

Extra twisted connected sums and their ν -invariants

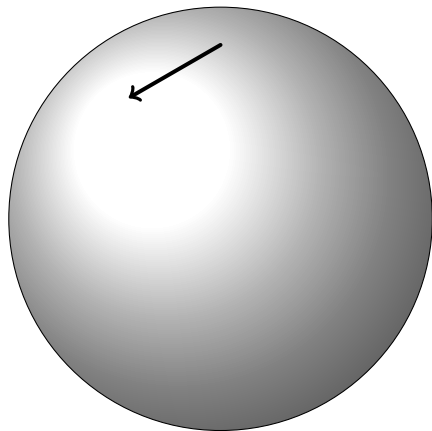
Sebastian Goette

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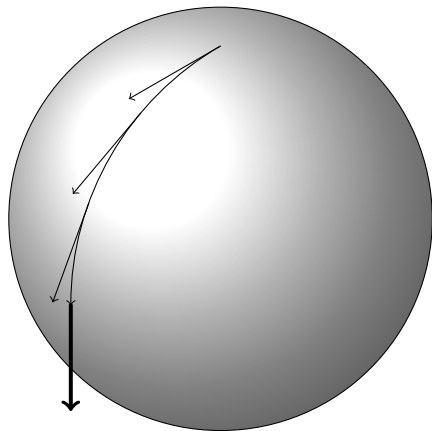
Geometry, Dynamics, and Zeta Functions
Nijmegen, 17–19. 6. 2026

- ▶ G_2 -Geometry
Intro, properties, questions
- ▶ The ν -invariant
Differential topology, definition of ν , properties, first examples
- ▶ Extra twisted connected sums
Construction, properties, problems
- ▶ Computation of the ν -invariant
Computations with η -invariants, examples, questions

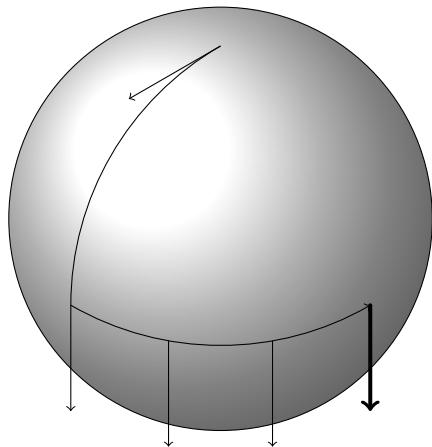
Consider parallel translation along a spherical triangle



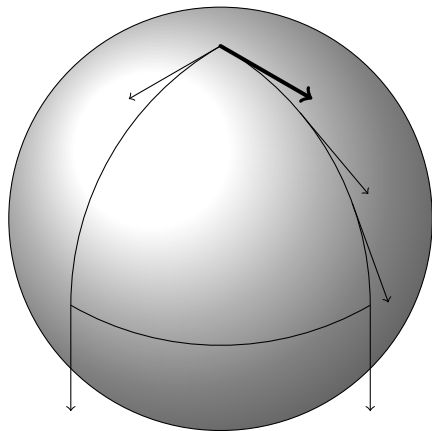
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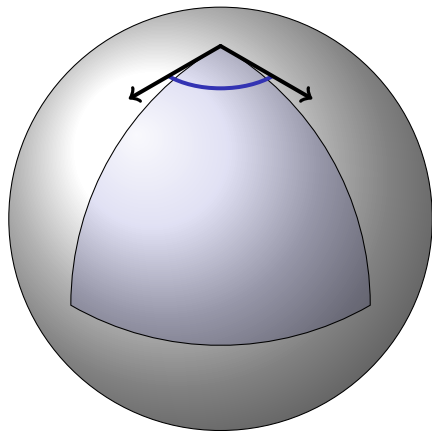


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([Gauß-Bonnet theorem](#))

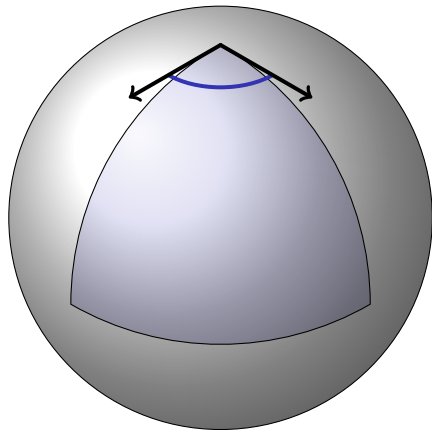


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$$SO(2) \cong \text{Hol}(S^2, g^{\text{rd}}) \subset \text{Aut}(T_p S^2)$$



