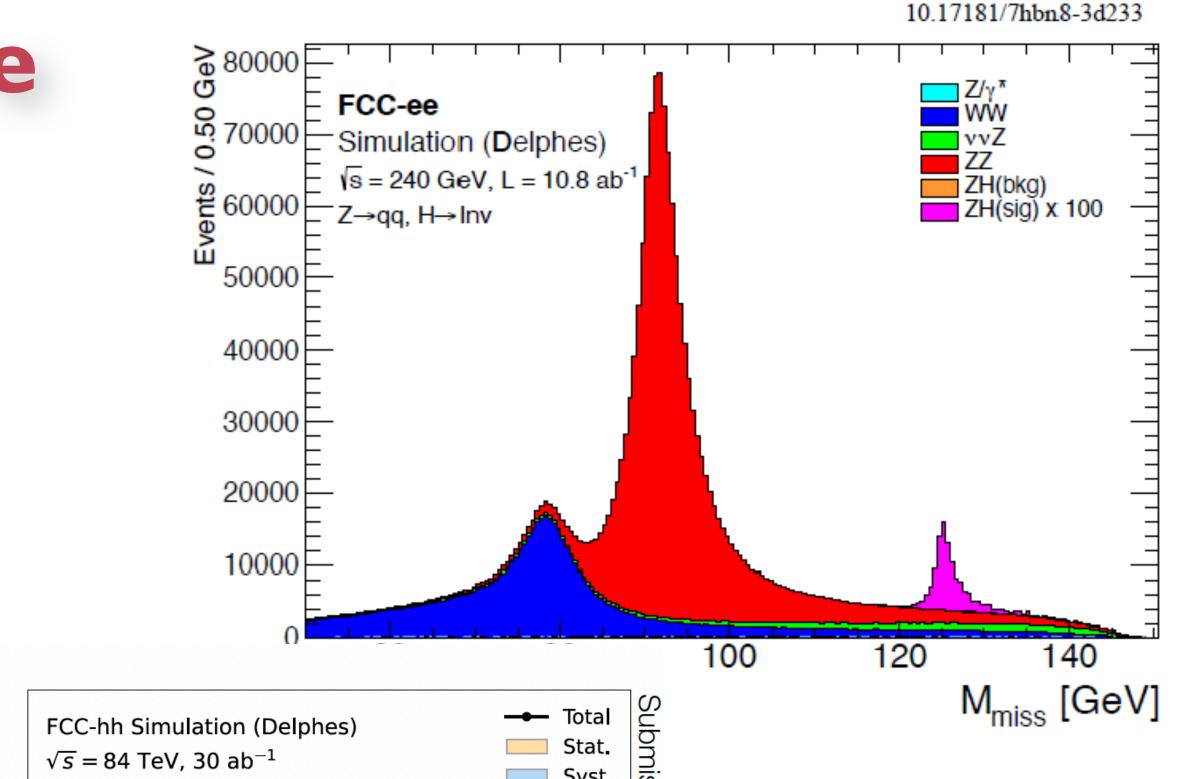
Cannot be done at pp because of invisible decays (e.g. Z to neutrinos)

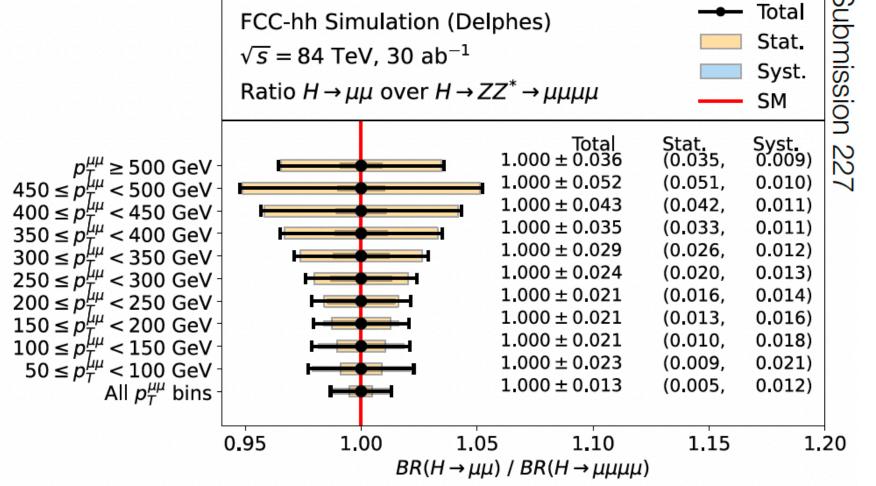
### Higgs to invisible future

### Fcc-ee gets to 0.05%-0.12% limit

Collider	Projected sensitivity (95% CL)
LHC (current)	~13%–19%
HL-LHC	2%
LEP3 / TLEP	~1% (exclude >2%)
FCC-ee / CEPC	≤ 0.1% (per mille level)
ILC (250 GeV)	~0.3%

### Fcc-hh Higgs to invisible limit of 0.02%

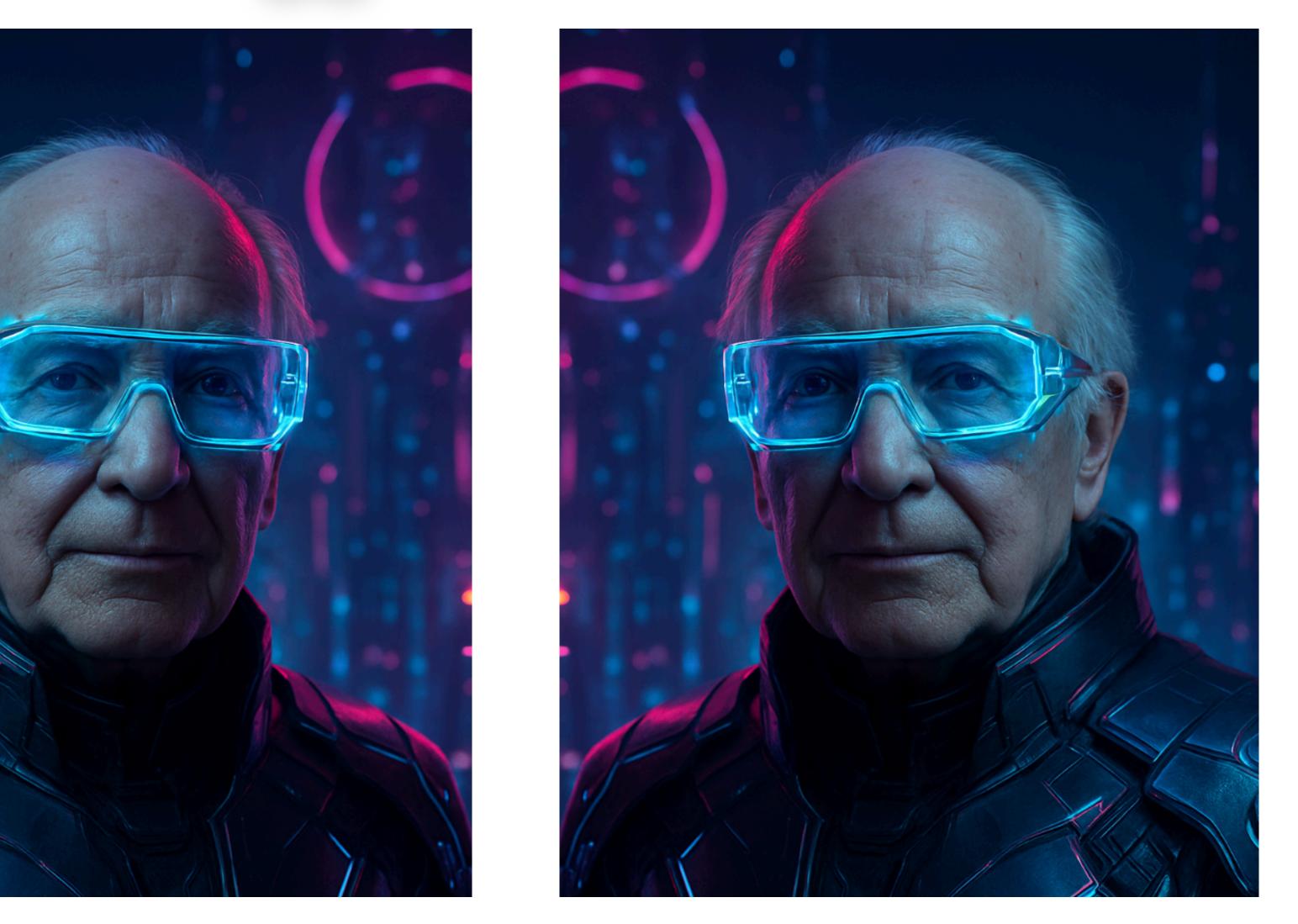




Uses H→µµµµ from FCCee to estimate cross section

### 2 Higgses Future





## Example of improvements vs time: ATLAS HH → bbtt

ATL-PHYS-PUB-2024-016

2015 HL\_LHC @3000 fb-1 0.6 sigma

ATL-PHYS-PUB-2015-046

2021 HL\_LHC @3000 fb<sup>-1</sup>
2.8 sigma

ATL-PHYS-PUB-2021-044

2024 HL\_LHC @3000 fb-1
3.5 sigma

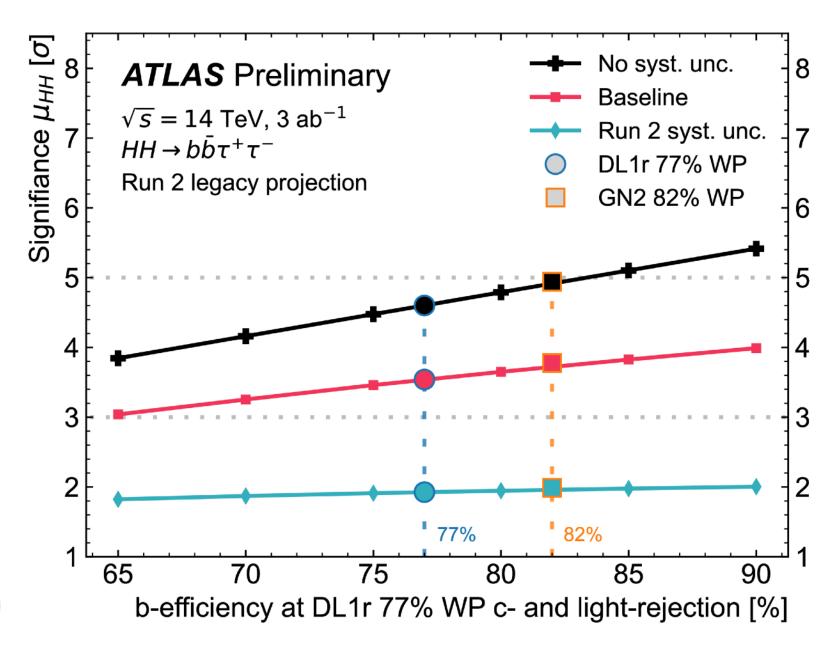
ATL-PHYS-PUB-2024-016

Table 9: Expected significance for several channel combinations, for a luminosity of 3 ab<sup>-1</sup>, including the expected uncertainties quoted in the text, using the asymptotic approximation. This table only takes into account the  $\tau_{lep}\tau_{had}$  and  $\tau_{had}\tau_{had}$  channels.

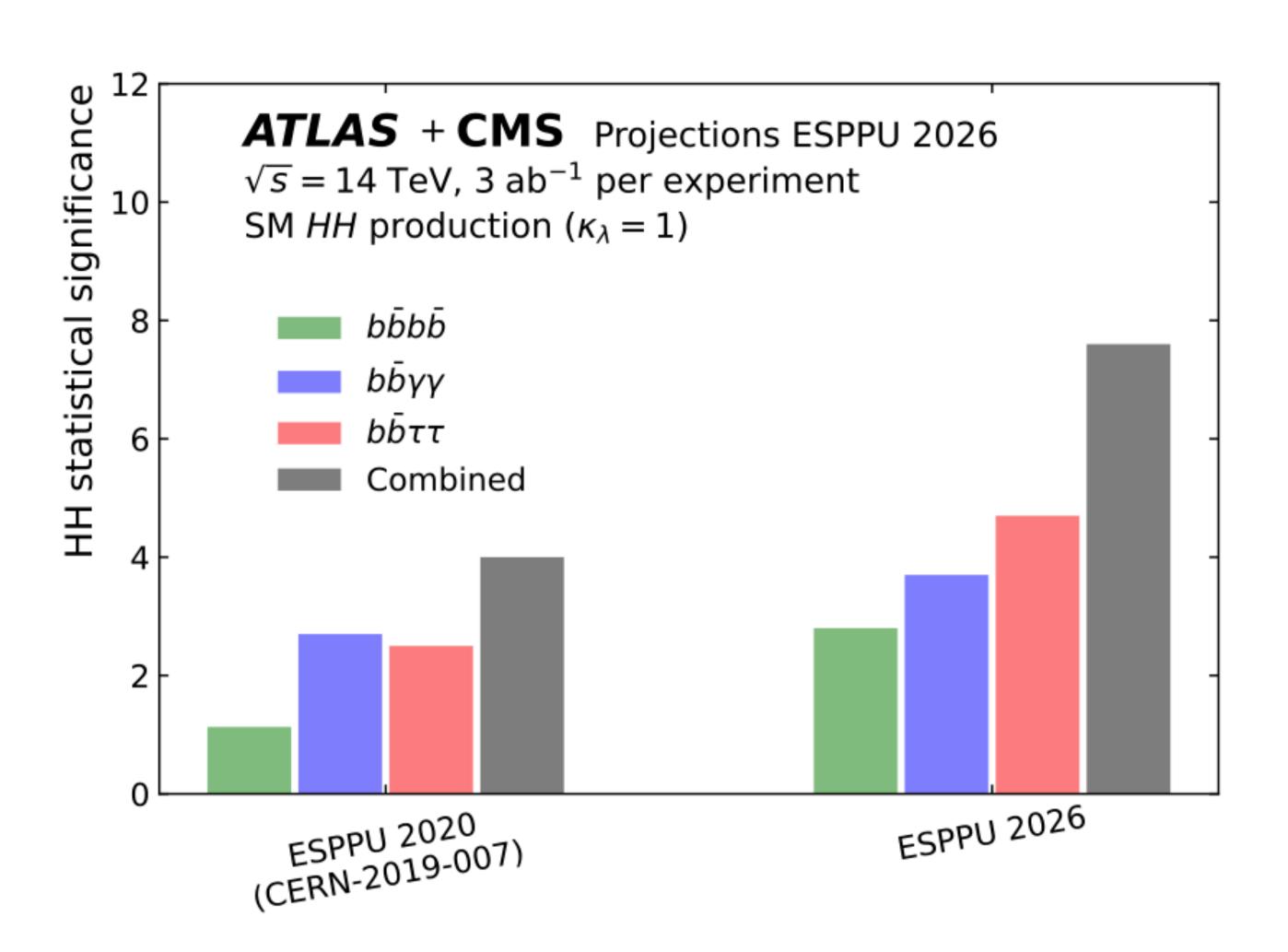
Channel	Significance	Combined in channel	Total combined	
e + jets	0.31	0.43		
μ+jets	0.30	0.43	0.60	
$ au_{ m had} au_{ m had}$	0.41	0.41		

assuming SM kinematics for the signal, the 95% CL upper limit on the HH signal strength is projected to be at 0.71 times the SM prediction with respect to the background-only hypothesis assuming an integrated luminosity of 3000 fb<sup>-1</sup> and  $\sqrt{s} = 14$  TeV. The signal significance is extrapolated to  $2.8\sigma$ , and assuming a true value of  $\kappa_{\lambda} = 1$ , the self-coupling modifier is constrained to the  $1\sigma$  CI [0.3, 1.9]  $\cup$  [5.2, 6.7]. If no HH

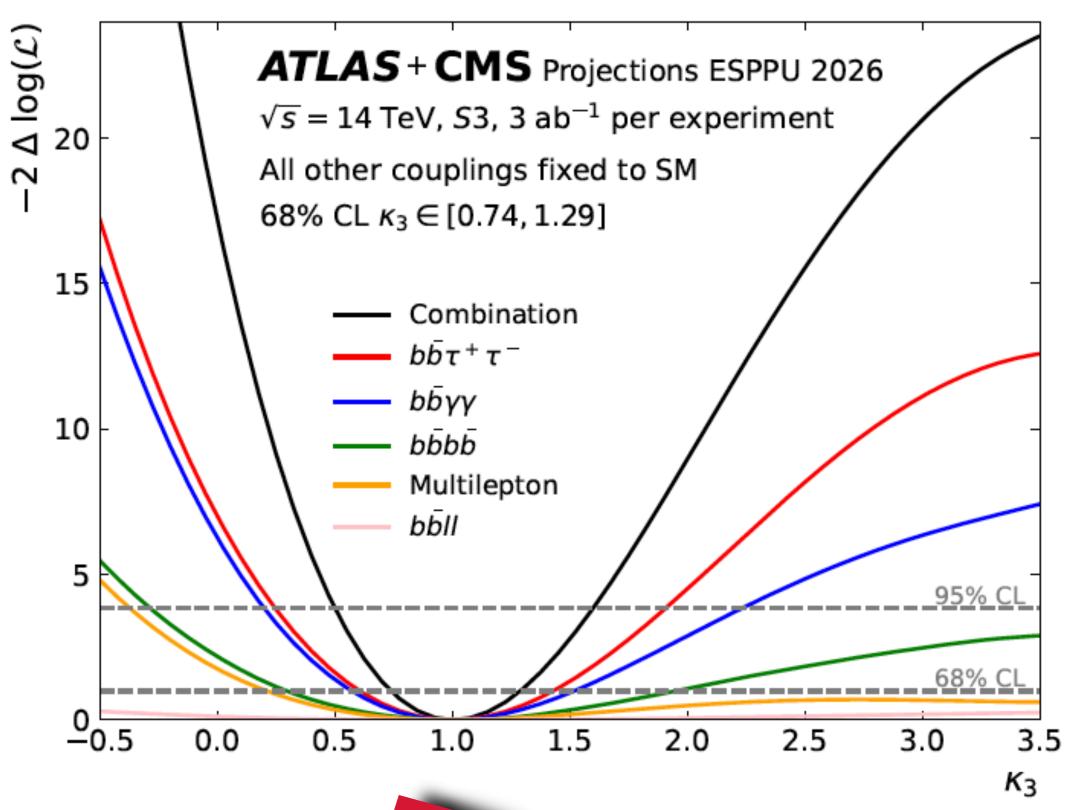
a variety of integrated luminosities ranging from 1000 to 3000 fb<sup>-1</sup>. Assuming SM HH production, a signal significance of 3.5  $\sigma$  (4.6  $\sigma$ ) is expected in the baseline (statistical only) extrapolation scenario for an integrated luminosity of 3000 fb<sup>-1</sup>. This translates into expected