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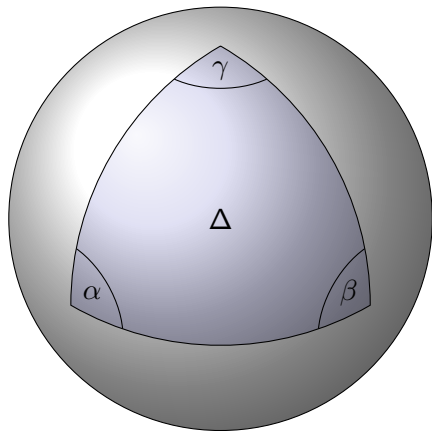
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Related: the spherical area formula

$$A(\Delta) = \alpha + \beta + \gamma - \pi .$$



Theorem (Berger)

The only possible holonomy groups of complete, simply connected Riemannian manifolds that are neither a product nor a symmetric space are

Holonomy group	dim	ric	Structure	Parallel spinors	Name
$SO(n)$	n				generic case
$U(k)$	$2k$		J		Kähler
$SU(k)$	$2k$	0	J, Ω	2	Calabi-Yau
$Sp(\ell) \cdot Sp(1)$	4ℓ	const	$\langle I, J, K \rangle$		Quat. Kähler
$Sp(\ell)$	4ℓ	0	I, J, K, Ω	$\ell + 1$	hyper Kähler
G_2	7	0	$\varphi \in \Omega^3$	1	exceptional
$Spin(7)$	8	0	$\psi \in \Omega^4$	1	exceptional

Mathematical motivation

- ▶ Only special holonomy group for odd dimensional manifolds
- ▶ Only G_2 and $\text{Spin}(7)$ holonomy have no direct relation to algebraic geometry

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Physical motivation

- ▶ In string theory, spacetime takes the form $\mathbb{R}^{3,1} \times V$, where V is Calabi-Yau
- ▶ In **M-theory**, spacetime takes the form $\mathbb{R}^{3,1} \times M$, where M is a G_2 -manifold
- ▶ Possible relations to other physical theories

Hence, many fruitful interactions possible

G_2 is simply connected \rightsquigarrow G_2 -manifolds are spin

The stabiliser of a nonzero spinor in $\text{Spin}(7)$ is G_2

A Riemannian 7-manifold (M, g) has $\text{Hol}(M, g) \subset G_2$

if and only if it is spin and there exists a nonzero **parallel spinor**

Let M be a simply connected compact oriented spin 7-manifold and define

$$\mathcal{X} = \{ g \mid \text{Hol}(M, g) \cong G_2 \}$$

Let $\mathcal{D} \subset \text{Diff}(M)$ be the connected component of id_M .

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It can be disconnected [Crowley-G-Nordström]
and contain nontrivial 2-spheres [G-Hertl]

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- ▶ Construct G_2 -metrics with prescribed singularities
Certain singularities allow massless chiral Fermions to appear in M -theory

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Idea. Use nowhere vanishing spinors to describe and distinguish G_2 -structures

Let σ_0, σ_1 be two nowhere vanishing spinors. Extend to $\bar{\sigma} \in \Gamma(\mathcal{S}^+(M \times [0, 1]))$
A generic $\bar{\sigma}$ will have nondegenerate isolated zeros because

$$\text{rk } \mathcal{S}^+(M \times [0, 1]) = 8 = \dim(M \times [0, 1])$$

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Orient $S^+(M \times [0, 1])$ and count with signs

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Let $F: M \rightarrow M$ be a spin diffeomorphism, then

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Can we write $\Delta\nu(M; \sigma_0, \sigma_1) = \nu(M, \sigma_0) - \nu(M, \sigma_1) \in \mathbb{Z}/48$?

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Problem. Given M , how to determine W with $M = \partial W$?

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- ▶ D_M —spin Dirac operator on $\Gamma(SM)$
- ▶ B_M —odd signature operator $*d \pm d*$ on $\Omega^{\text{ev}}(M)$
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Theorem (Crowley-G-Nordström)

$$\nu(M, \sigma) = 2 \int_M \sigma^* \psi(\nabla^{SM}, g^{SM}) - 24(\eta + h)(D_M) + 3\eta(B_M) \in \mathbb{Z}/48$$

Proof.

Use $2e(\nabla^{S^+W}) = e(\nabla) + 48\hat{A}(\nabla)^{[8]} - 3L(\nabla)^{[8]} \in \Omega^8(W)$ □

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- ▶ $\bar{\nu}(M, g) = 0$ if M admits an orientation reversing isometry

What about the known examples by Joyce and Kovalev?

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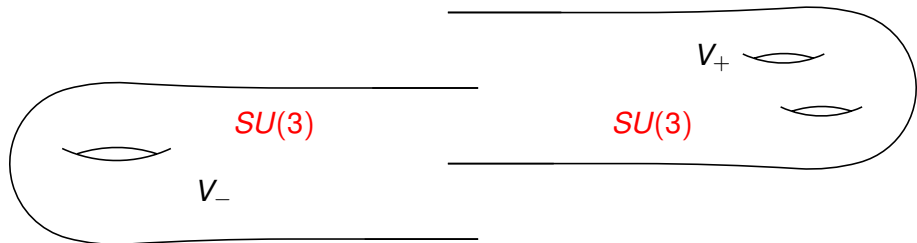
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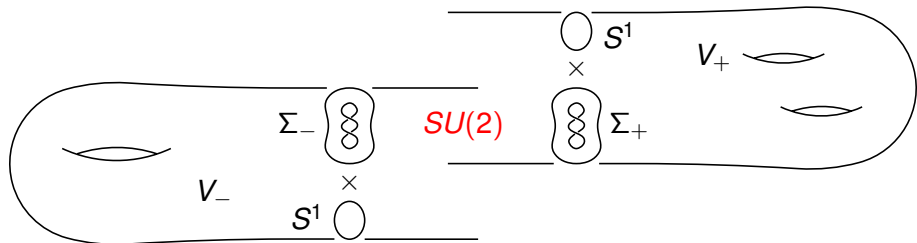
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the G_2 -moduli space \mathcal{M} has several connected components
—even with the same underlying G_2 -structure

Twisted connected sums à la Kovalev and Corti-Haskins-Nordström-Pacini



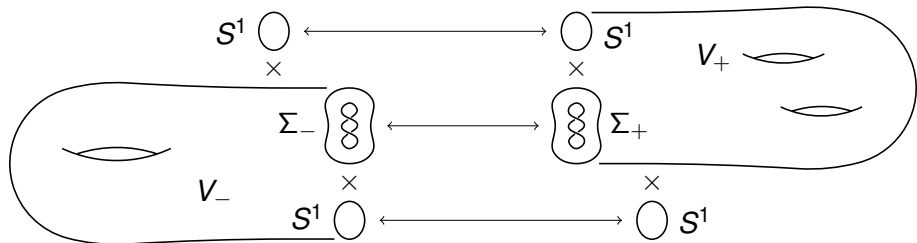
Let V_+ , V_- be **asymptotically cylindrical Calabi-Yau threefolds**

Twisted connected sums à la Kovalev and Corti-Haskins-Nordström-Pacini



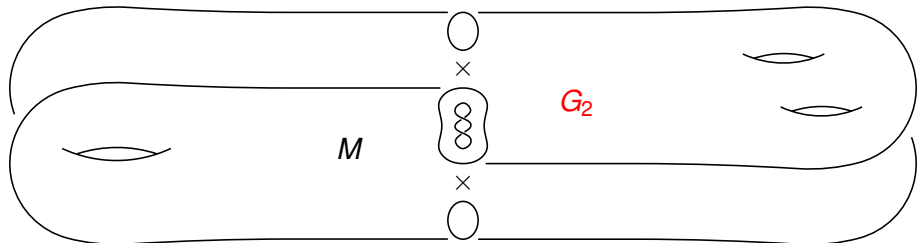
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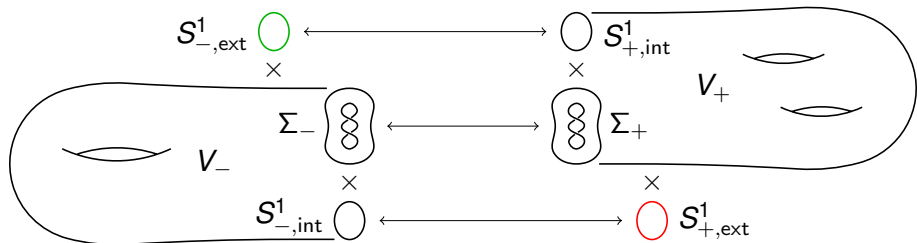