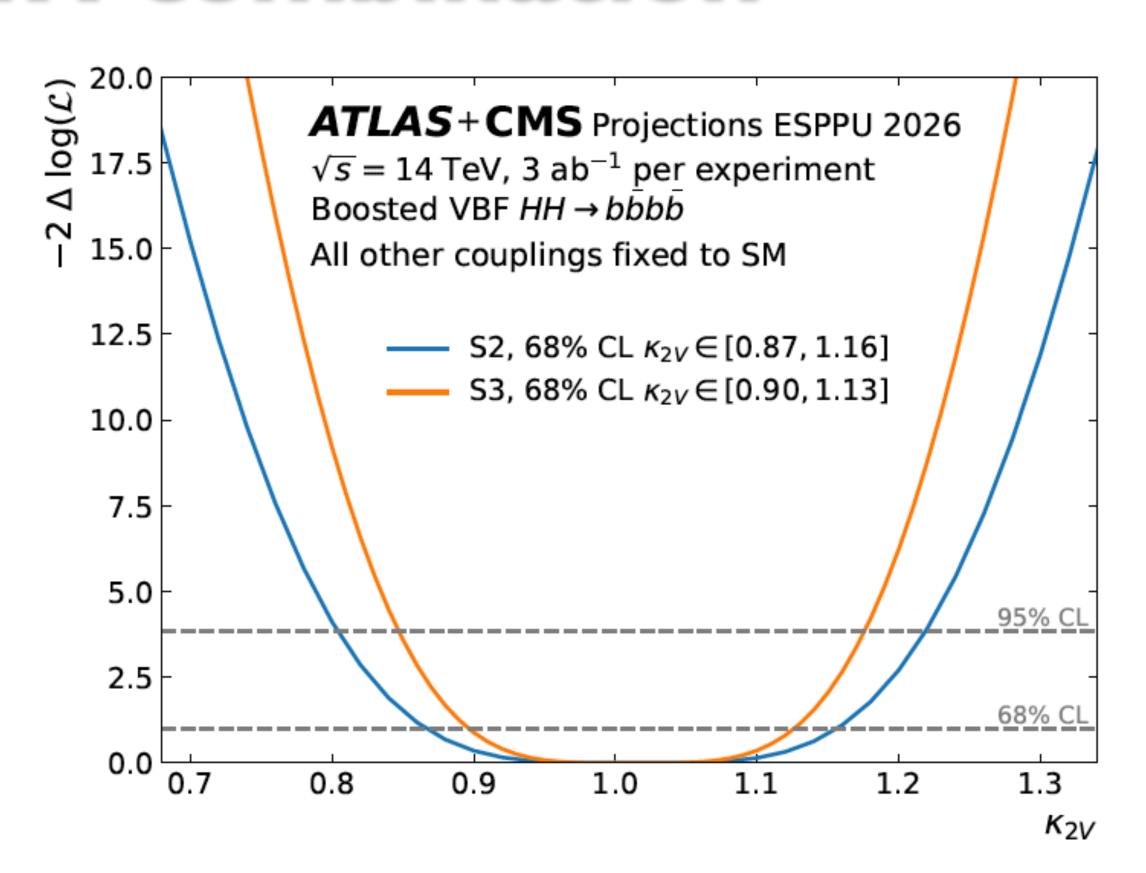


#### Improvements in projections since 2020



#### ATLAS+CMS HH combination







 $HH \rightarrow bb\gamma\gamma$  both ATLAS and CMS,

HH→bbττ, →bbll, →multileptons ATLASx2

HH→bbbb resolved and boosted CMSx2

#### Non-resonant HH combination

	$2 \text{ ab}^{-1} (S2)$		3 ab	<sup>1</sup> (S2)	$3 \text{ ab}^{-1} (S3)$	
	<b>ATLAS</b>	CMS	ATLAS	CMS	ATLAS	CMS
HH statistical significance						
$b\bar{b}\tau^+\tau^-$	$3.0^{\dagger}$	1.9	3.5 <sup>†</sup>	2.4	3.8 <sup>†</sup>	2.7
$b ar{b} \gamma \gamma$	$2.1^{\dagger}$	$2.0^{\dagger}$	$2.4^{\dagger}$	$2.4^{\dagger}$	$2.6^{\dagger}$	$\boldsymbol{2.6}^{\dagger}$
$b ar{b} b ar{b}$ resolved	0.9	$1.0^{\dagger}$	1.0	$1.2^{\dagger}$	1.0	$1.3^{\dagger}$
$b\bar{b}b\bar{b}$ boosted	_	$1.8^{\dagger}$	_	$2.2^{\dagger}$	_	$\boldsymbol{2.2}^{\dagger}$
Multilepton	$0.8^{\dagger}$	_	$1.0^{\dagger}$	_	$\boldsymbol{1.0}^{\dagger}$	_
$b\bar{b}\ell^+\ell^-$	$0.4^{\dagger}$	_	$0.5^{\dagger}$	_	$0.5^{\dagger}$	_
Combination	3.7	3.5	4.3	4.2	4.5	4.5
ATLAS+CMS	6.0		7.2		7.6	
		$\kappa_3$ 68	8% confidence i	nterval		
$b\bar{b}\tau^+\tau^-$	$[0.3, 1.8]^{\dagger}$	[0.1, 3.0]	$[0.4, 1.7]^{\dagger}$	[0.2, 2.2]	$\left[0.5,\ 1.6\right]^{\dagger}$	$[0.3, \ 2.0]$
$b \overline{b} \gamma \gamma$	$[0.3, \ 2.0]^{\dagger}$	$[0.2, \ 2.3]^{\dagger}$	$[0.4, 1.8]^{\dagger}$	$[0.3, \ 2.0]^{\dagger}$	$\left[0.5,\ 1.7\right]^{\dagger}$	$\left[0.4,\ 1.9\right]^{\dagger}$
$b \bar{b} b \bar{b}$ resolved	[-0.7, 6.3]	$[-0.6, 7.6]^{\dagger}$	[-0.5, 6.1]	$[-0.3, 7.3]^{\dagger}$	[-0.5, 6.1]	$\left[-0.3,\ 7.2\right]^{\dagger}$
$b\bar{b}b\bar{b}$ boosted	_	$[-0.6, 8.5]^{\dagger}$	_	$[-0.4, 8.2]^{\dagger}$	_	$\left[-0.4,\ 8.2\right]^{\dagger}$
Multilepton	$[-0.2, 4.9]^{\dagger}$	_	$[-0.1, 4.7]^{\dagger}$	_	$\left[-0.1,\ 4.7\right]^{\dagger}$	_
$b\bar{b}\ell^+\ell^-$	$[-2.4, 9.3]^{\dagger}$	_	$[-2.2, 9.2]^{\dagger}$	_	$[-2.1,\ 9.1]^{\dagger}$	_
Combination	[0.6, 1.5]	[0.4, 1.7]	[0.6, 1.5]	[0.5, 1.6]	[0.6, 1.4]	$[0.6,\ 1.5]$
ATLAS+CMS uncertainty  † used in the AT	-32% / +37%		-27% ,	/ +31%	$-26\%\ / +29\%$	

Close to single experiment observation!

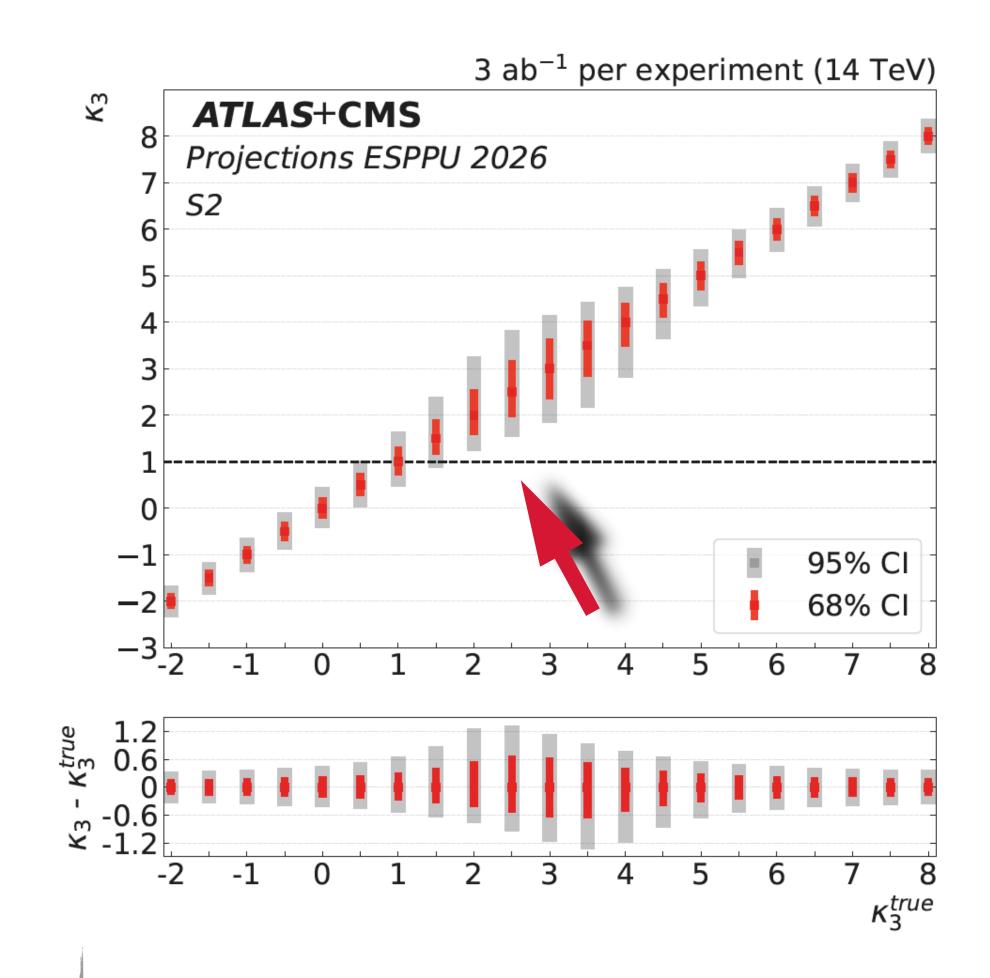


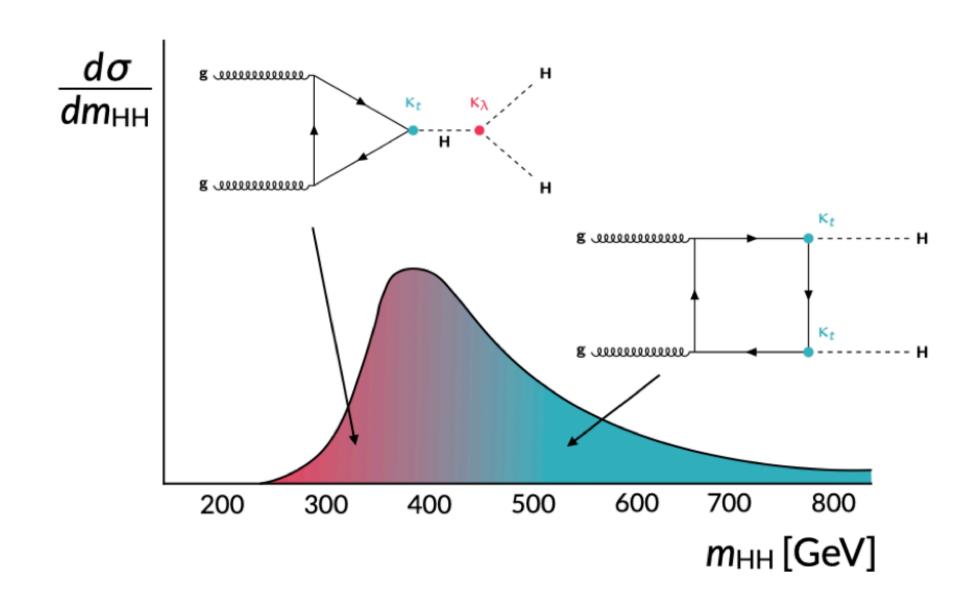
■  $7.6\sigma \Leftrightarrow 4.0\sigma$  ESU2020

† used in the ATLAS+CMS combination

@3ab-1 ATLAS and CMS will get to less than 30% accuracy on Higgs self couplings, with conservative assumptions!

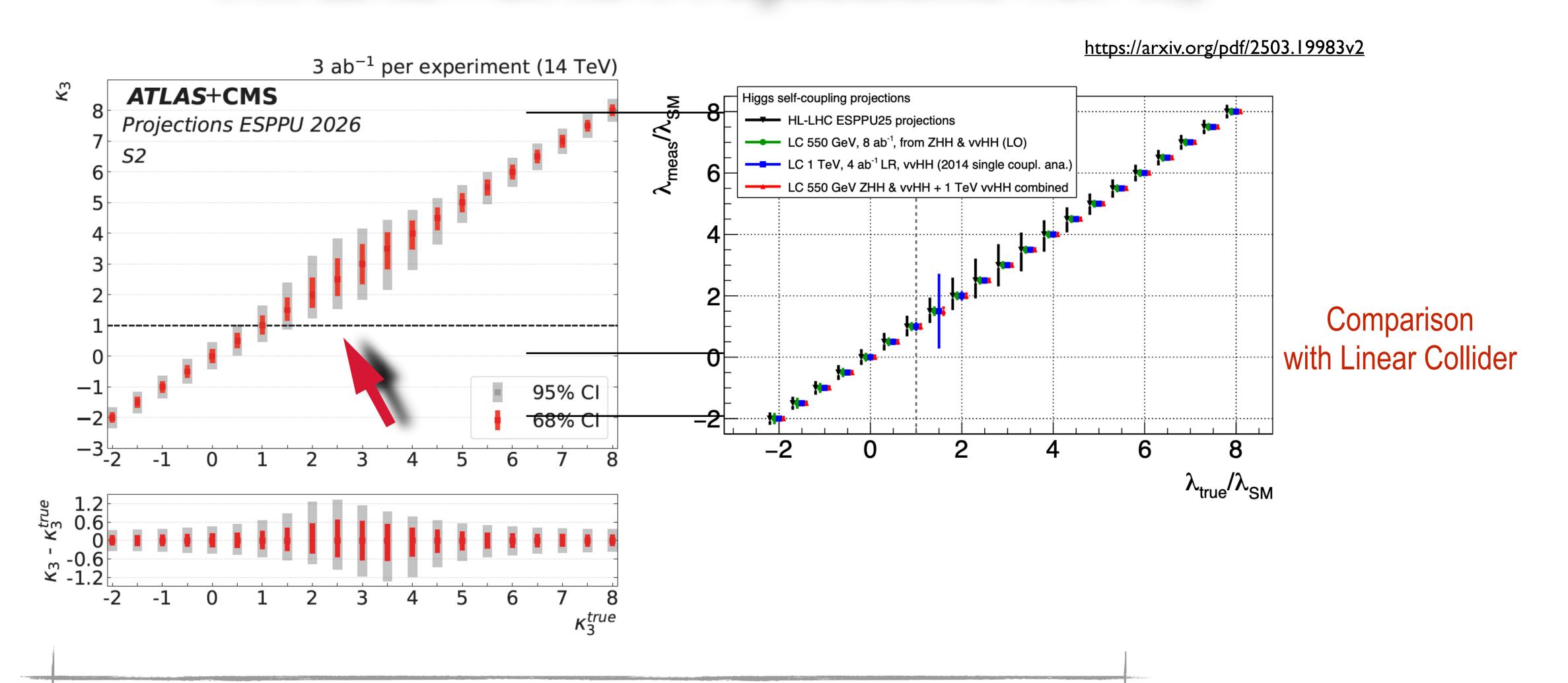
#### ATLAS+CMS Projections for non SM k<sub>3</sub>





We can exclude at more than 2 sigma the  $k_3$ =1 around  $k_3$ ~2.5, where we have maximal interference among box and triangle diagrams

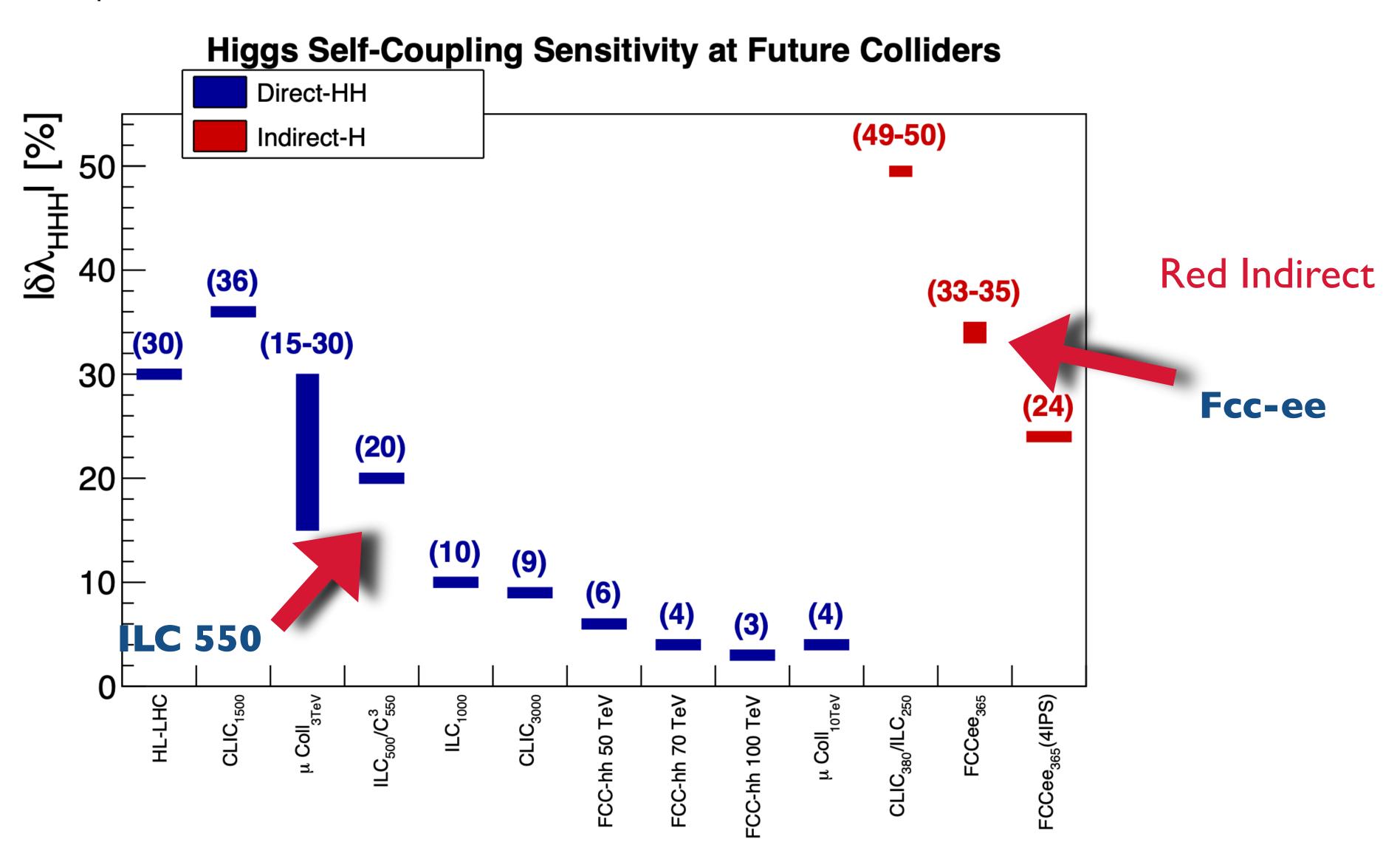
#### ATLAS+CMS Projections for k<sub>3</sub>



We can exclude at more than 2 sigma the  $k_3$ =1 around  $k_3$ ~2.5, where we have maximal interference among box and triangle diagrams

#### di-higgs precision European Strategy 2025

https://indico.cern.ch/event/1439855/contributions/6461643/attachments/3046021/5381998/NearTermProtonCollider.pdf



Blue direct

#### Fcc-hh % precision on k<sub>λ</sub>

• assuming s = 100TeV,  $\mathcal{L}$ int = 30 ab-1@68%CL

Mangano, Ortona, Selvaggi, Eur. Phys. J. C80 (2020)]

I00 TeV	s l	s II	s III
stat	3.0	<b>4.</b> I	5.6
syst	1.6	3.0	5.4
tot	3.4	<b>5.1</b>	7.8

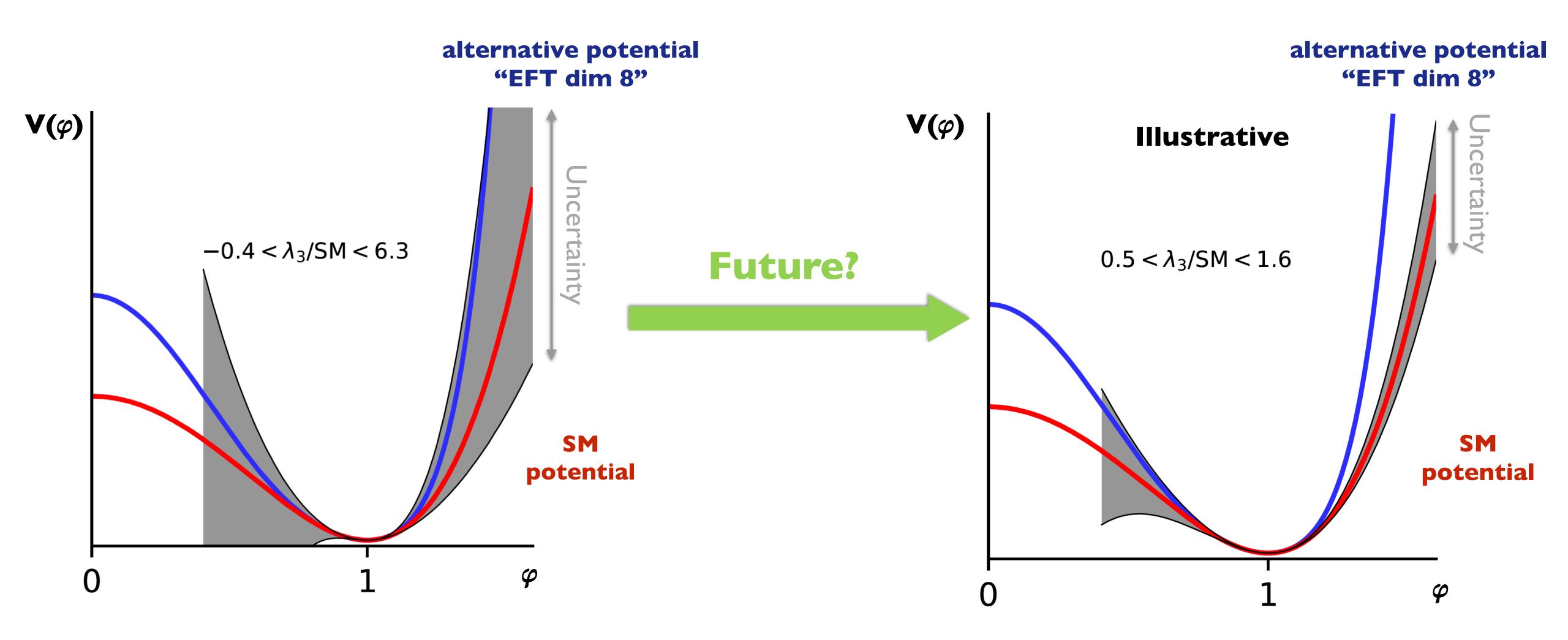
- Scenario I: Optimistic target detector performance, similar to Run 2 LHC conditions.
- Scenario II: Realistic intermediate detector performance.
- Scenario III: Conservative pessimistic detector performance, assuming extrapolated HL-LHC performance using present-day algorithms.