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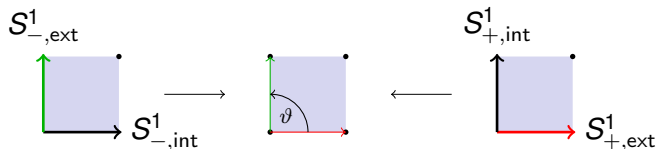
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There is a **torsion-free** G_2 -structure close to the one obtained by gluing

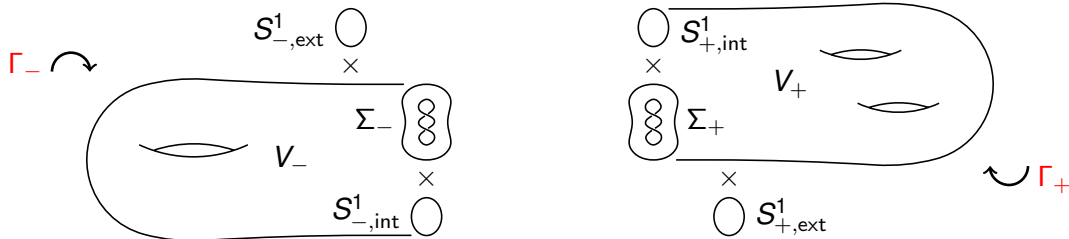
Recall twisted connected sums



Gluing of tori at angle $\vartheta = \frac{\pi}{2}$ between exterior circles



Extra twisted connected sums

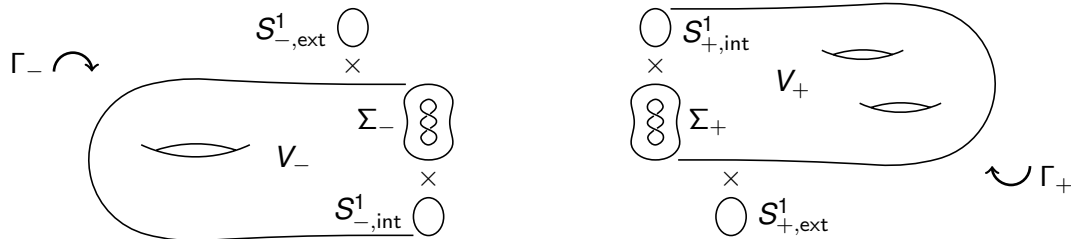


Assume that $\Gamma_{\pm} \cong \mathbb{Z}/k_{\pm}$ acts both on V_{\pm} and on $S^1_{\pm,ext}$

The induced action on ∂V_{\pm} has to fix Σ_{\pm} pointwise

The actions on $S^1_{\pm,int}$ and $S^1_{\pm,ext}$ have to be free

Extra twisted connected sums



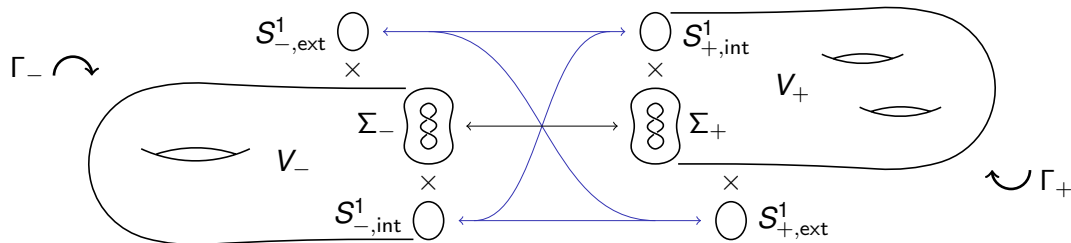
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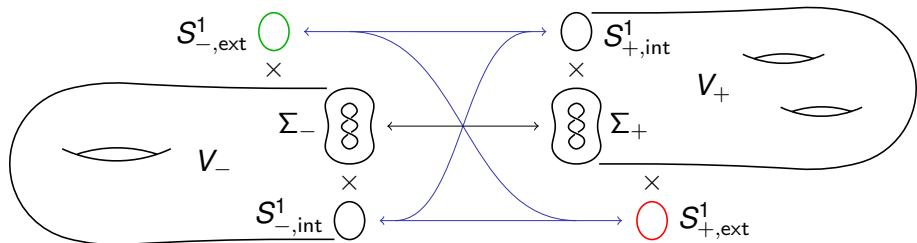
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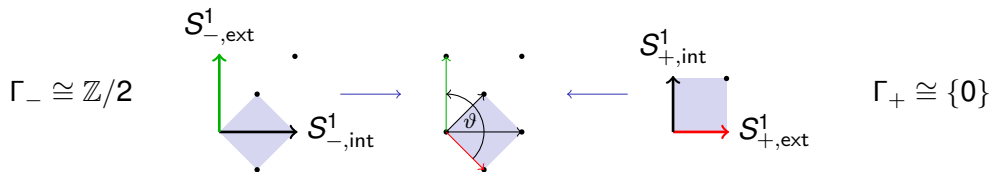
If both the tori and the K3 surfaces are isometric,

we can glue $M_{\pm} = (V_{\pm} \times S^1_{\pm, \text{ext}})/\Gamma_{\pm}$ at various angles ϑ

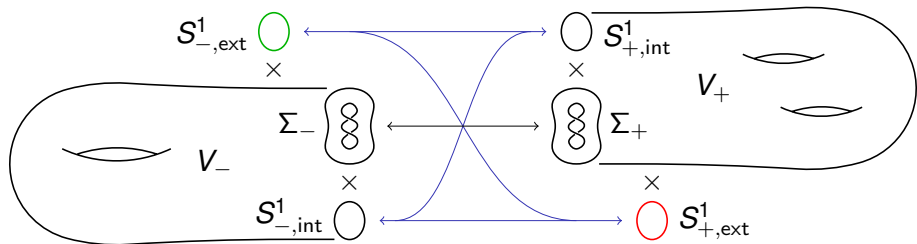
Extra twisted connected sums



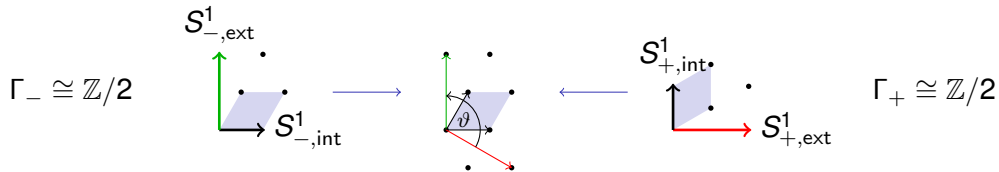
Modified gluing of tori at angle $\vartheta = \frac{3}{4}\pi$



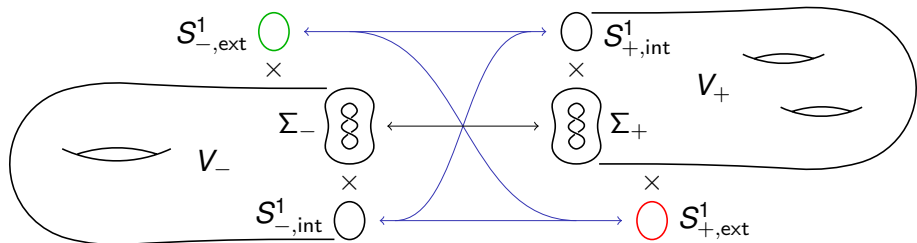
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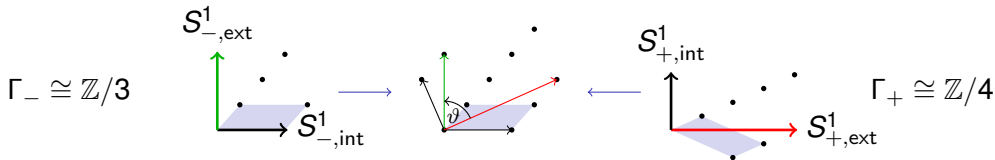
Modified gluing of tori at angle $\vartheta = \frac{2}{3}\pi$



Extra twisted connected sums



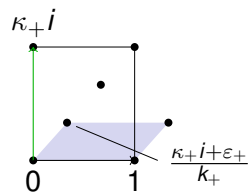
Modified gluing of tori at angle $\vartheta = \arccos\left(\frac{1}{\sqrt{6}}\right) \notin \mathbb{Q}\pi$



Assume that $\Gamma_{\pm} \cong \mathbb{Z}/k_{\pm}$
 acts on $\tilde{T} = S_{\pm, \text{int}}^1 \times \kappa_{\pm} S_{\pm, \text{ext}}^1$

A **torus matching** is described by

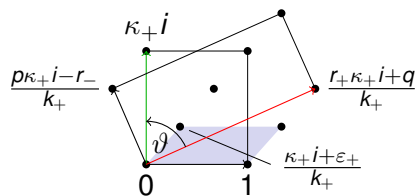
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 with $\det G = -k_+ k_-$ and $r_+ r_-, pq \geq 0$