Mangano, FCC-hhStudiesfornextEuropStrategyKick.off, 9/24

80 TeV	s l	s II	s III
stat	3.5	4.7	6.4
syst	1.6	3.0	5.4
tot	3.8	5.6	8.4

- → decreasing to 80 TeV the centre of mass energy.
- → Interesting options also at 60 TeV
- ⇒ both can be achieved with 12 Tesla existing magnet technology)

#### the Higgs Potential



What we know today

## Generic BSM potential & First Order Phase Transition (FOPT)

Consider a sufficiently general class of (non-SM) EWSB potentials, for which it is known whether they lead to a first-order phase transition, and determine their predictions for k3, the Higgs self coupling.

#### Effective Field Theories (EFT):

- EFT dim 6 (k3 and k4 are not independent)
- EFT dim 8 (k3 and k4 independent)

Logarithmic Potential

**Exponential Potential** 

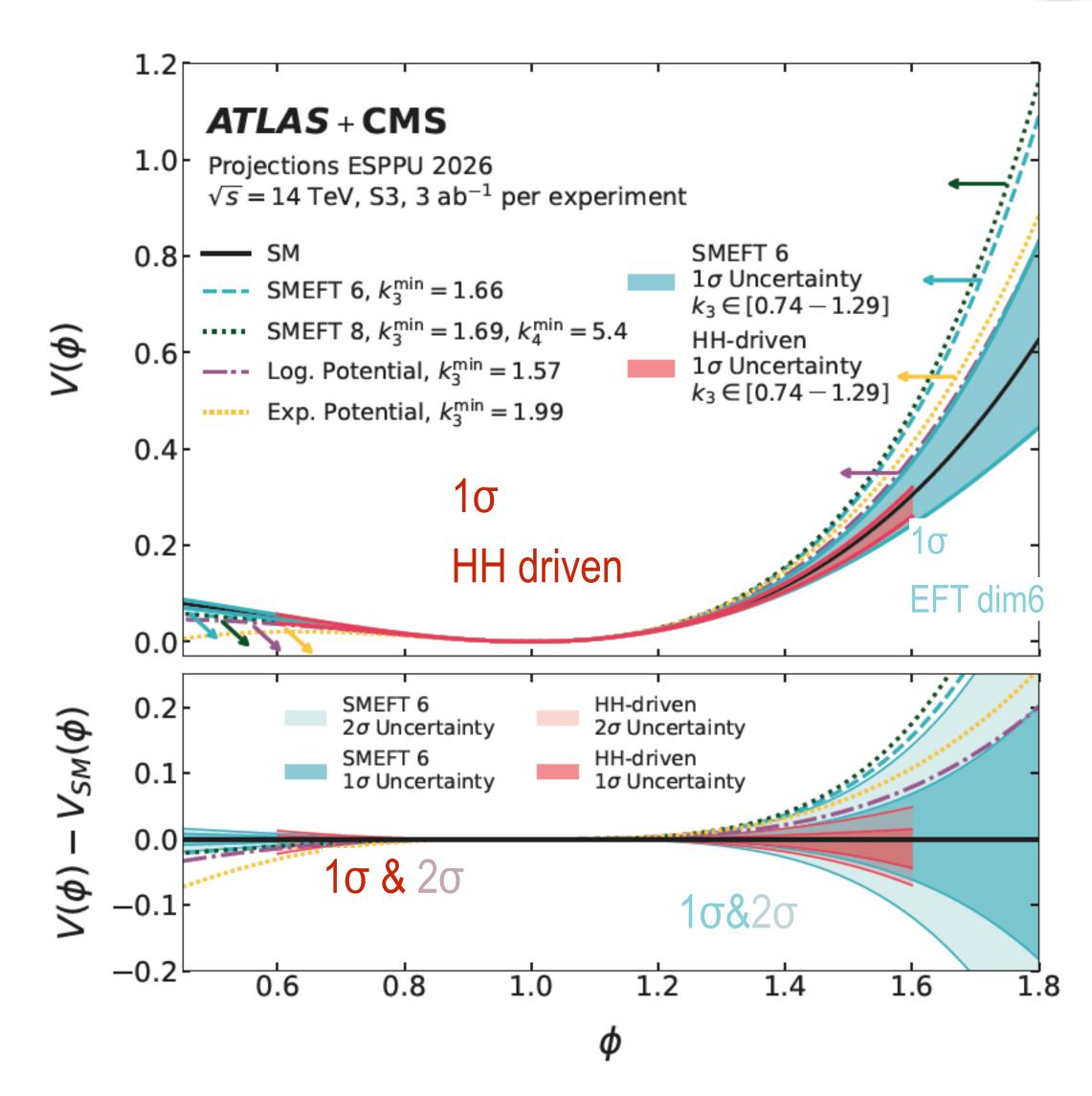
$$V_6^{ ext{SMEFT}}(\phi) = rac{1}{8} \left(\phi^2 - 1\right)^2 + rac{k_3 - 1}{16} \left(\phi^2 - 1\right)^3$$

$$V_8^{\text{SMEFT}}(\phi) = \frac{1}{8} \left( \phi^2 - 1 \right)^2 + \frac{k_3 - 1}{16} \left( \phi^2 - 1 \right)^3 + \frac{k_4 - 6k_3 + 5}{128} \left( \phi^2 - 1 \right)^4 .$$

$$V^{\text{CW}}(\phi) = \frac{1}{8} \left( \phi^2 - 1 \right)^2 + \frac{3}{8} (k_3 - 1) \left[ \left( \frac{1}{2} \phi^4 - 1 \right) - \phi^2 \log \phi^2 \right].$$

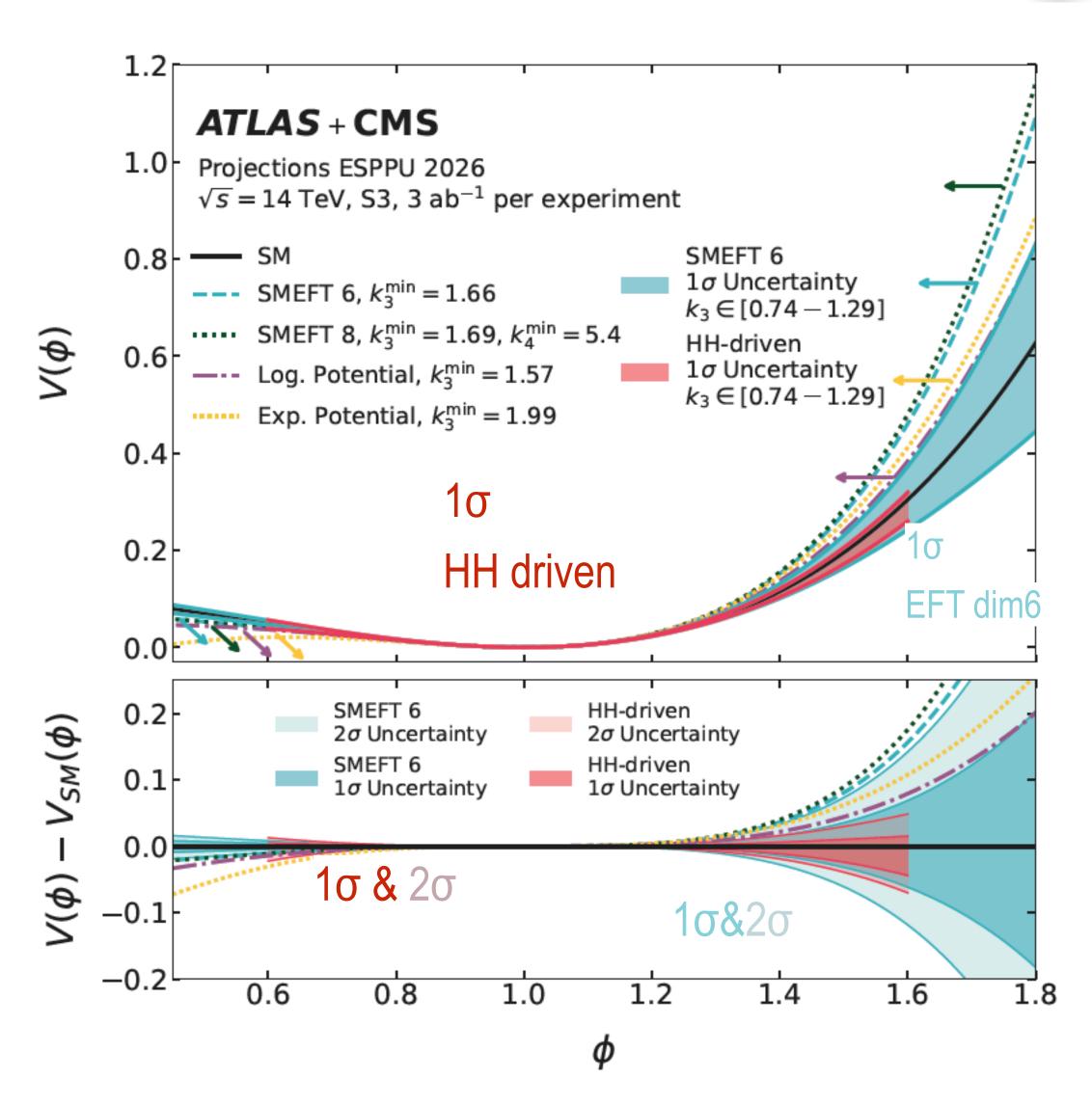
$$V^{I}(\phi) = \frac{1}{8} \left( \phi^{2} - 1 \right)^{2} \pm \phi^{4} e^{-\frac{2}{c_{I}\phi^{2}}} + \frac{e^{-2/c_{I}} \pm \left( -(c_{I}(c_{I} + 2) + 2)\phi^{4} + 2(c_{I} + 2)\phi^{2} - 2 \right)}{c_{I}^{2}}$$

## Higgs potential and generic model strong FOPT



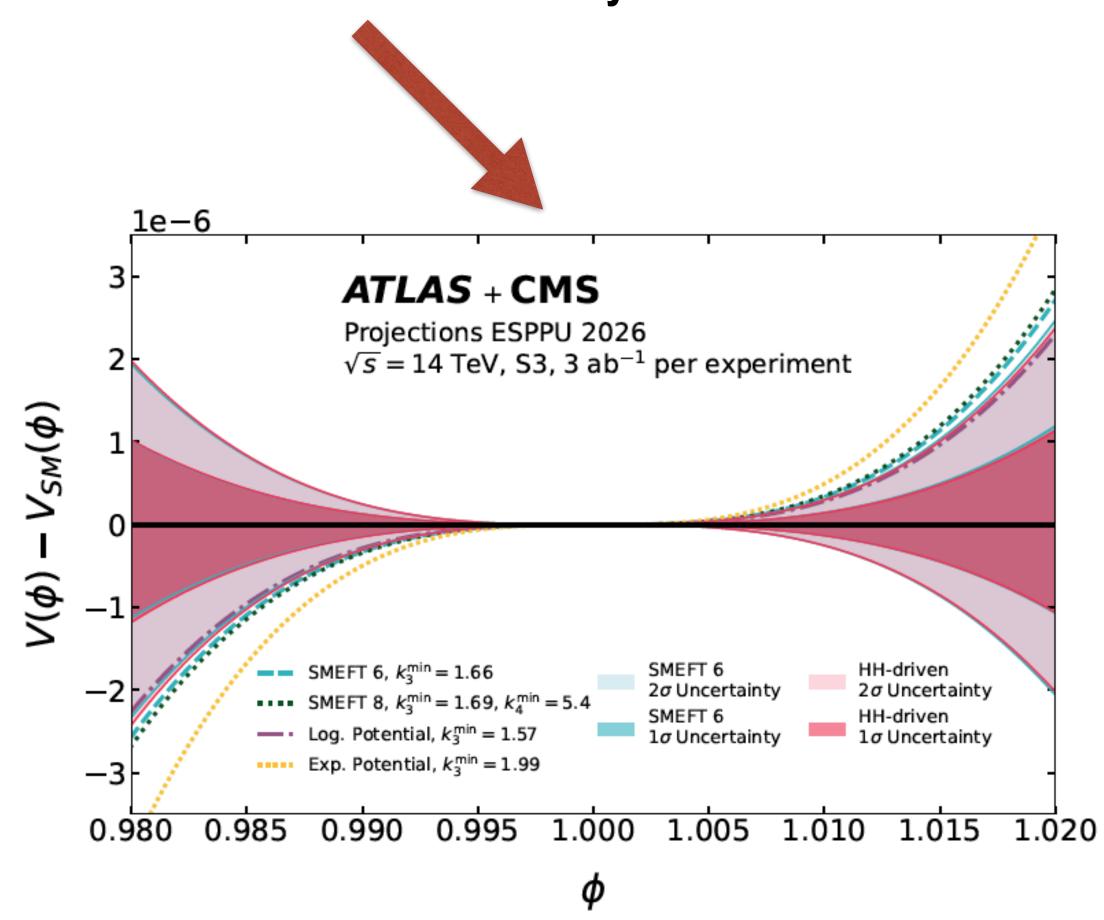
Around the minimum: all potentials coincide, the HH driven curve is valid only around the minimum. It is experimentally driven and agnostic of k<sub>4</sub> values

## Higgs potential and generic model strong FOPT

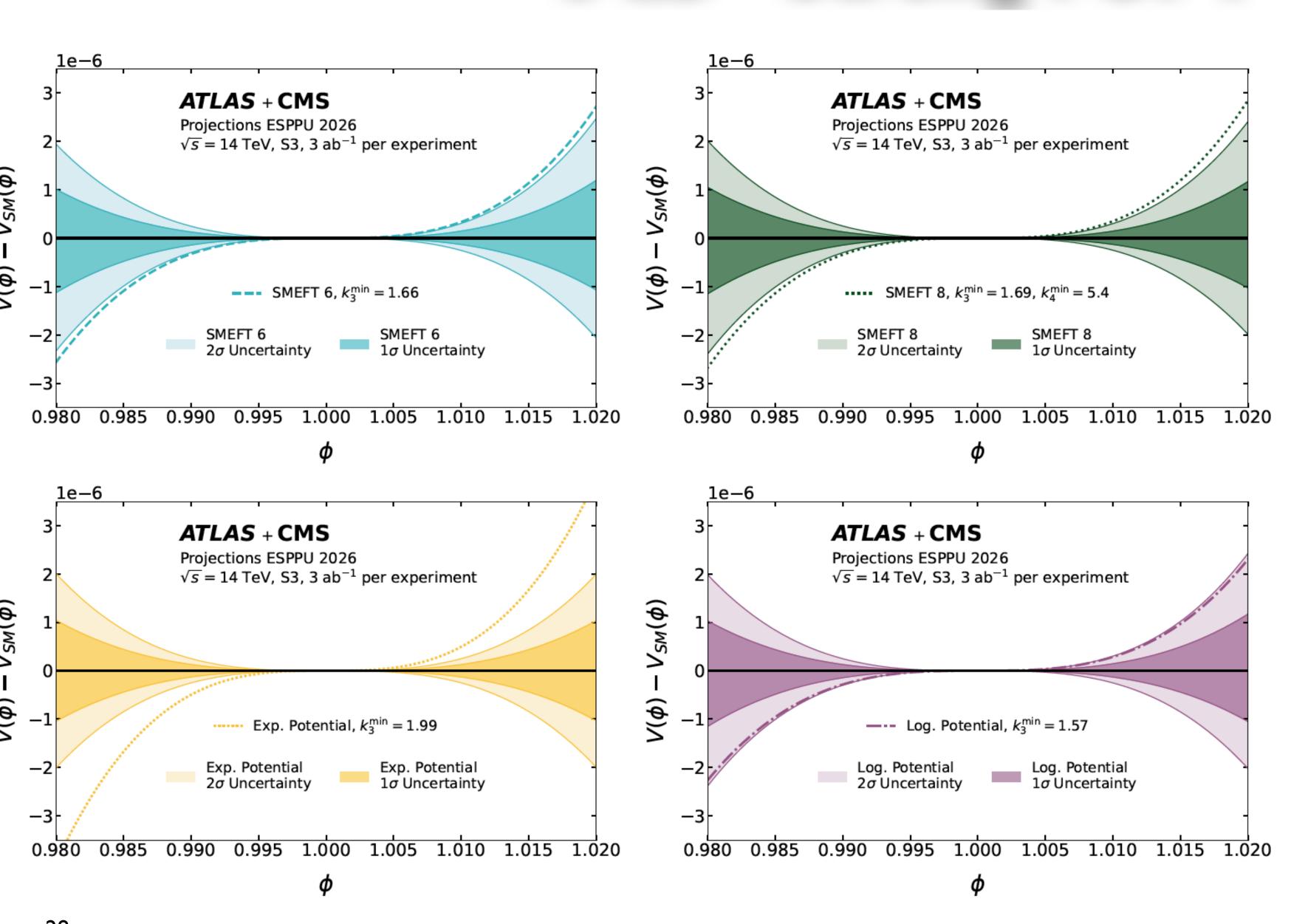


#### Around the minimum: all potentials coincide,

Zoomed version shows that the red curve that is data driven excludes fully all FOPT!