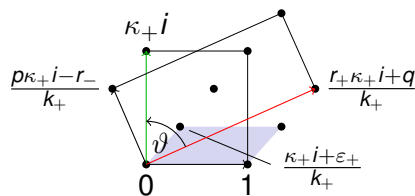


satisfying some extra conditions (only finitely many choices for ε_+ , G possible)

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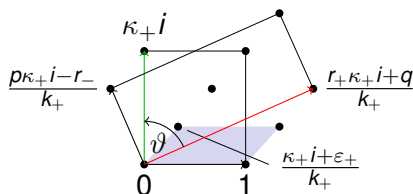
From G , recover

- ▶ The **aspect ratios** $\kappa_+ = \frac{\ell(S^1_{+, \text{ext}})}{\ell(S^1_{+, \text{int}})} = \sqrt{\frac{qr_-}{pr_+}}$ and $\kappa_- = \frac{\ell(S^1_{-, \text{ext}})}{\ell(S^1_{-, \text{int}})} = \sqrt{\frac{qr_+}{pr_-}}$
- ▶ The **gluing angle** $\vartheta = \arg(r_+ \kappa_+ + iq) \in (-\pi, \pi]$
- ▶ The **fundamental group** $\pi_1(M) \cong \mathbb{Z}/p$

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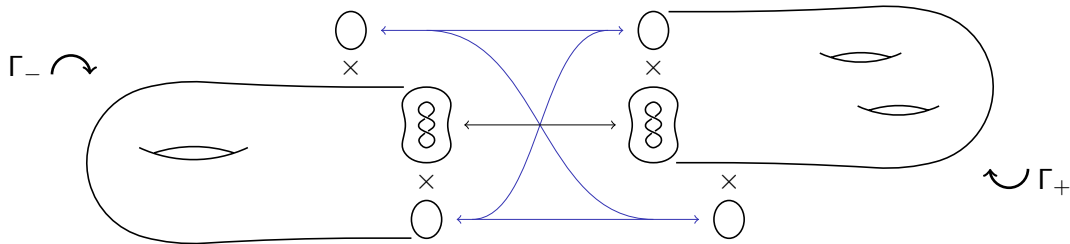


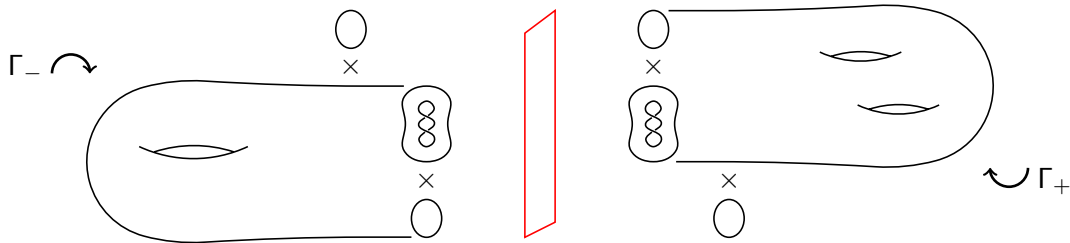
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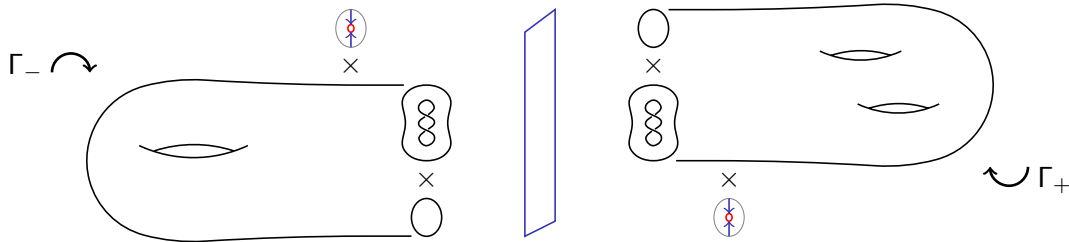
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There is a similar **K3 matching problem** involving a hyperkähler rotation by ϑ