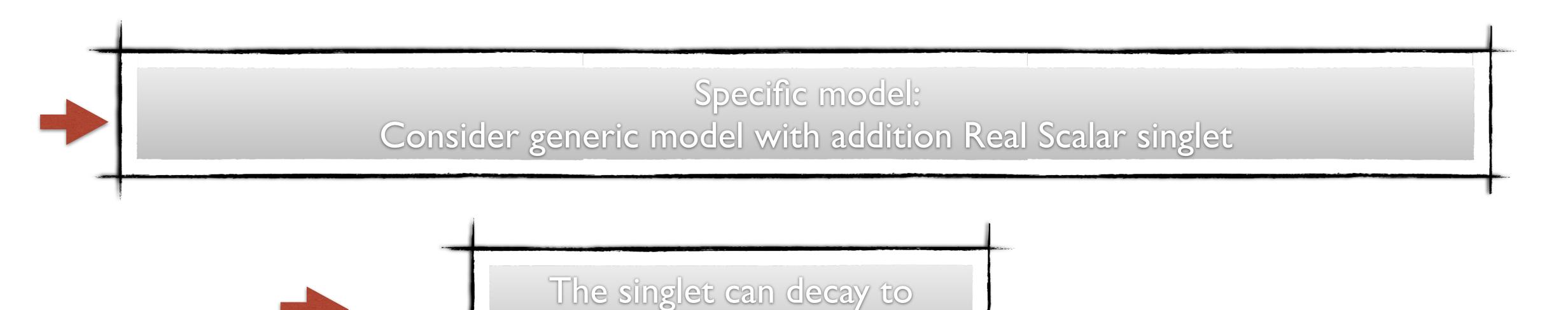
#### 3 ab-1 Strong FOPT



For the exclusion within each individual potential "model" we have to compare the dashed lines with the respective error band

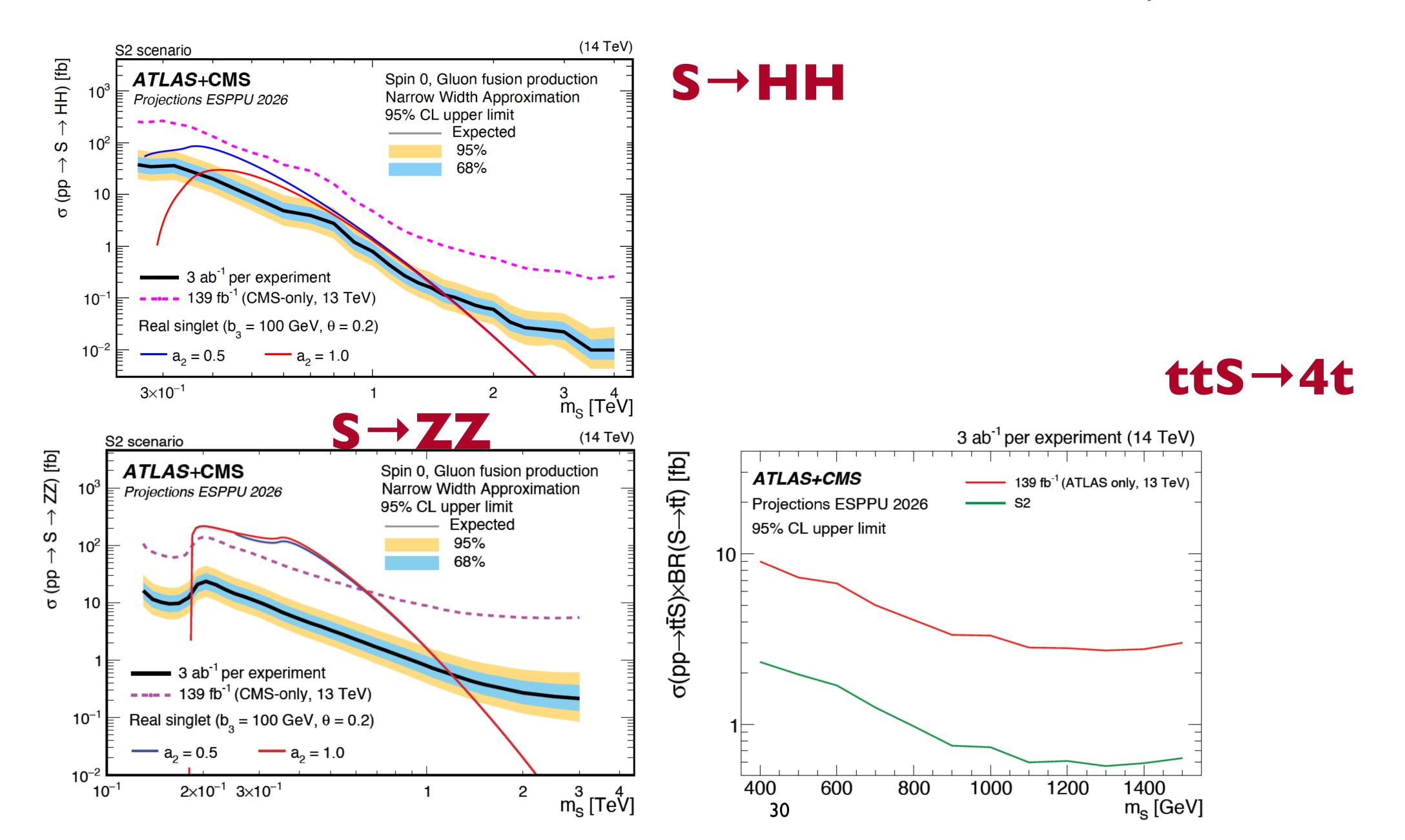
# A specific model implementation Scalar Singlet Model and FOPT

$$V(\Phi, S) = -\mu_H^2 |\Phi|^2 + \lambda_H |\Phi|^4 + b_1 S - \frac{\mu_S^2}{2} S^2 + \frac{b_4}{4} S^4 + \frac{a_2}{2} |\Phi|^2 S^2 + \frac{b_3}{3} S^3 + \frac{a_1}{2} |\Phi|^2 S.$$



 $S \rightarrow HH/ZZ$ , ttS $\rightarrow$ tttt

#### Scalar resonant searches S→HH/ZZ, ttS→4t



#### Scalar Singlet Model and FOPT

 $m_S = 300 \text{ GeV}, b_3 = 0 \text{ GeV}, b_4 = 0.25$ 

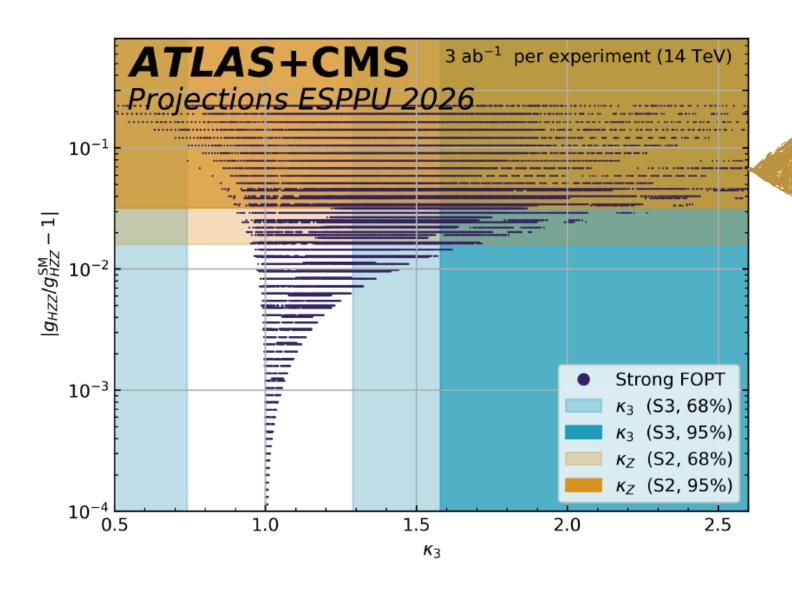
Conservative choice of model parameters

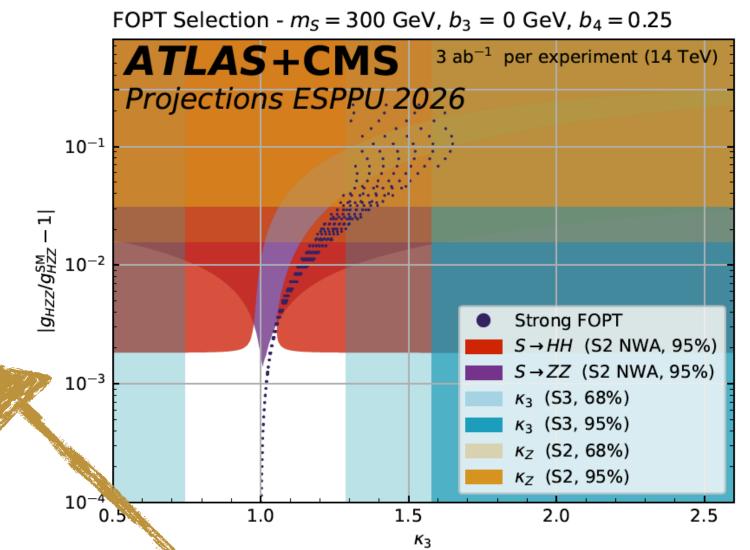
Full FOPT exclusion in other regions!

ATLAS+CMS
7 Projections ESPPU 2026

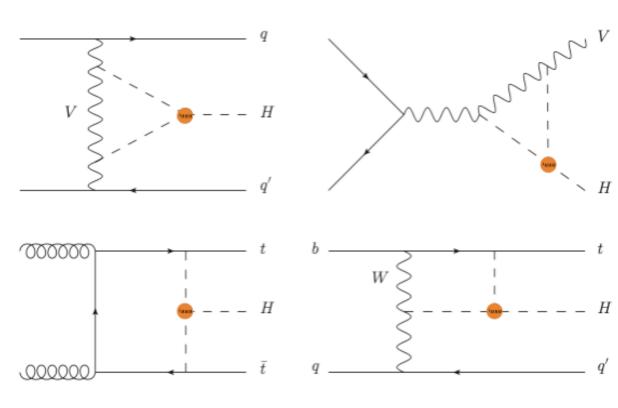
S→HH (S2, NWA, 95%)
S→ZZ (S2, NWA, 95%)
K₂ (S2, 95%)
Strong FOPT

4
3
2
1
0
-0.25-0.20-0.15-0.10-0.05 0.00 0.05 0.10 0.15 0.20 0.25





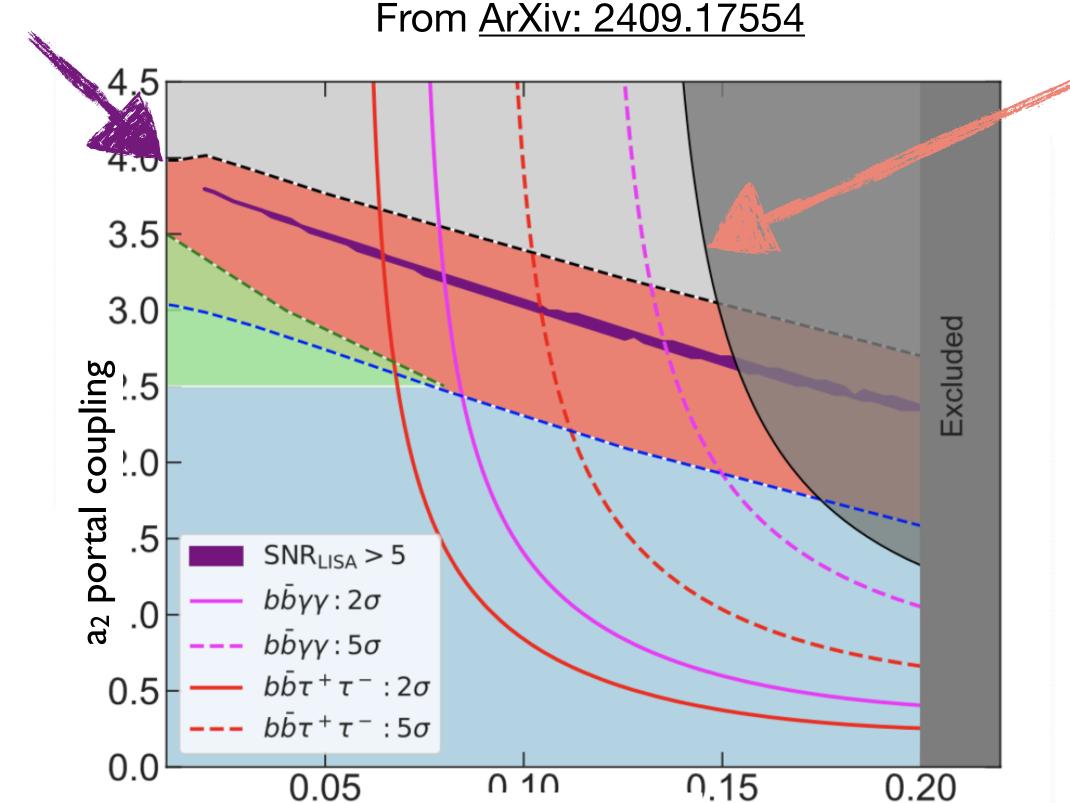
Single H boson production includes processes sensitive to the Higgs self couplings at NLO EW correction level. Constraint from k<sub>z</sub> useful



#### Comparison with Lisa sensitivity

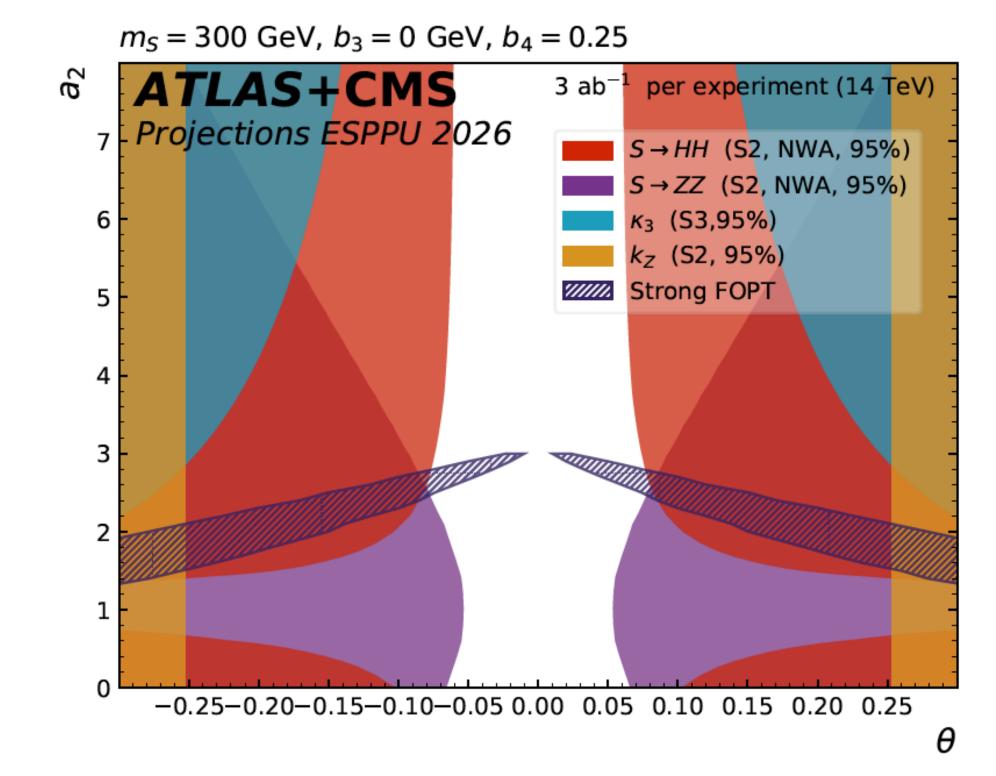
Gravitational waves experiment:

Lisa sensitivity



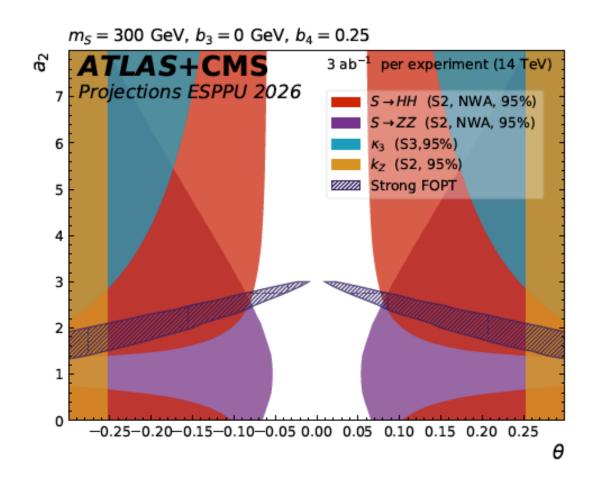
Mixing angle.

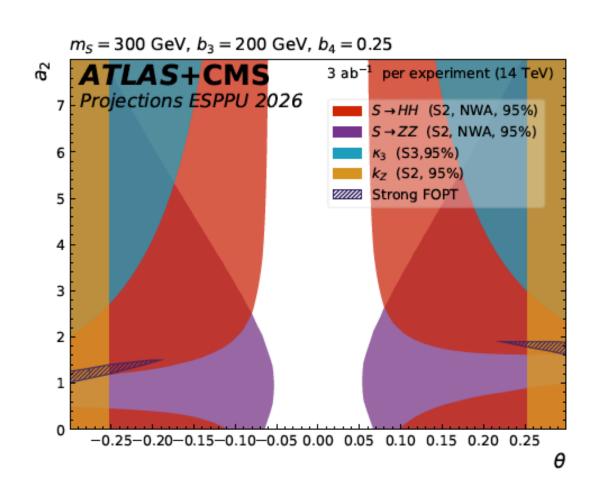
Models implying the existence of an additional scalar singlet that would imply a strong 1st order phase transition (orange region)

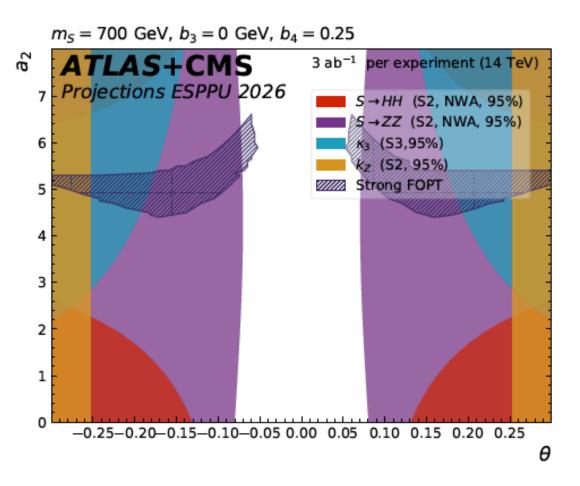


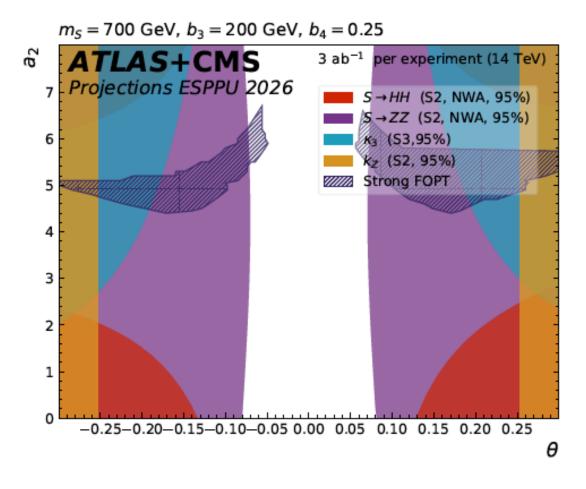
Nice complementarity with Gravitational Wave experiments

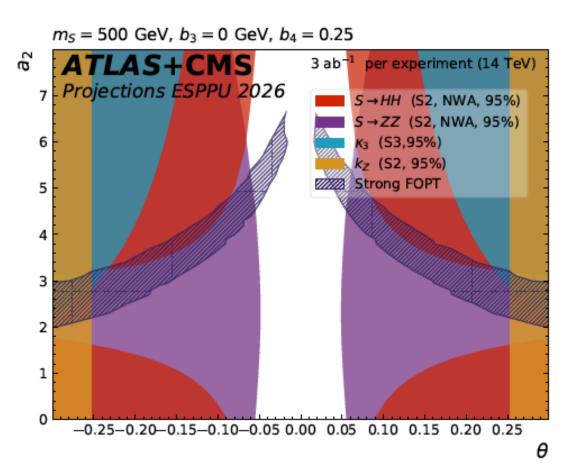
#### sample of different model parameters

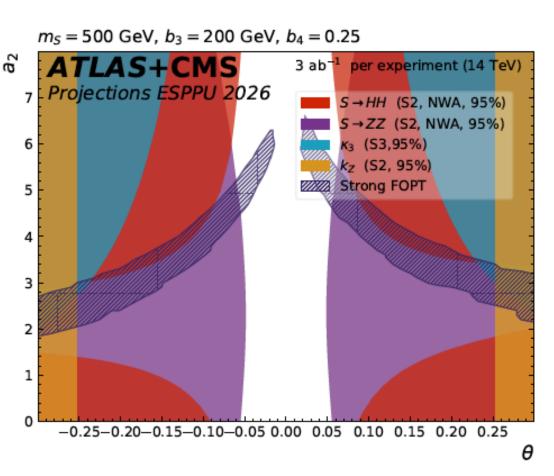


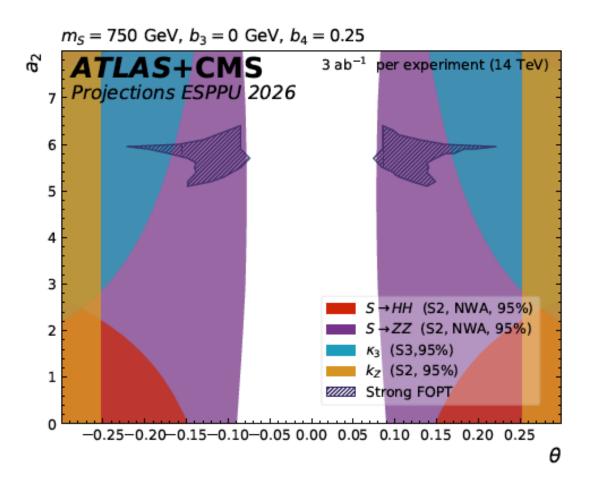


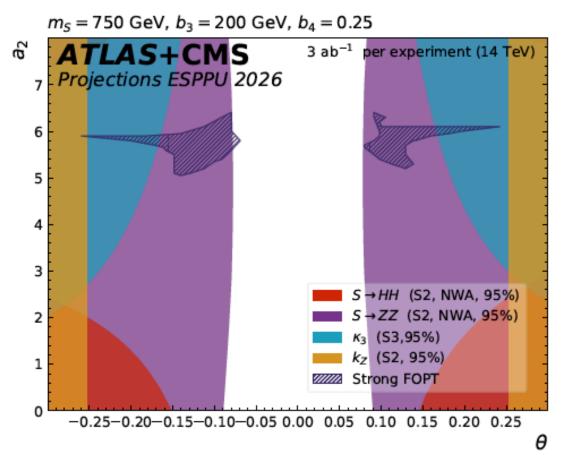






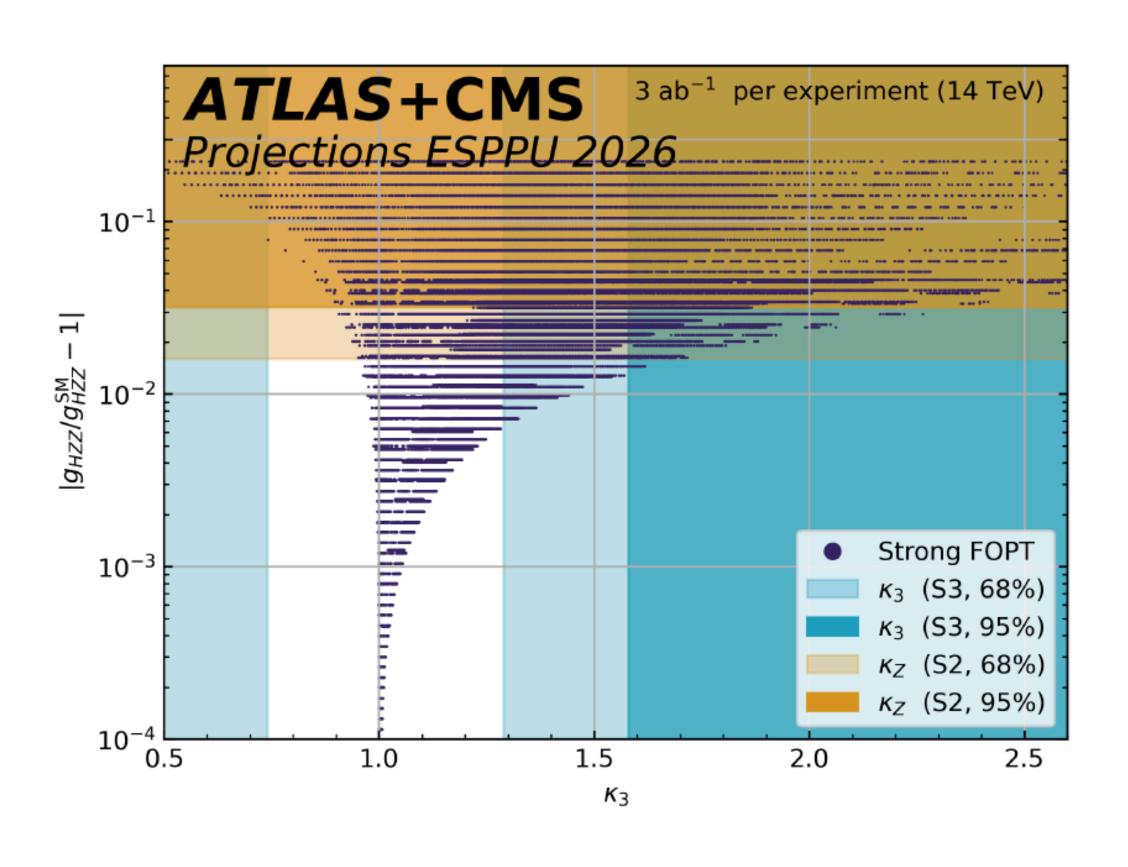


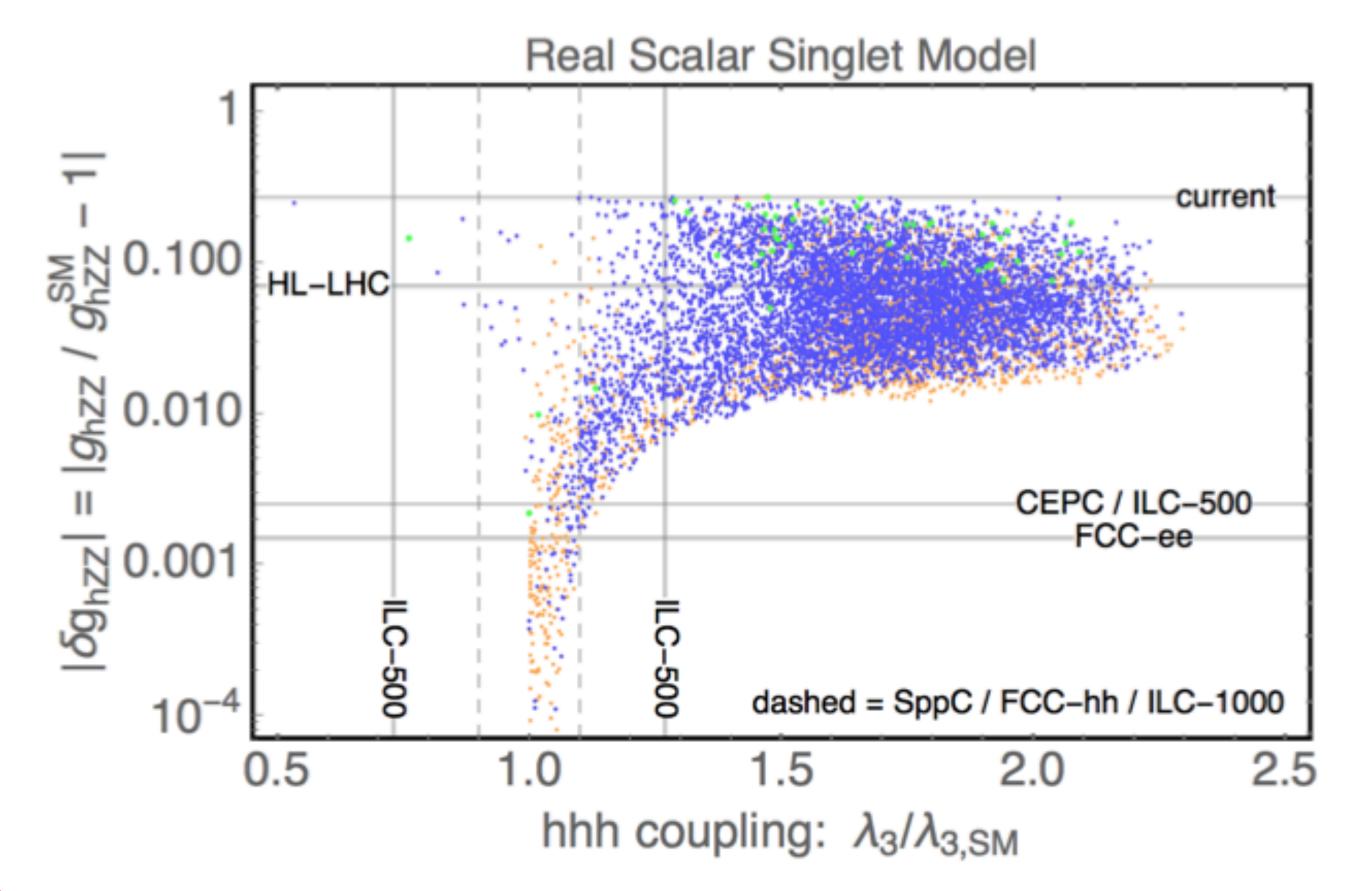




## Scalar Singlet FOPT exclusion by HZZ and k<sub>3</sub>

arXiv1608.06619

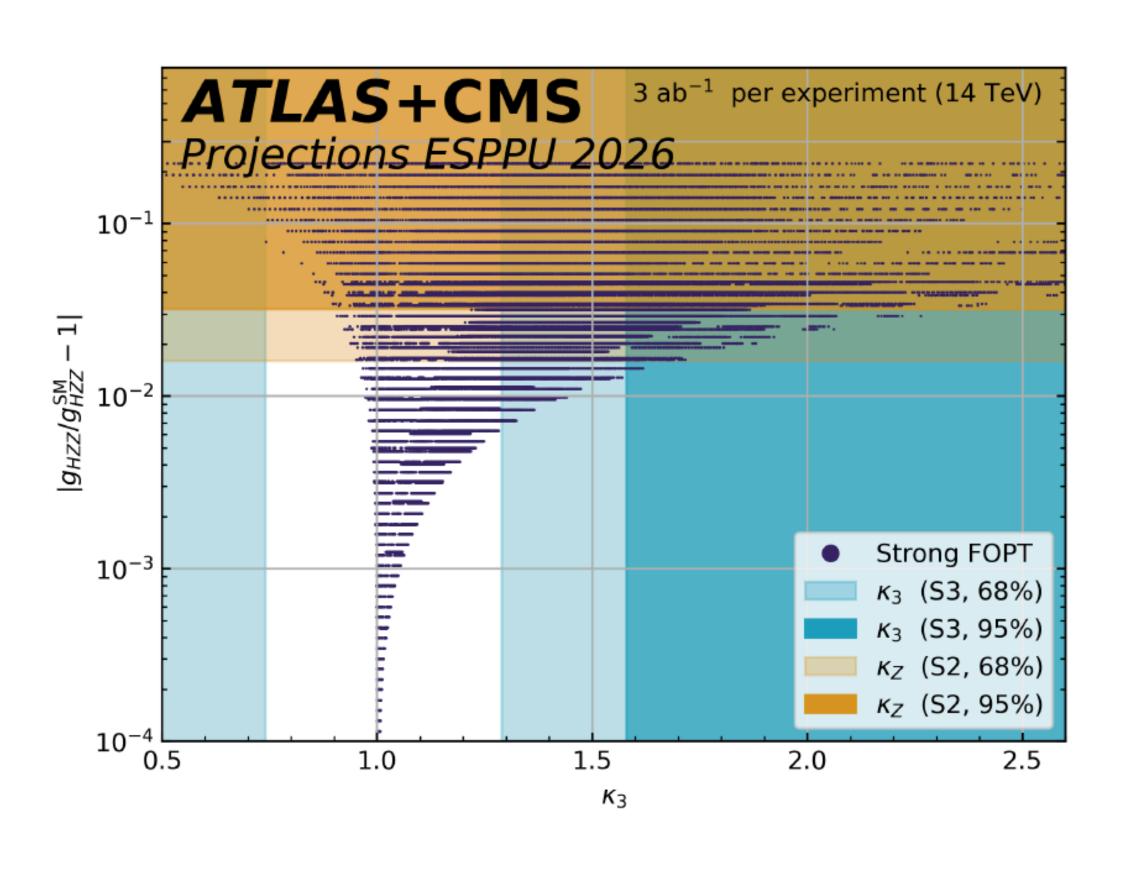


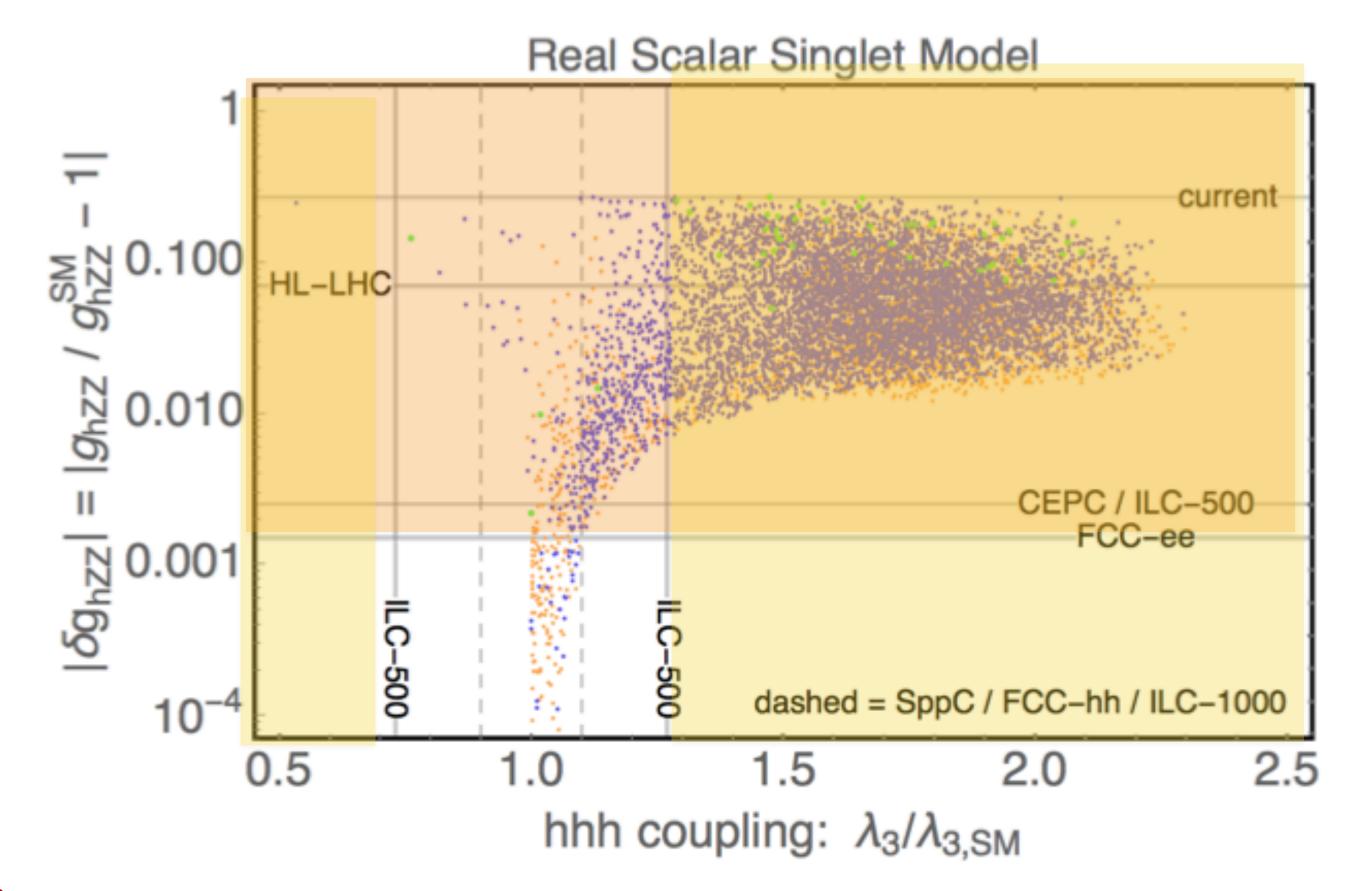


Global exclusion, without constraint from resonances that depend on specific values of model parameters

## Scalar Singlet FOPT exclusion by HZZ and k<sub>3</sub>

arXiv1608.06619

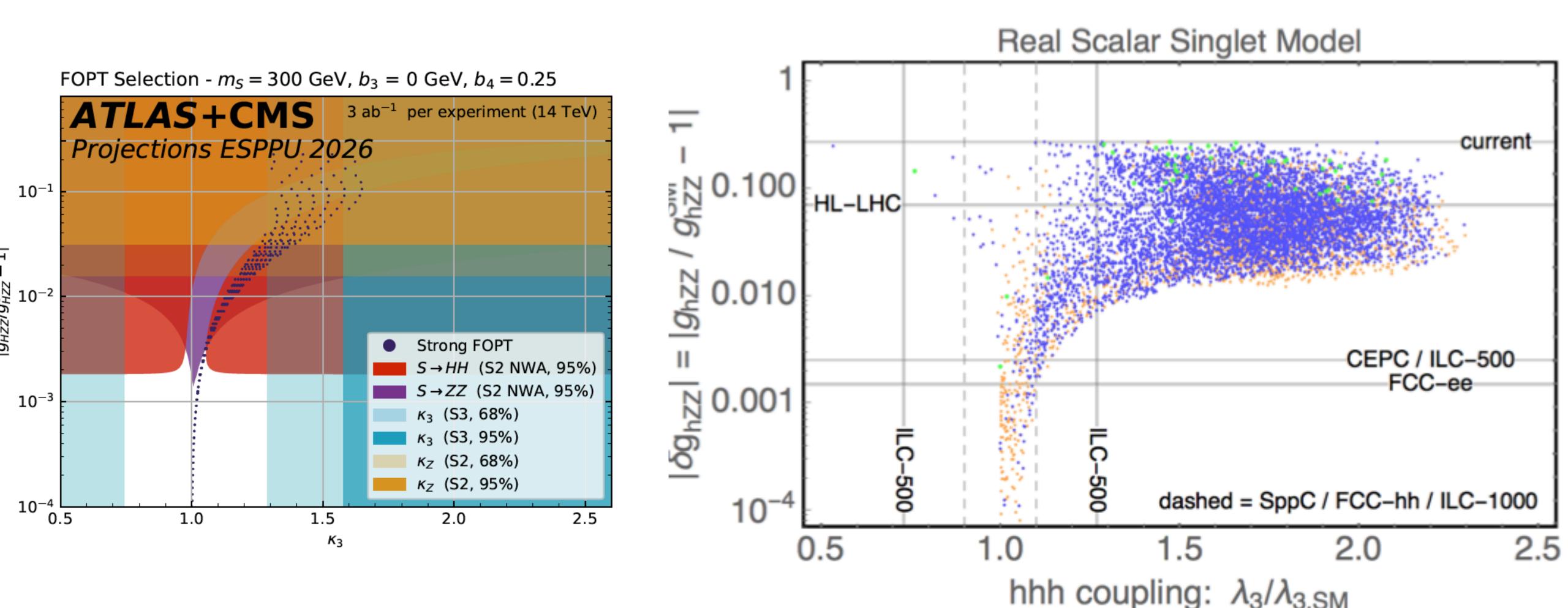




Global exclusion, without constraint from resonances that depend on specific values of model parameters

#### Scalar Singlet FOPT exclusion

arXiv1608.06619

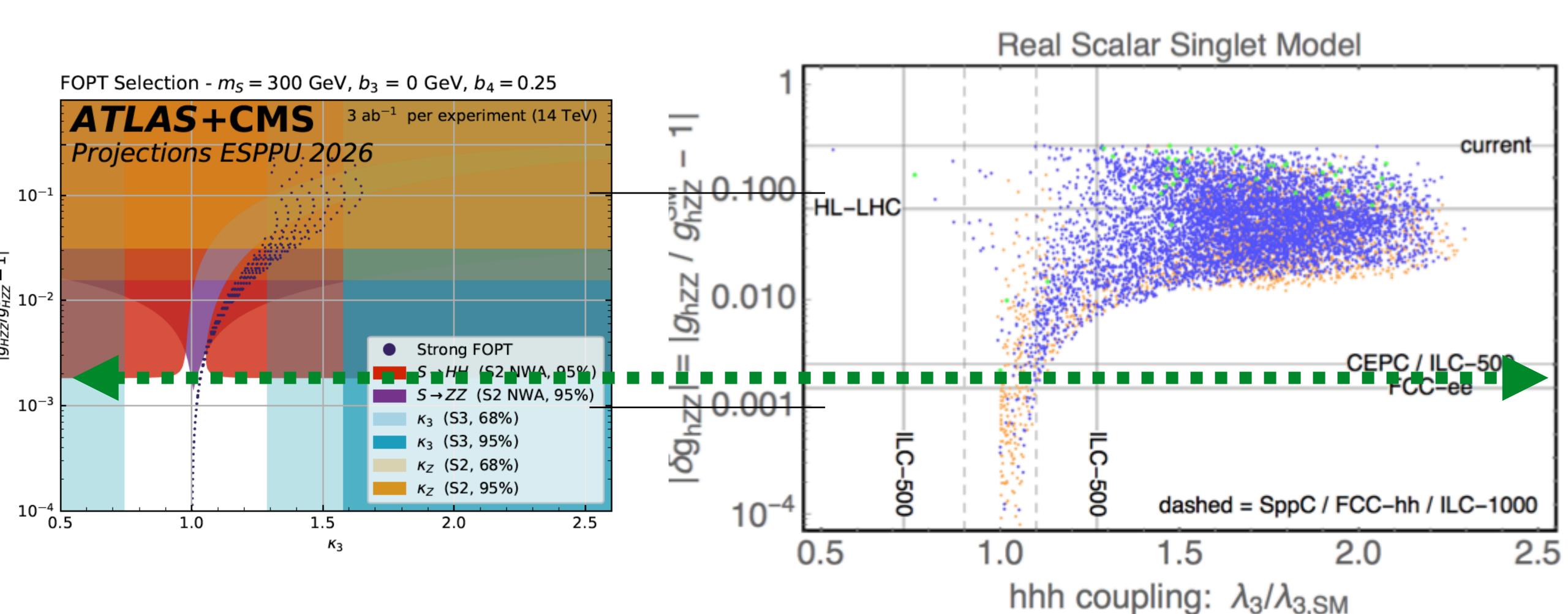


Exclusion in each slice shows also constraints from resonant searches

The exclusion is as strong as FCC-ee and ILC.