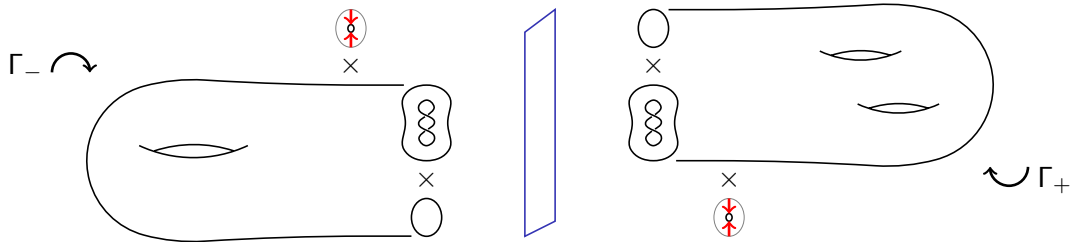




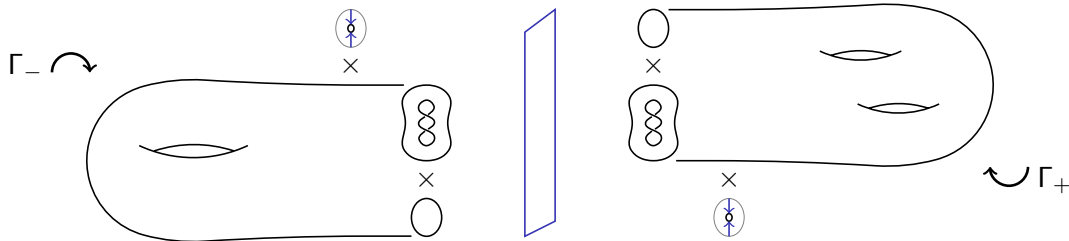
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Write $M = M_+ \cup_{M_0} M_-$ with $M_0 = \Sigma \times T^2$

Define Dirac operators D_{M_0} and B_{M_0} on M_0

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Using modified Atiyah-Patodi-Singer boundary conditions, define

$$\bar{\nu}(M_{\pm}, g) = 3\eta_{\text{APS}}(B_{M_{\pm}}; L_{B,\pm}) - 24\eta_{\text{APS}}(D_{M_{\pm}}; L_{D,\pm}) \in \mathbb{R}$$

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There exists $m(L_{B,+}, L_{B,-}) \in \mathbb{Z}$ depending only on the K3 matching such that

$$\bar{\nu}(M, g) = \bar{\nu}(M_+, g) + \bar{\nu}(M_-, g) + 144 \frac{\vartheta}{\pi} - 72 + 3m(L_{B,+}, L_{B,-})$$

Recall that the gluing angle ϑ depends only on the gluing matrix

Compute adiabatic limit ($r \rightarrow 0$) of η -invariants on manifolds with boundary
We use the adiabatic limit theorems of [Bismut-Cheeger '92, Dai '01, G '14]

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Except over the singular set of V_{\pm}/Γ_{\pm} ,
 the manifold M_{\pm} is locally isometric to a product

The adiabatic limit depends only on isolated fixpoints of elements of Γ_{\pm}

Write

$$\begin{aligned}
 D_{\gamma_{\pm}}(V_{\pm}) &= \lim_{r \rightarrow 0} \bar{\nu}(M_{\pm, r}) \\
 &= \frac{3}{k_{\pm}} \sum_{j=1}^{k_{\pm}-1} \cot \frac{\pi j}{k_{\pm}} \sum_{p \in V_{\pm, j}} \frac{\cos \frac{\alpha_{j,1}(p)}{2} \cos \frac{\alpha_{j,2}(p)}{2} \cos \frac{\alpha_{j,3}(p)}{2} - 1}{\sin \frac{\alpha_{j,1}(p)}{2} \sin \frac{\alpha_{j,2}(p)}{2} \sin \frac{\alpha_{j,3}(p)}{2}} \in \mathbb{Q}
 \end{aligned}$$

Shrink S^1 -fibres of the flat orbibundle $M_{\pm} \rightarrow V_{\pm}/\Gamma_{\pm}$ to 0

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Let $\tilde{\eta}(\mathbb{A})$ be the η -form of the family of tori $(S^1_{\pm,\text{int}} \times rS^1_{\pm,\text{ext}})/\Gamma_{\pm}$ for $r \in (0, \kappa_{\pm})$

The boundary contribution is given by

$$\bar{\nu}(M_{\pm}) - \lim_{r \rightarrow 0} \bar{\nu}(M_{\pm,r}) = -\frac{144}{\pi} F_{\kappa_{\pm}, \varepsilon_{\pm}}(\kappa_{\pm}) = 288 \int_0^{\kappa_{\pm}} \tilde{\eta}(\mathbb{A})$$

Here, κ_{\pm} are the aspect ratios of the tori and depend only the gluing matrix