#### Scalar Singlet FOPT exclusion

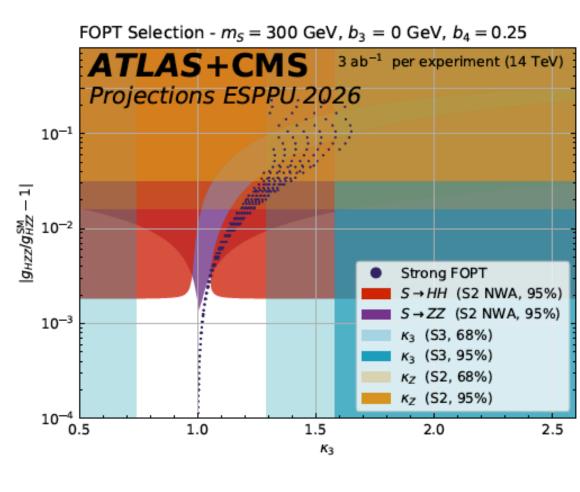
arXiv1608.06619

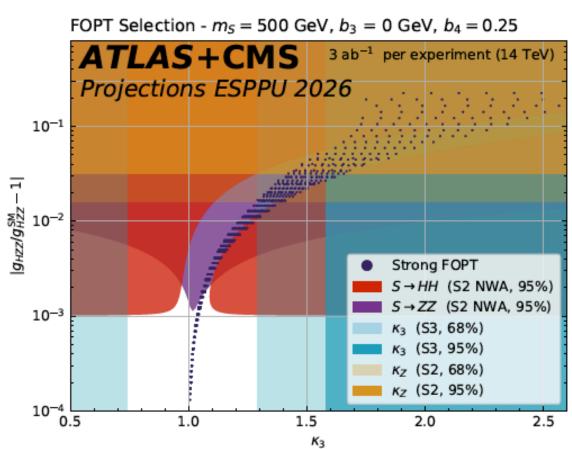


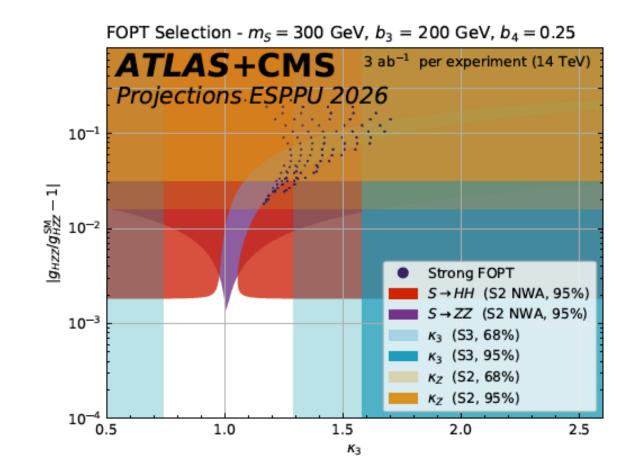
Exclusion in each slice shows also constraints from resonant searches
The exclusion ~ FCC-ee and ILC.

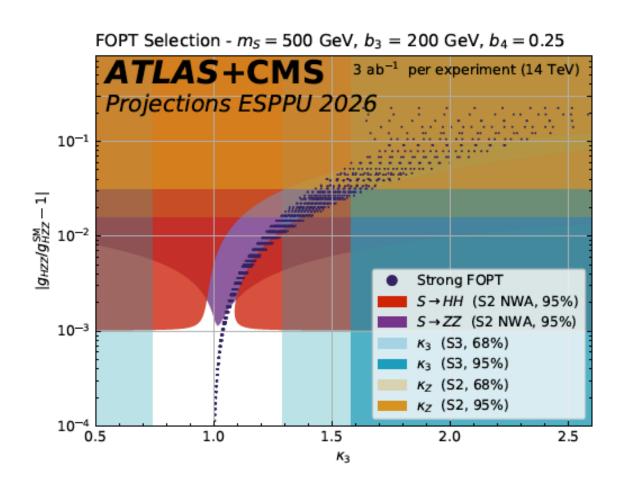
37

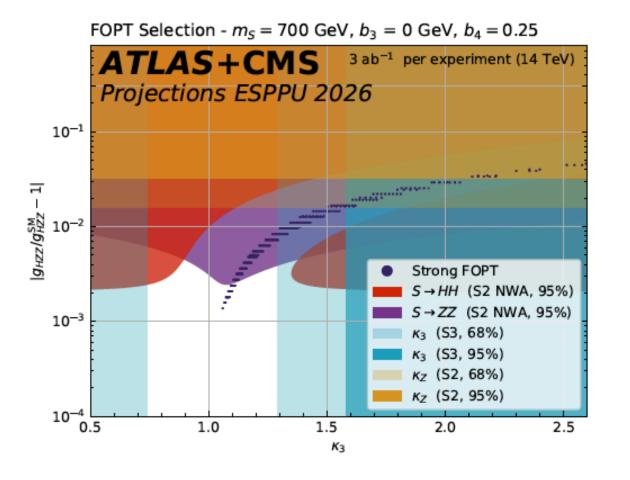
#### sample of different model parameters

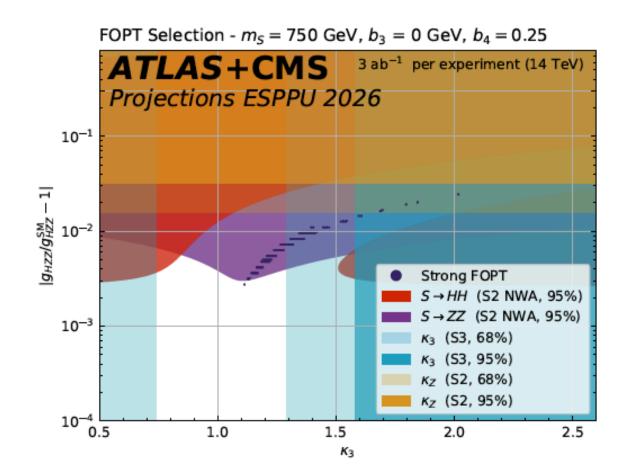


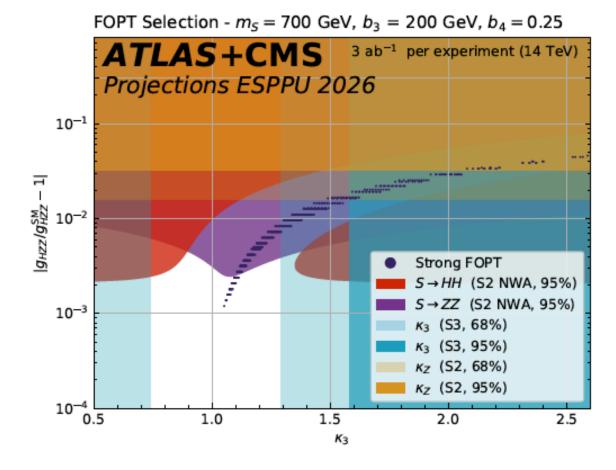


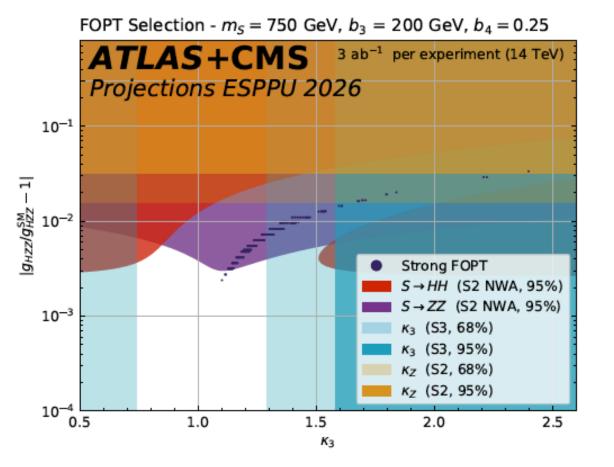








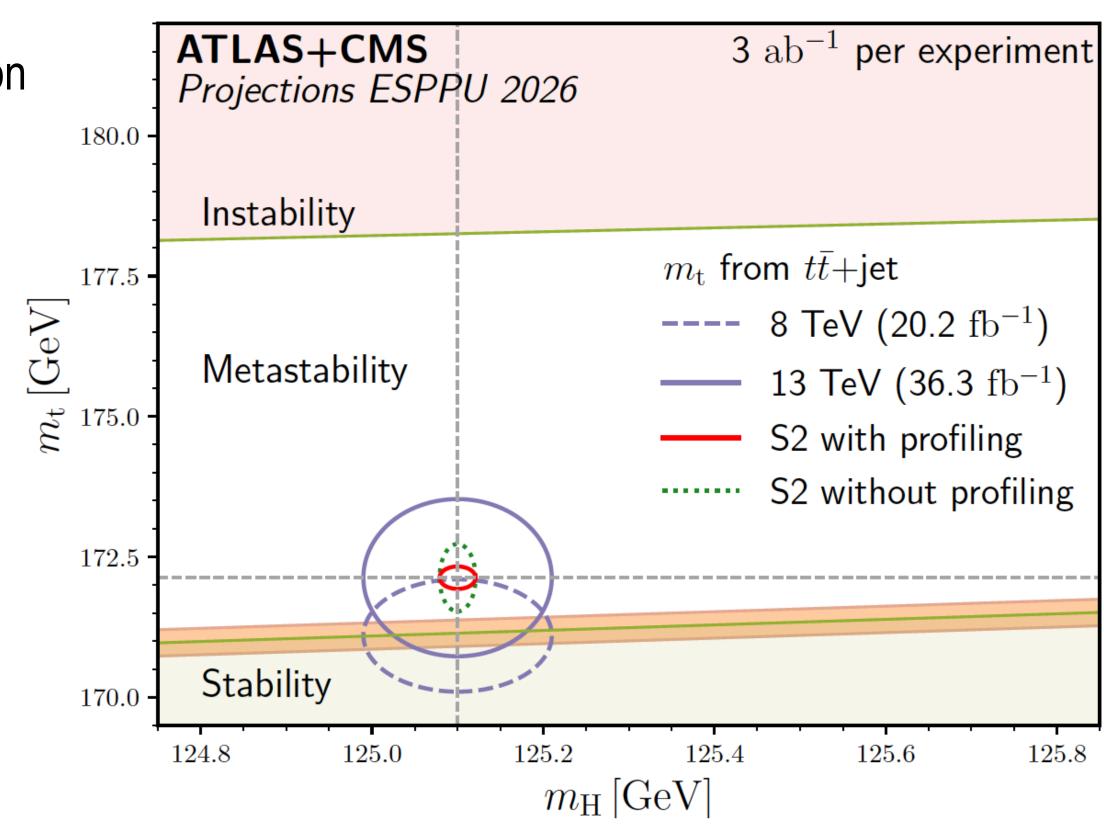




#### EW Vacuum Stability

In case the SM is valid up to the Plank scale we have indirect indication that our Universe is metastable. This plots provides important information about tensions in the SM.

if central value stays where it is now, we will have more than  $2\sigma$  exclusion of the Stability



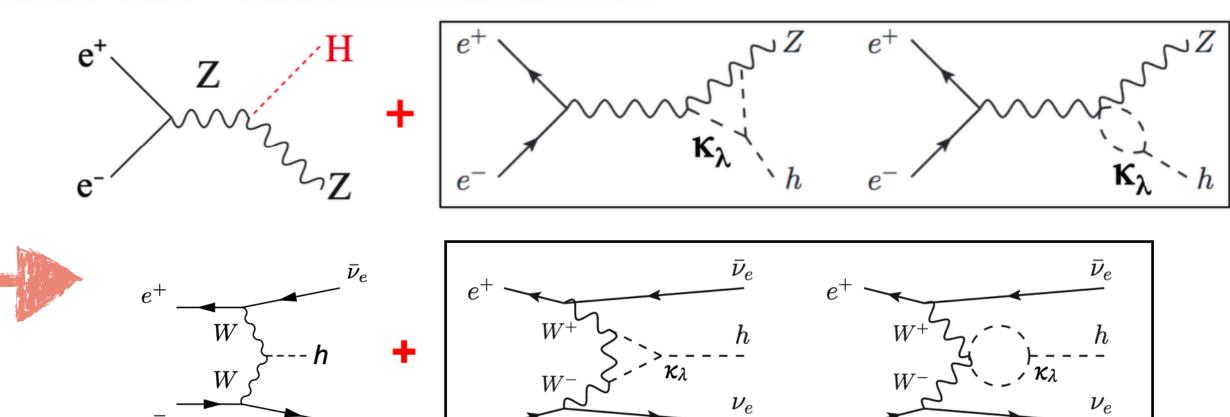
200 MeV precision from tt+jets 13 TeV by CMS (link) profiling

600 MeV precision from tt+1jets 8 Tev by ATLAS (link) no profiling

mHZZ=ATLAS+CMS 21 MeV

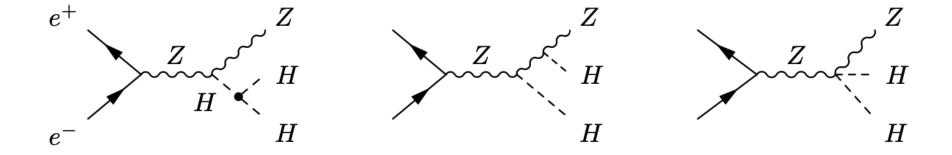
#### Higgs potential and $k_{\lambda}$ at e<sup>+</sup>e<sup>-</sup> Colliders

 At e+e- Colliders below di-Higgs threshold constraints on the Higgs potential come from indirect (or better at 1 loop) single Higgs production processes

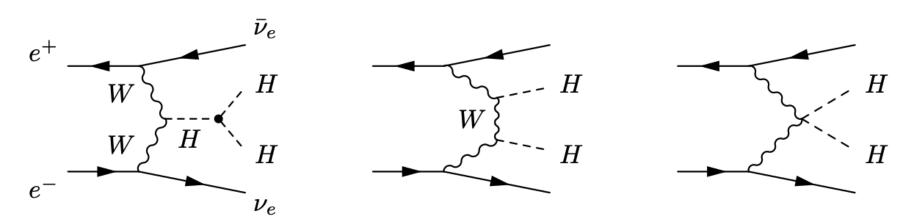


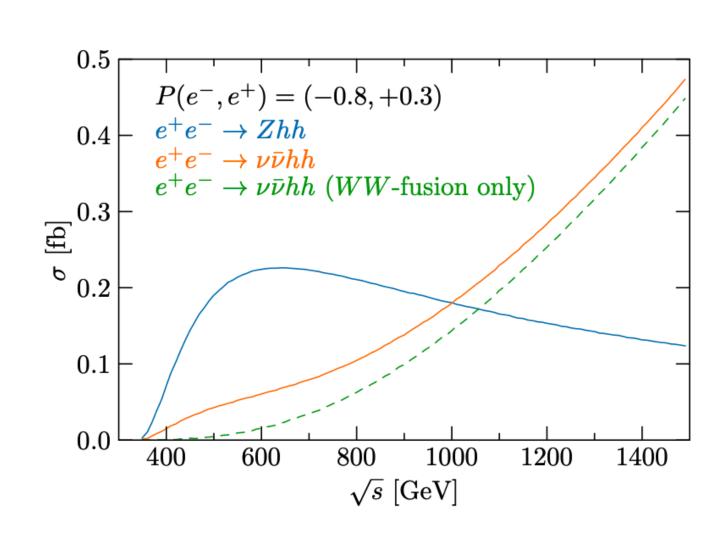
• At e+e-Colliders above di-Higgs threshold, direct di-Higgs processes

double Higgs-strahlung:  $e^+e^- \rightarrow ZHH$ 



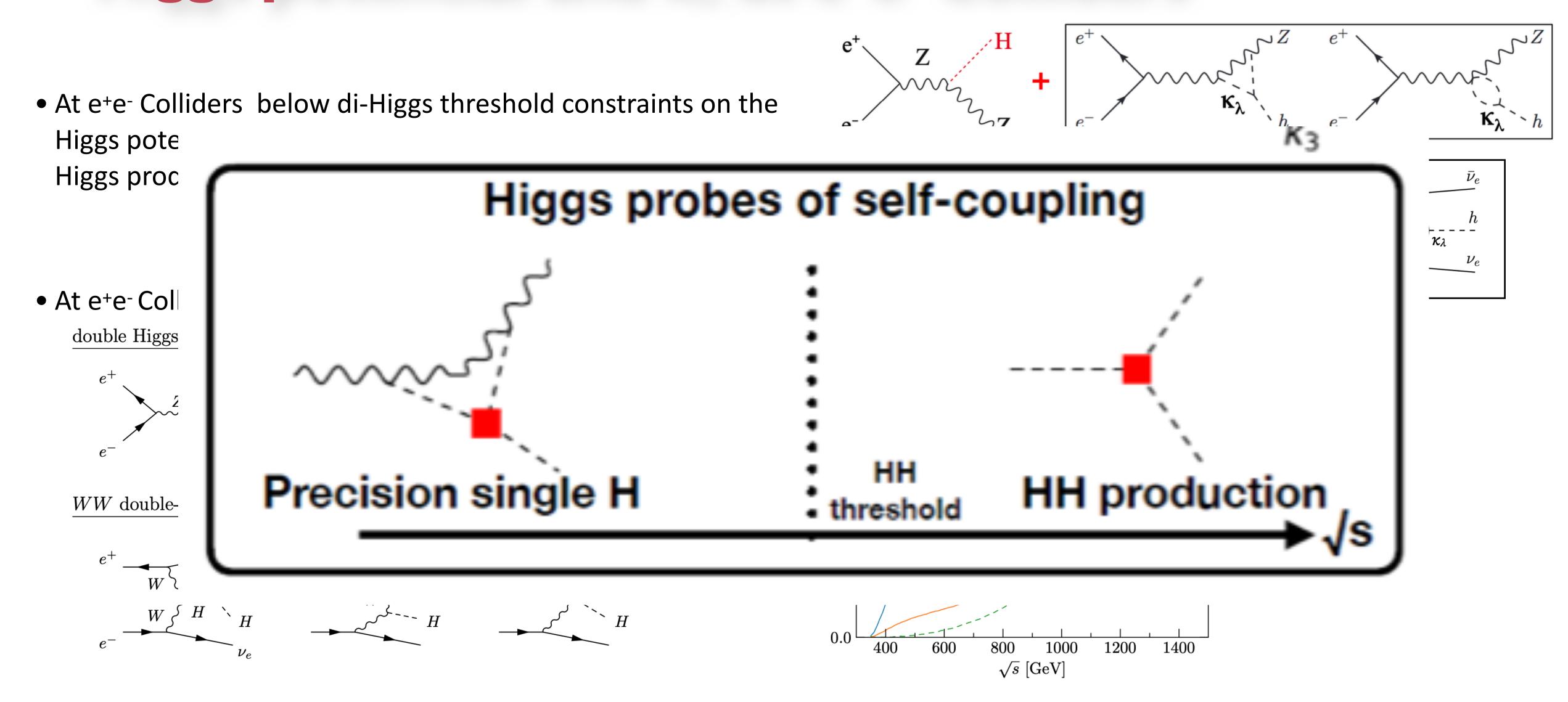
WW double-Higgs fusion:  $e^+e^- \rightarrow \bar{\nu}_e\nu_e HH$ 





• Indirect sensitivity to limited by: experimental uncertainties on considered observables, on input parameters, theoretical uncertainties due to unknown Higher Order corrections

#### Higgs potential and $k_{\lambda}$ at e<sup>+</sup>e<sup>-</sup> Colliders



• Indirect sensitivity to limited by: experimental uncertainties on considered observables, on input parameters, theoretical uncertainties due to unknown Higher Order corrections

#### Single&di-Higgs constraints on ka

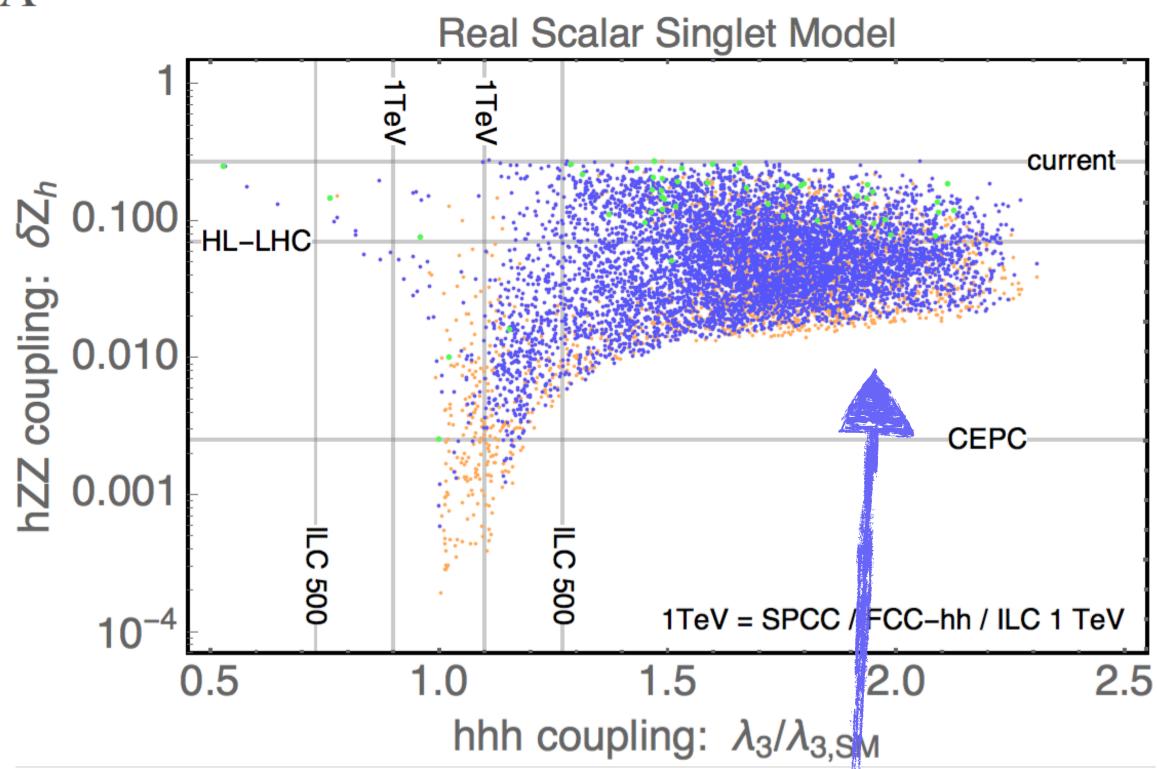


$$\frac{1}{\Lambda^2} (H^{\dagger} \partial H)^2$$
 Modifies H-Z coupling  $\Rightarrow \delta Zh$ 

- The HZ Coupling is better measured, but deviations in the Zh coupling ( $\delta$ Zh) are proportional to gz (the Z SM coupling)
- deviations in  $\lambda$  ( $\delta\lambda$ ) are not related to  $\lambda_{SM.-}$  There are Models iwhere  $\delta\lambda >> \delta Zh$ , notably the ones with a first order phase transition, with large self-interaction!

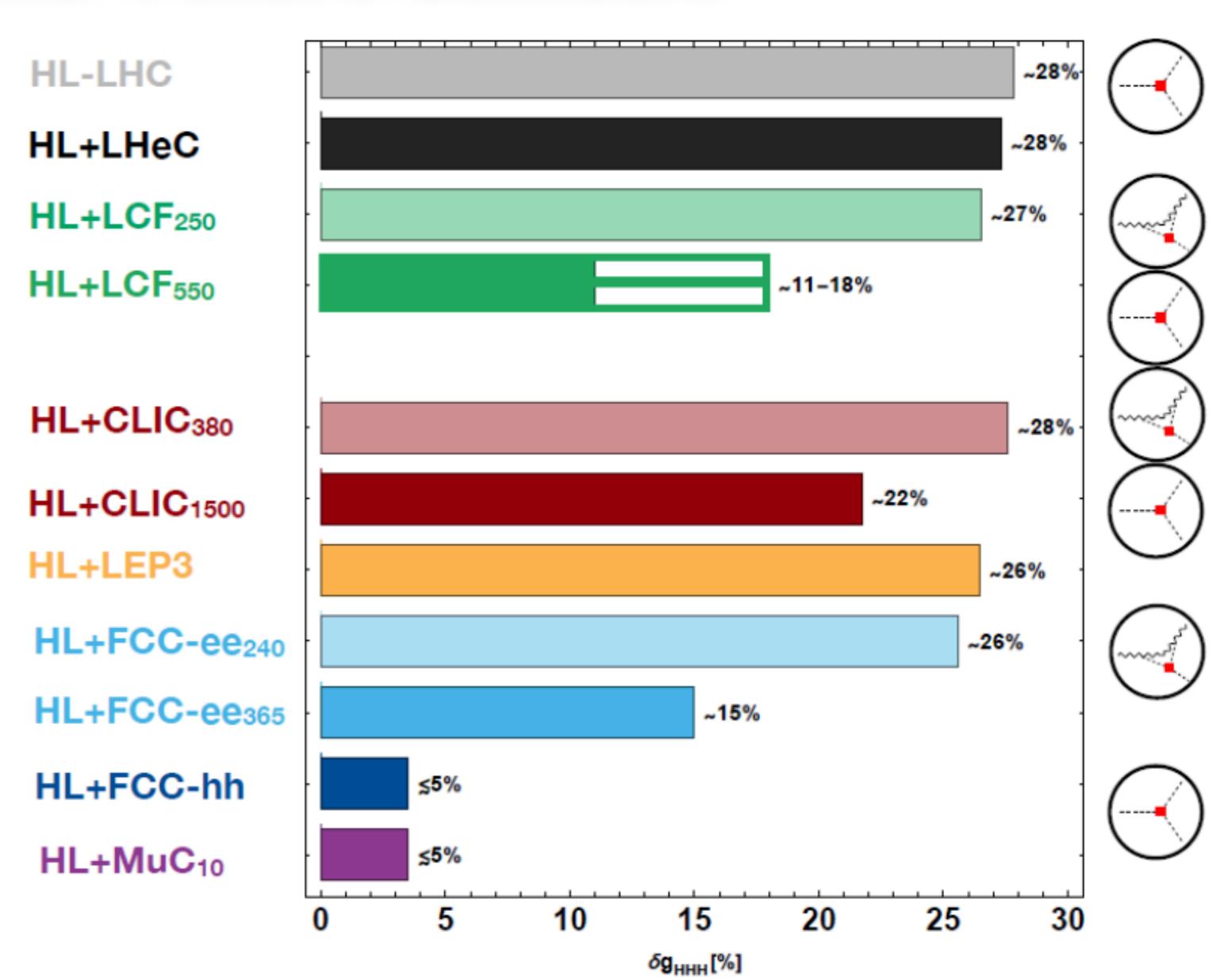
di-Higgs

$$\frac{1}{\Lambda^2}(H^{\dagger}H)^3$$
 modifies Higgs self coupling  $\delta\lambda$ 

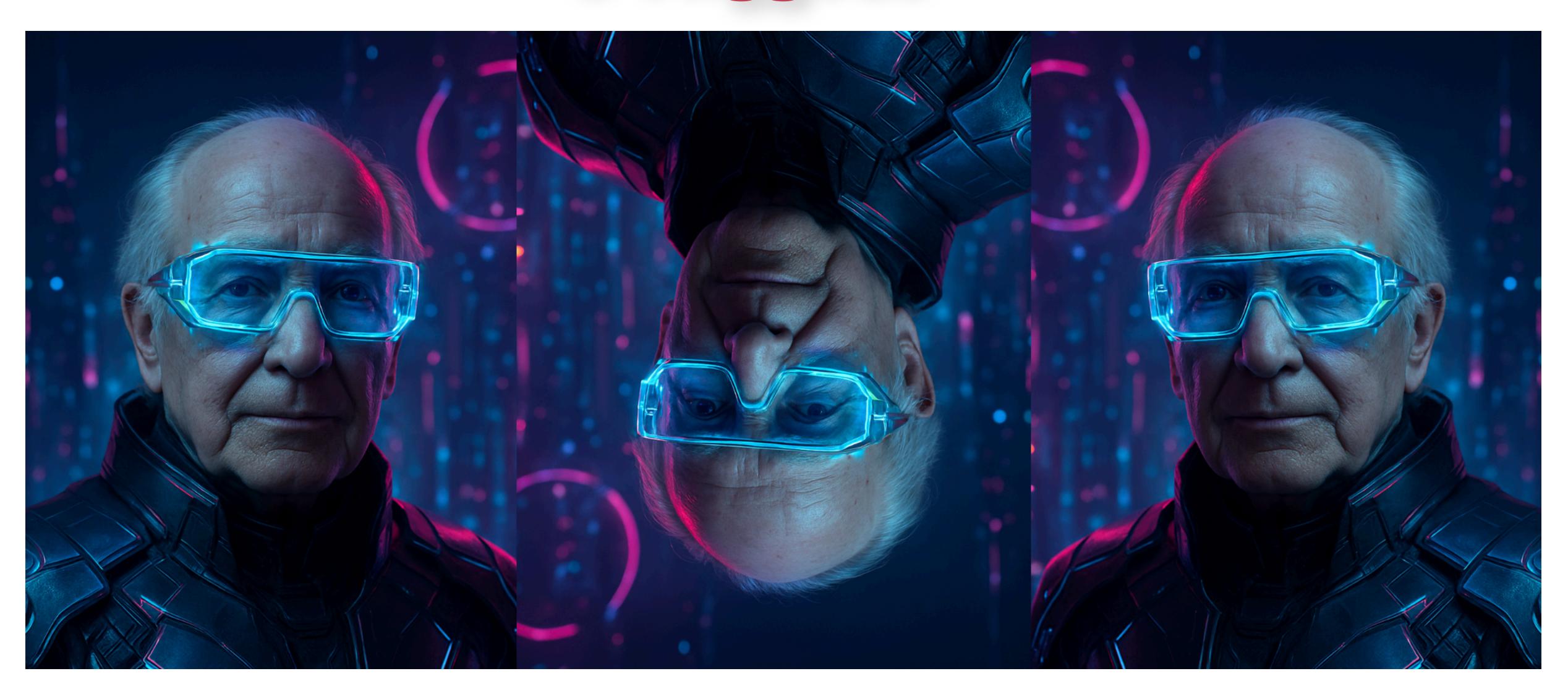


#### k<sub>λ</sub> and Future Colliders

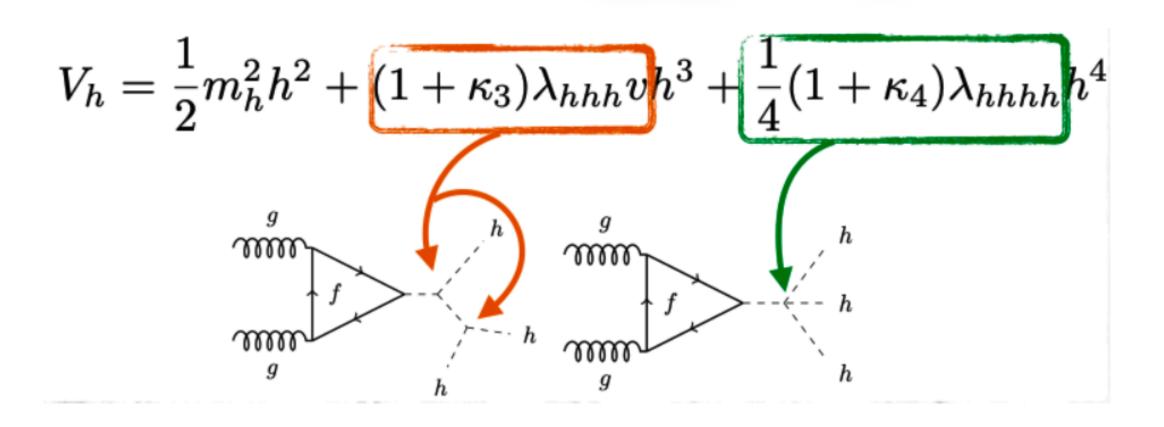
At e<sup>+</sup>e<sup>-</sup> Colliders below HH threshold come from indirect single Higgs constraints



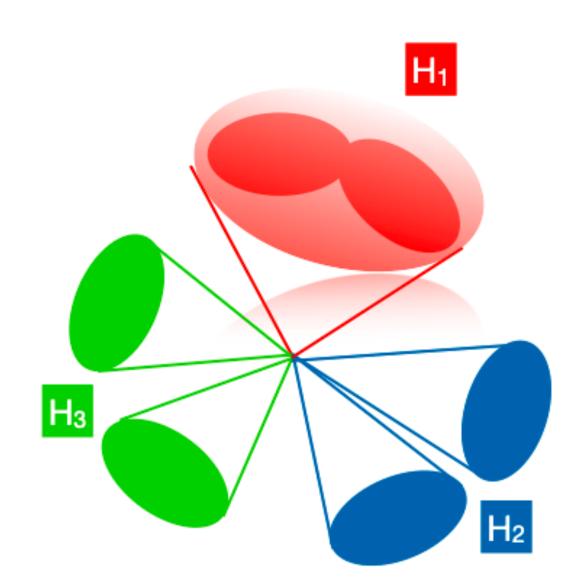
## 3 Higgses

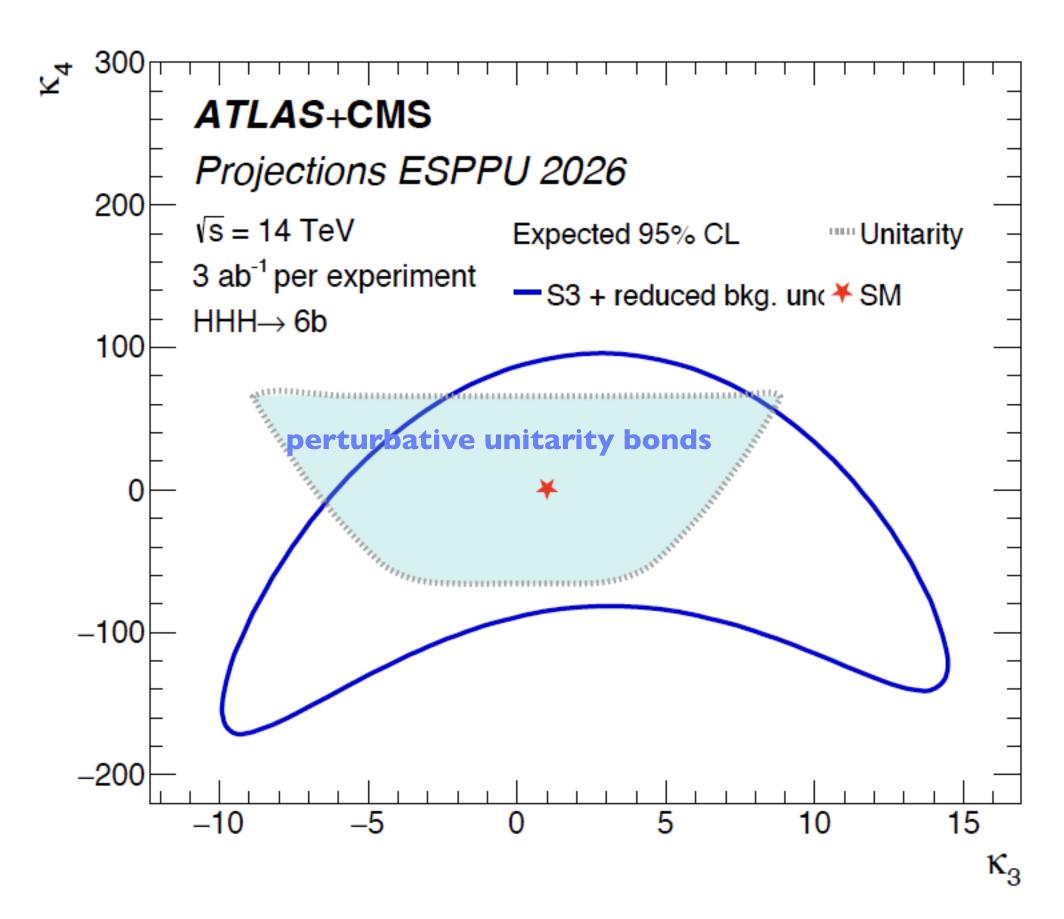


### Higgs quartic coupling



First time that the constraints are set on the quartic Higgs coupling.





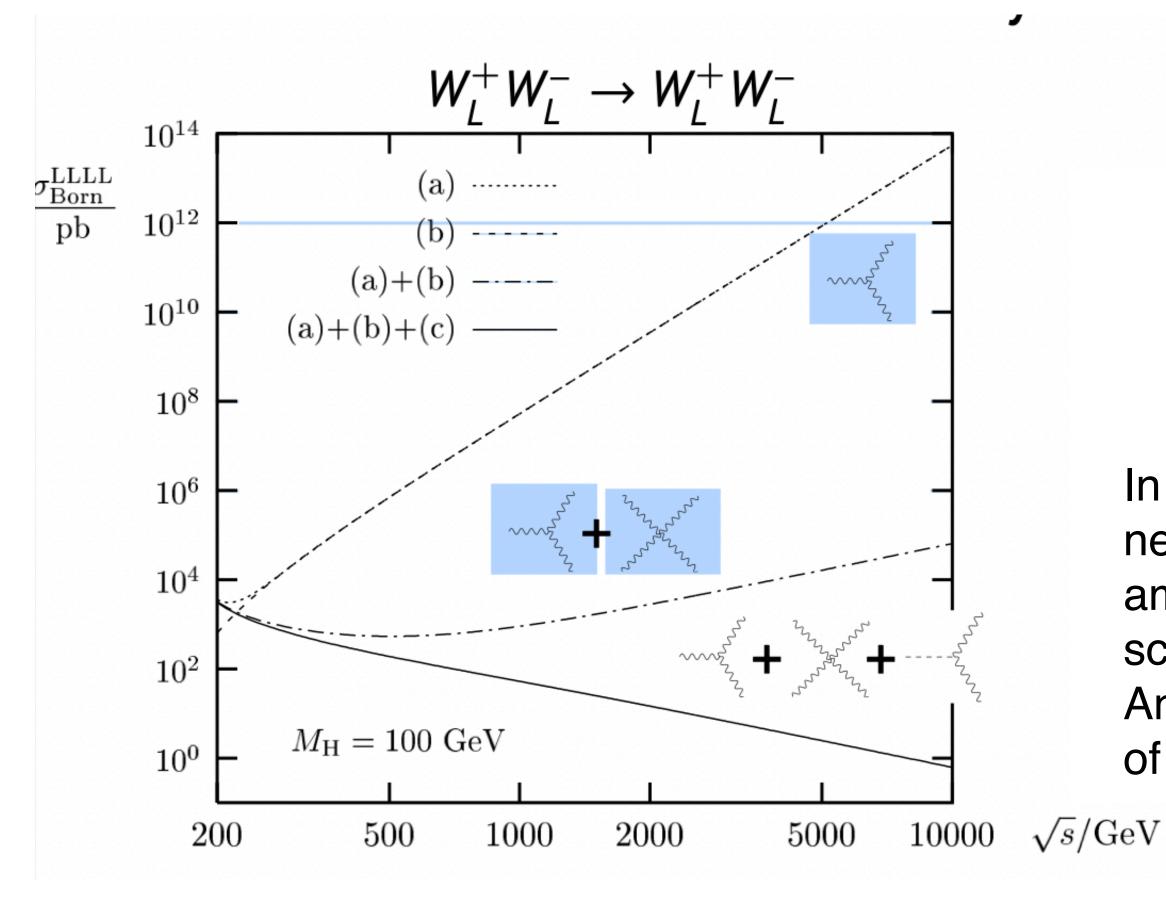
95% CL limit on cross-section < 109xSM More channels will be added in the future!