Represent the torus $(S_{\pm, \text{int}}^1 \times rS_{\pm, \text{ext}}^1)/\Gamma_{\pm}$ by $\tau \in \mathcal{H} = \{z \in \mathbb{C} \mid \text{Im}(z) > 0\}$
Then $\tilde{\eta}(\mathbb{A}) \in \Omega^1(\mathcal{H})$ is $SL(2, \mathbb{Z})$ -invariant and satisfies $d\tilde{\eta}(\mathbb{A}) = -\frac{1}{4\pi} dA_{\text{hyp}}$
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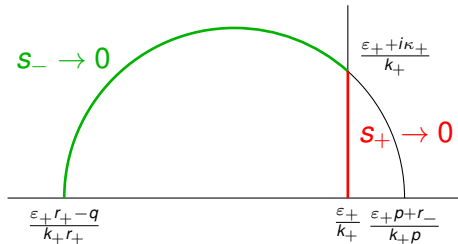
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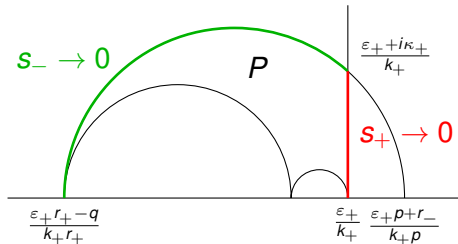
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$\tilde{\eta}(\mathbb{A}) = 0$ along rectangular families

Determine P using continued fractions



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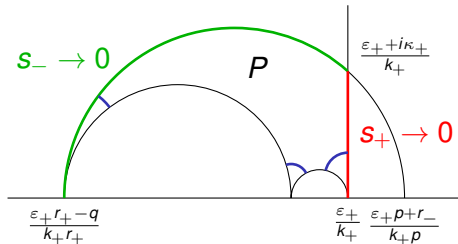
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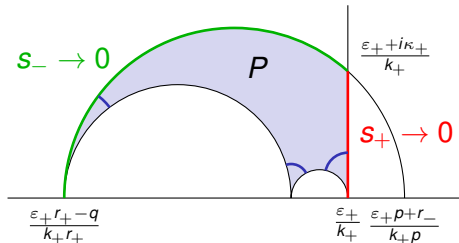
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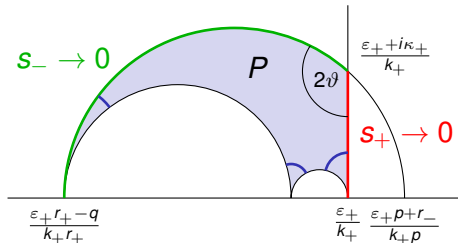
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Use the hyperbolic area formula for P

The cusp contributions add up to a classical Dedekind sum [Zagier '75]

The angle 2ϑ at the finite corner cancels $144 \frac{\vartheta}{\pi}$ in the gluing formula



An alternative approach uses Dedekind's η -function

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Compute the sum of the variational terms using the functional equations

$$H(\tau + 1) = \frac{\pi i}{12} + H(\tau) \quad \text{and} \quad H\left(-\frac{1}{\tau}\right) = \frac{1}{2} \log \frac{\tau}{i} + H(\tau)$$

The correction terms add up to give a classical Dedekind sum

Let M be an extra twisted connected sum
with gluing matrix $G = \begin{pmatrix} r_+ & p \\ q & -r_- \end{pmatrix}$ and gluing angle ϑ
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For integers $n > 0$ and a define the classical Dedekind sum

$$S(a, n) = \sum_{j=1}^{n-1} \left(\left(\frac{j}{n} \right) \right) \left(\left(\frac{aj}{n} \right) \right) \in \frac{1}{6n} \mathbb{Z} \quad \text{with} \quad ((x)) = \begin{cases} 0 & x \in \mathbb{Z} \\ x - [x] - \frac{1}{2} & x \notin \mathbb{Z} \end{cases}$$

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Theorem (G-Nordström)

Assume $q > 0$. Then $a = \frac{r_+ - \varepsilon_+^* q}{k_+} \in \mathbb{Z}$ and

$$\bar{\nu}(M, g) = D_{\gamma_+}(V_+) + D_{\gamma_-}(V_-) + 3m(L_{B_+}, L_{B_-}) + 24 \left(12 S(a, q) - \frac{r_-}{k_- q} - \frac{r_+}{k_+ q} \right)$$

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Better maybe: $D_{\gamma_+}(V_+) - 24 \frac{\varepsilon_+^*}{k_+}, D_{\gamma_-}(V_-) - 24 \frac{\varepsilon_-^*}{k_-} \in \mathbb{Z}$

But ε_{\pm}^* is not well defined—and what does this mean?

Thanks for your attention!