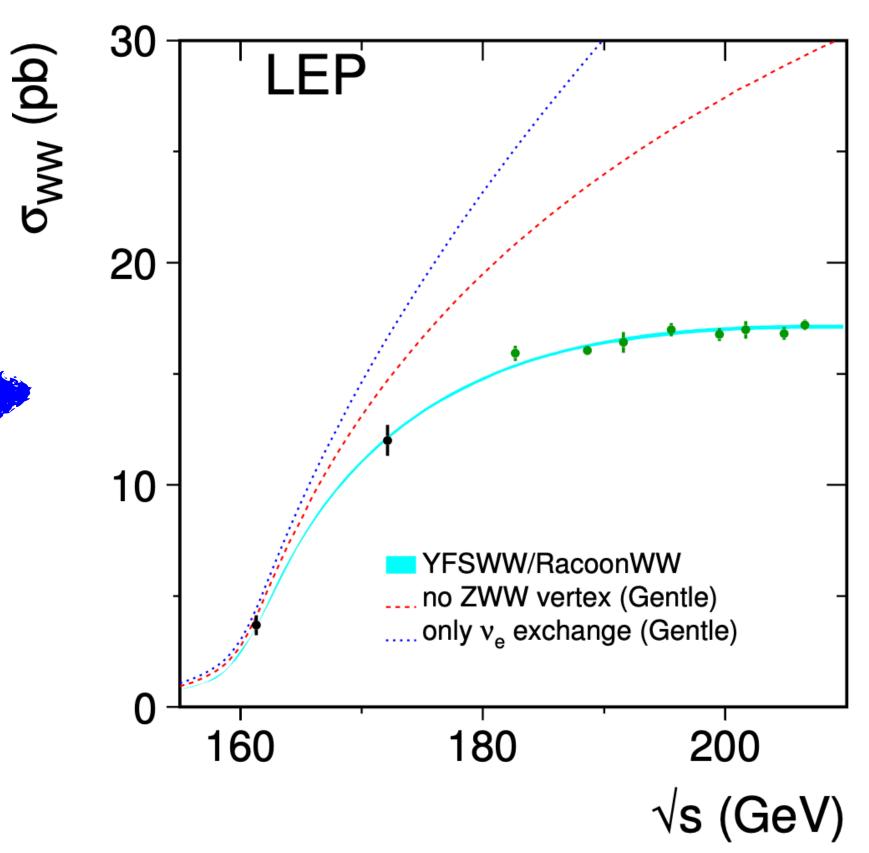
# Testing the Electroweak symmetry breaking via Vector Boson Scattering: Another approach

#### Electroweak symmetry breaking

Gauge-boson self interactions play a crucial role for the renormalisability of the electroweak theory

Large cancellations of divergences arising in individual diagrams are exact if couplings take the values of the SM

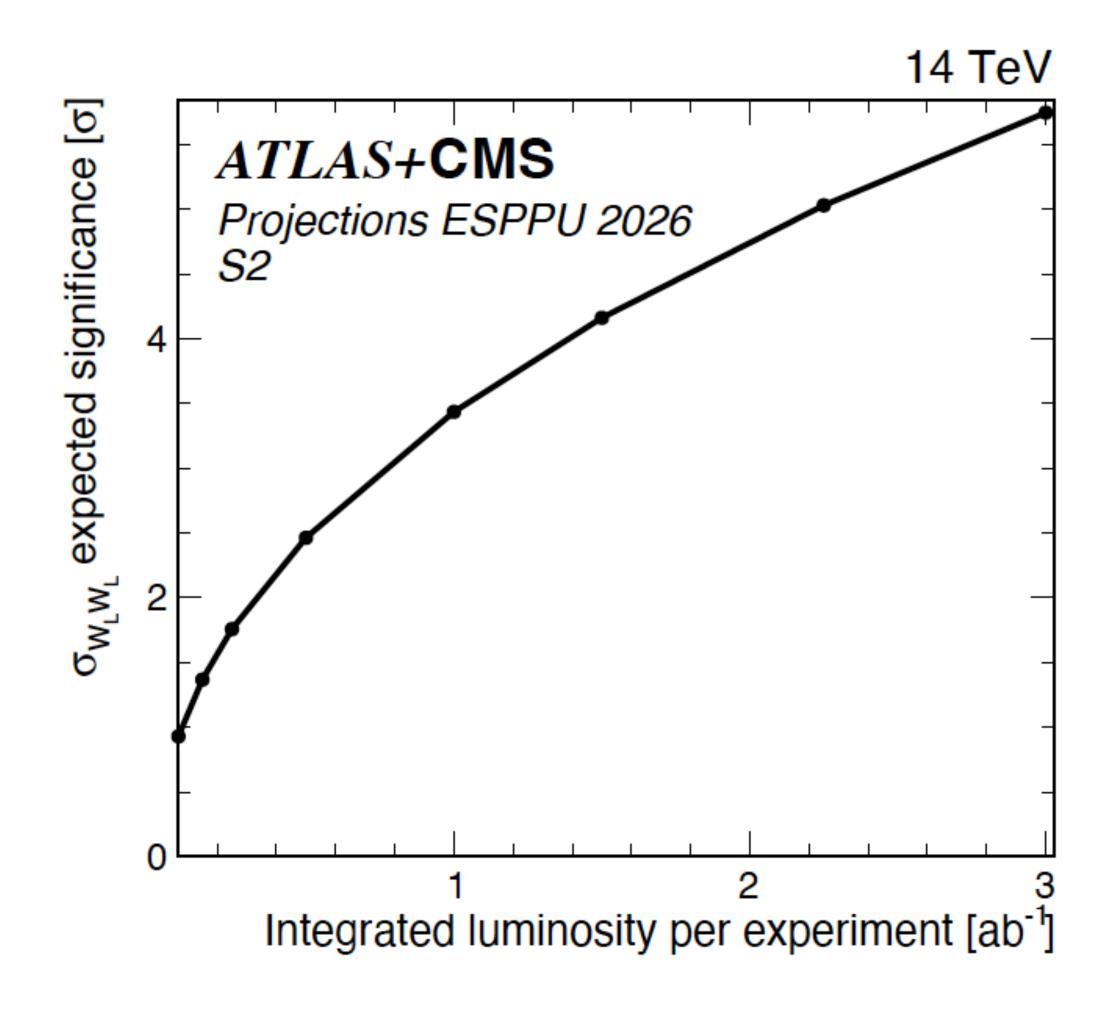




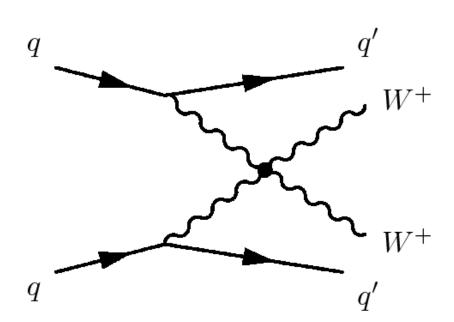
In vector boson scattering, the presence of the Higgs boson is needed to exactly cancel out the otherwise diverging scattering amplitudes at high energies and prevent unitarity violation at the TeV scale.

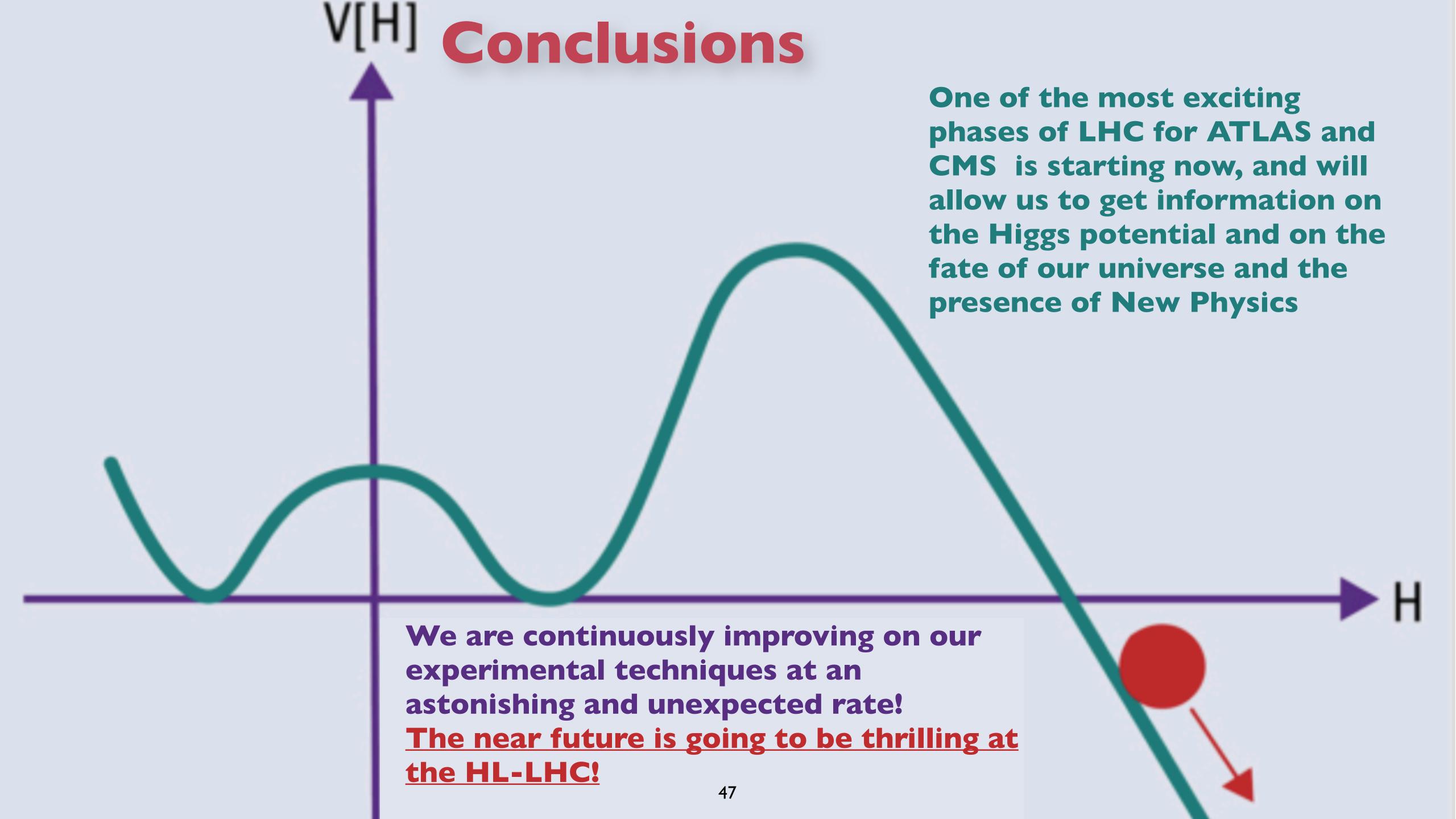
Any significant deviation from the predicted high-energy behaviour of vector boson scattering would point to new phenomena.

#### di-boson scattering: a complementary path

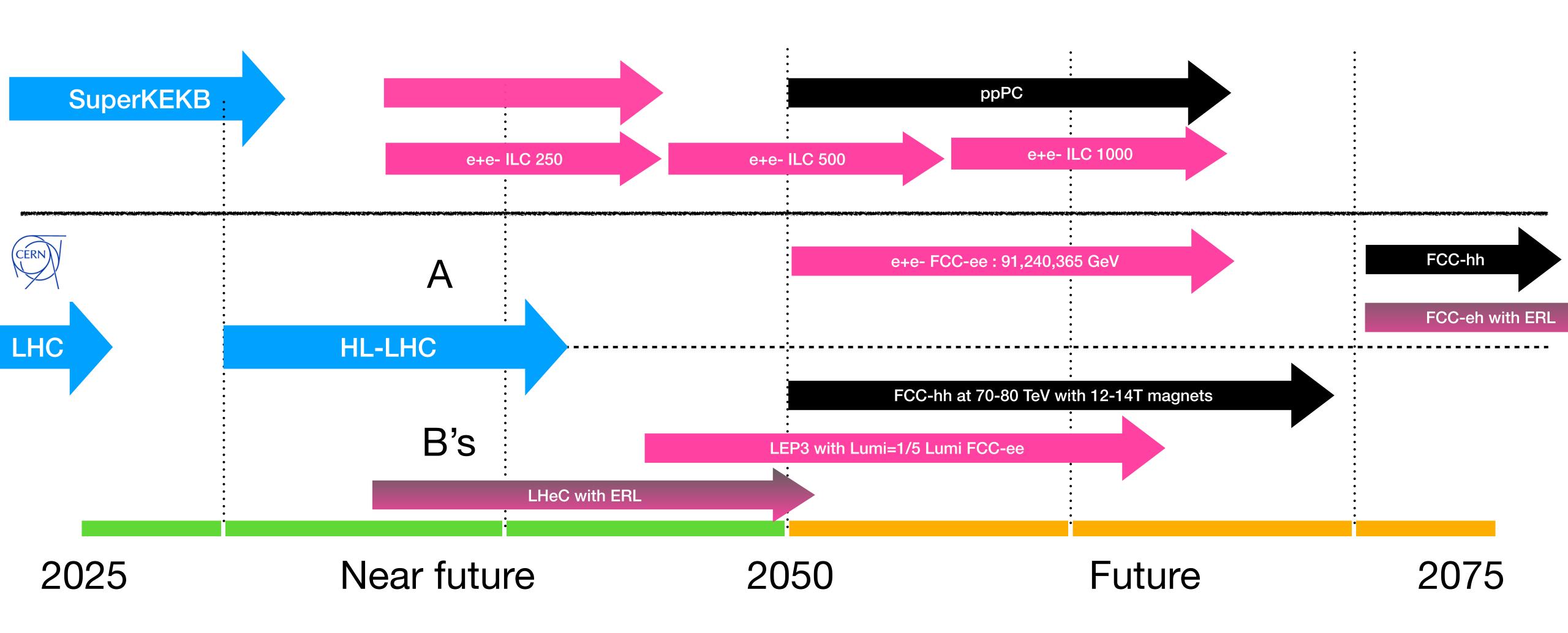


Observation of longitudinally polarised vector boson scattering WLWL process with better than 20% precision





## Going back from where we started.. are our opinions changed?

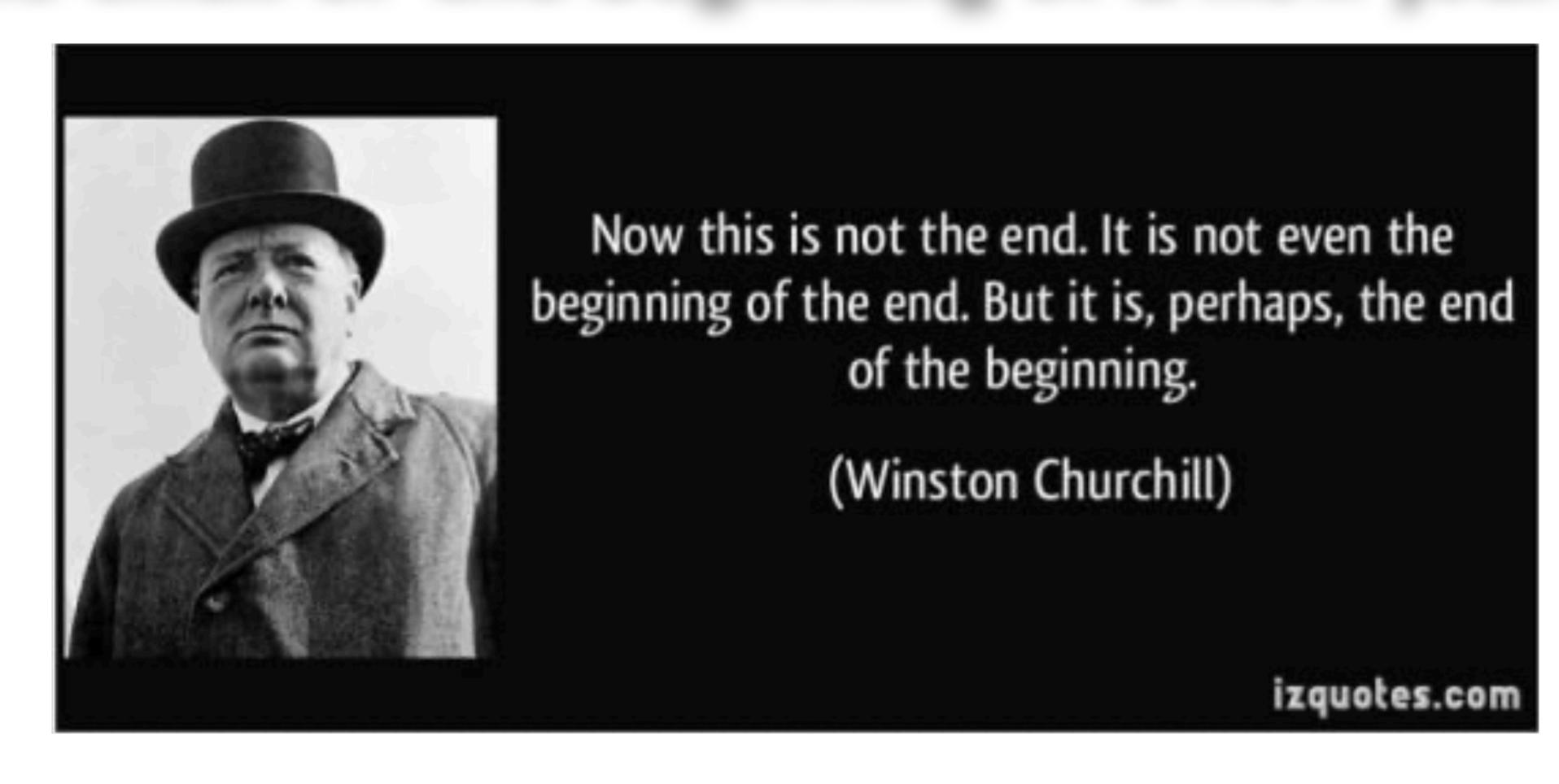


#### cost estimates

Fcc-ee	+Fcc-hh 35	<b>BCHF</b>
Fcc-hh	27 BCHF	

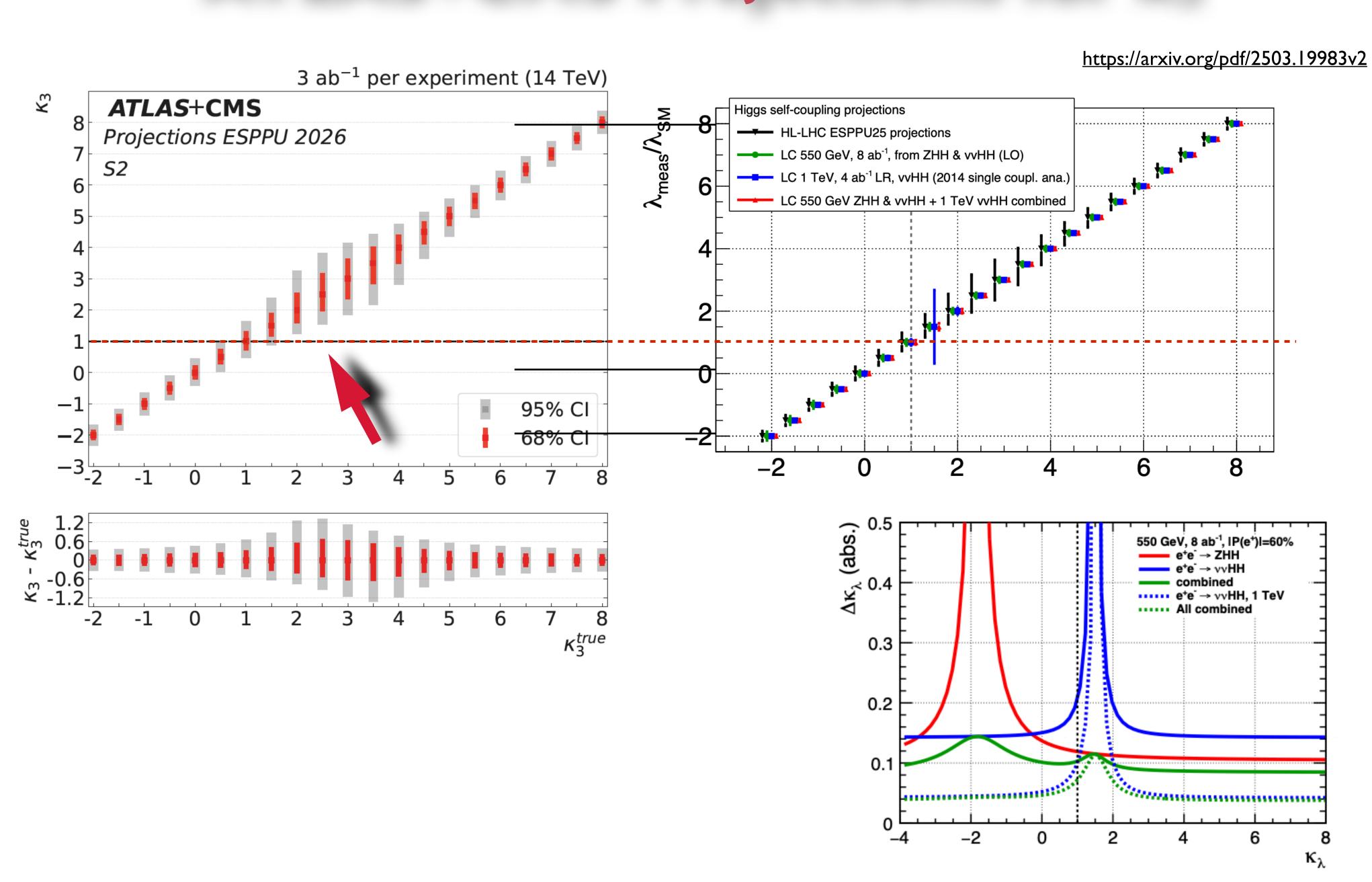
project	scenario / energy / details	cost	roject	scenario	cost
FCC-ee	Z, WW, ZH, + 4 exp. CERN part	14.0 BCHF	.HeC	9 km new tunnel, 50 GeV e <sup>-</sup>	2 BCHF
	+ 4 exp. non CERN part	+1.3 BCHF	Muon	3.2 TeV at CERN	12.0 BCHF
	+ upgrade to ttbar	+1.3 BCHF	Collider	J.Z IEV at CLINI	12.0 DCH
FCC-hh	85 TeV pp, as second stage	+19.1 BCHF		7.6 TeV at CERN	17.0 BCHF
FCC-hh	85 TeV pp, standalone	27.3 BCHF		10 TeV green field	19 BCHF
LCF	250 GeV	8.5 BCHF	HALHF	e <sup>-</sup> 375 GeV, e <sup>+</sup> 41.7	3.8 BCHF
	550 GeV	+5.5 BCHF		GeV (250 GeV c.m.) 380 GeV	4.9 BCHF
CLIC	380 GeV	7.3 BCHF		550 GeV	6.3 BCHF
	1500 GeV	+7,1 BCHF			
I ED2	74 of 220 CoV + 2 over CEDN	3.9 BCHF	ALIVE	p-driven, 10 TeV	?
LEP3	ZH at 230 GeV + 2 exp CERN part	ALE	ALEGRO	e⁺e⁻ & γγ, ≥10 TeV	?
	+ 2 exp non CERN part	+0.9 (0.3) BCHF			

#### The end.. or the beginning of a new journey



#### BACK-UP

#### ATLAS+CMS Projections for k<sub>3</sub>



#### Ratios of K's reduce theory systematics

